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# Nonlinear Electron Resonance Heating in Asymmetric Capacitive Discharges

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**Abstract**: In low-pressure capacitive discharges the concept of Nonlinear Electron Resonance Heating (NERH) becomes important to enhance Ohmic dissipation. Particularly in geometrically asymmetric capacitive discharges (generation of a DC self-bias), the nonlinearities of the boundary plasma sheaths lead to a strongly non-sinusoidal radio frequency current. The Fourier spectra of such a current can indicate harmonics which are in resonance with the Plasma Series Resonance (PSR). The scenario of PSR was investigated by different models (e.g. zero-dimensional, spatially resolved), which assume that the bulk current is carried by the electron current and the sheath current by the displacement current. Although these models are able to resolve the nonlinear behavior between bulk and sheath a detailed kinetic picture is missing. In this work, we discuss the particle dynamics in the regime of NERH on a nanosecond timescale by means of a self-consistent kinetic simulation. We use a 1d3v spherical Particle-In-Cell code in order to simulate an asymmetric discharge. It is shown that the excitation of harmonics is connected to the generation of multiple electron beams accelerated by the expanding plasma sheath. Furthermore, the interaction of these beams with bulk electrons leads to significant plasma oscillations and thus, to a moderate displacement current in the center of the discharge. These kinetic effects should be taken into account for future models in order to understand the comprehensive electron heating in capacitive discharges.

## Asymmetric CCRF Discharge



## **Particle-In-Cell Simulation**

### • 1d3v PIC simulation, spherical shell model[2][4]

### **Parameter Variation**

↑ voltage ⇒ ↑ ionization ⇒ ↑ n<sub>e</sub> ⇒ ↑ ω<sub>pe</sub> ⇒ ↓ τ<sub>pe</sub> = 2π/ω<sub>pe</sub>
electrons respond to local perturbations on faster timescales
excitation of higher harmonics is possible
higher pressure damps the nonlinear excitation
sheath edge criterion [5]: \$\int\_{0}^{s(t)} n\_i(x) dx = \int\_{s(t)}^{\infty} (n\_i(x) - n\_e(x)) dx\$

#### Current and voltage at the electrode







- argon chemistry: 3 electron-neutral (elastic, excitation, ionization) and 2 ion-neutral (isotropic, backward elastic scattering) collisions
- base case setup:
- $-V_{rf} = 300 \text{ V}, \ \omega_{rf} = 2\pi \cdot 13.56 \text{ MHz}$ , p = 1 Pa = 7.5 mTorr
- $-r_i = 20$  mm,  $r_a = 100$  mm,  $r_a r_i = 80$  mm gap size
- electrode area ratio  $A_a/A_i \approx 25.01$ ,  $C_b = 0.3$  nF

### **Base Case**



#### • significant charge densities in the bulk due to electron modulation (ions constant over time)



ullet harmonic oscillation in the rf current  $\Longrightarrow$  excitation of plasma series resonance

• stronger oscillation for higher voltages, oscillations are damped for higher pressures

#### Fourier spectra at the electrode



1 Pa, 600 V

Time (ns)

#### Electrons above 15.76 eV in m<sup>-3</sup>



#### Charge density in $\mu$ Asm<sup>-3</sup>





• electron beam penetrates into the bulk and cold bulk electrons are accelerated back to the expanding sheath



#### • charge density in the bulk region vanishes for higher pressures

### Conclusion

kinetic analysis is required in order to understand local kinetic phenomena which lead to resonance effects
local interaction between energetic beam electrons and cold bulk electrons ⇒ (bulk) electron-sheath interaction [6]
displacement current in the bulk region exist due to this interaction

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