



Nonlocal and Nonlinear Electron Dynamics in Capacitively Coupled Radio Frequency Discharges

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[R. P. Brinkmann](#)¹, [Z. Donkó](#)⁴, and [T. Mussenbrock](#)²

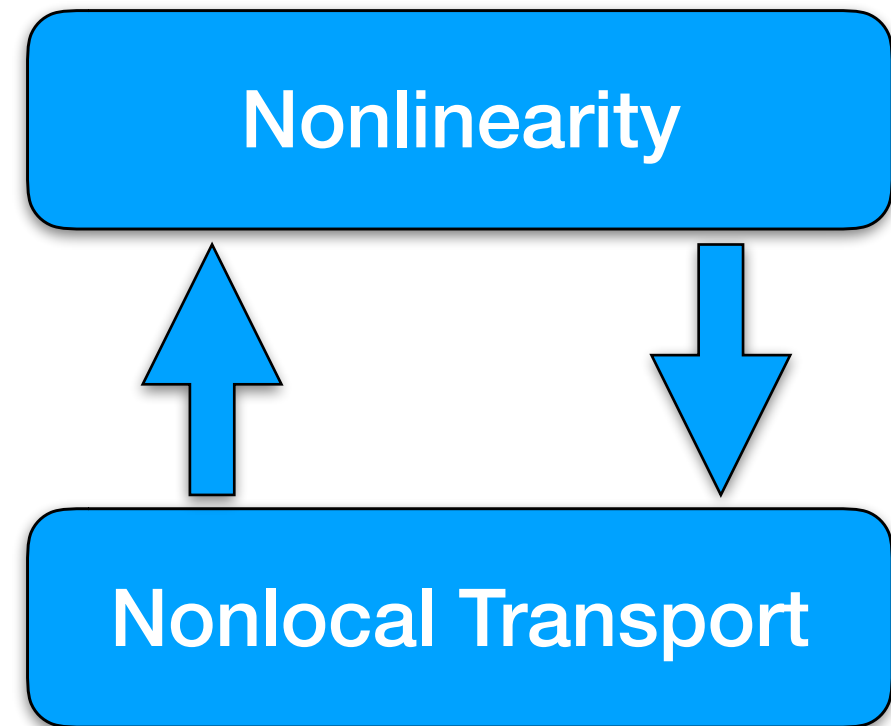
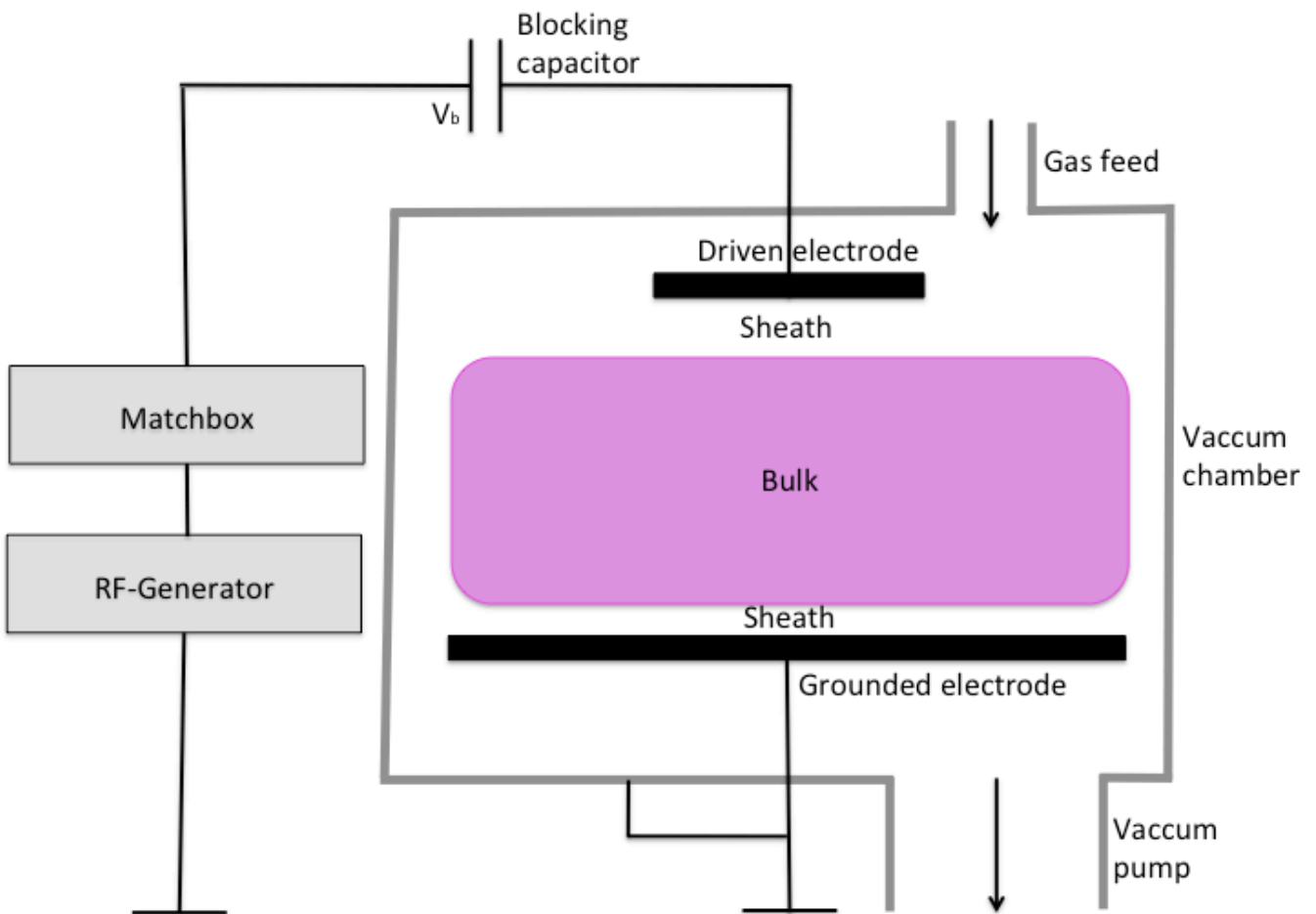
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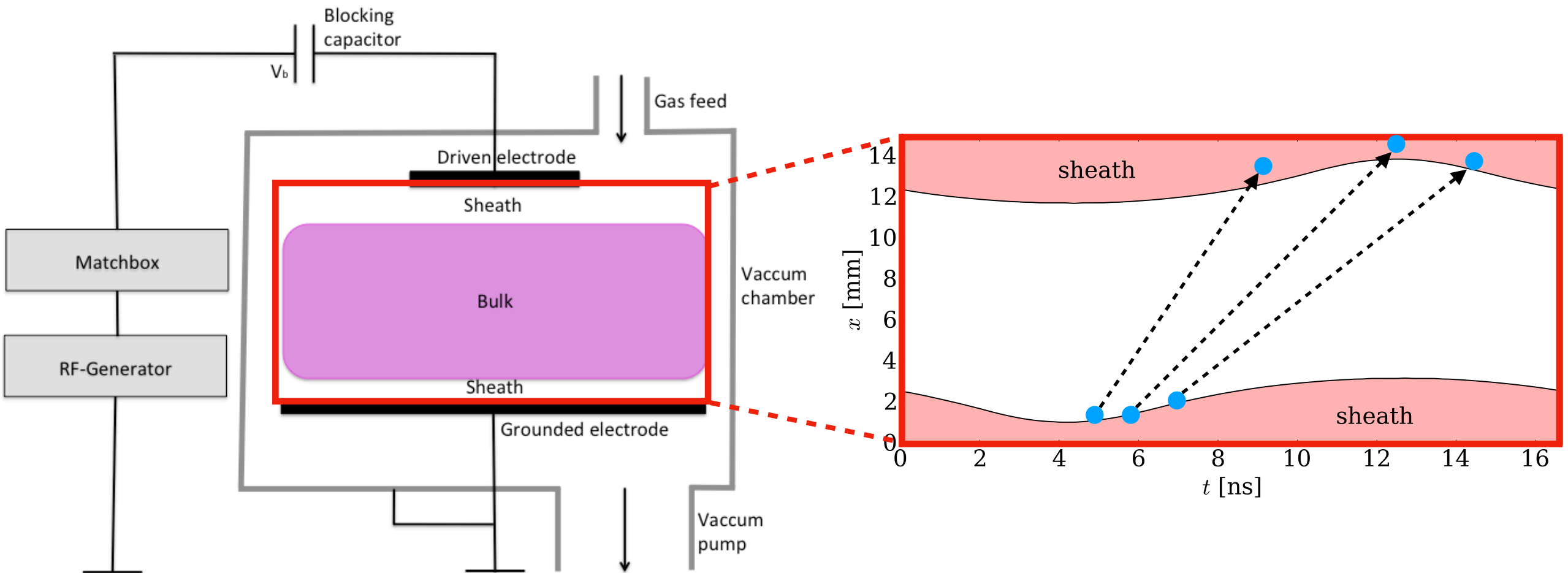
Motivation: Electron Dynamics



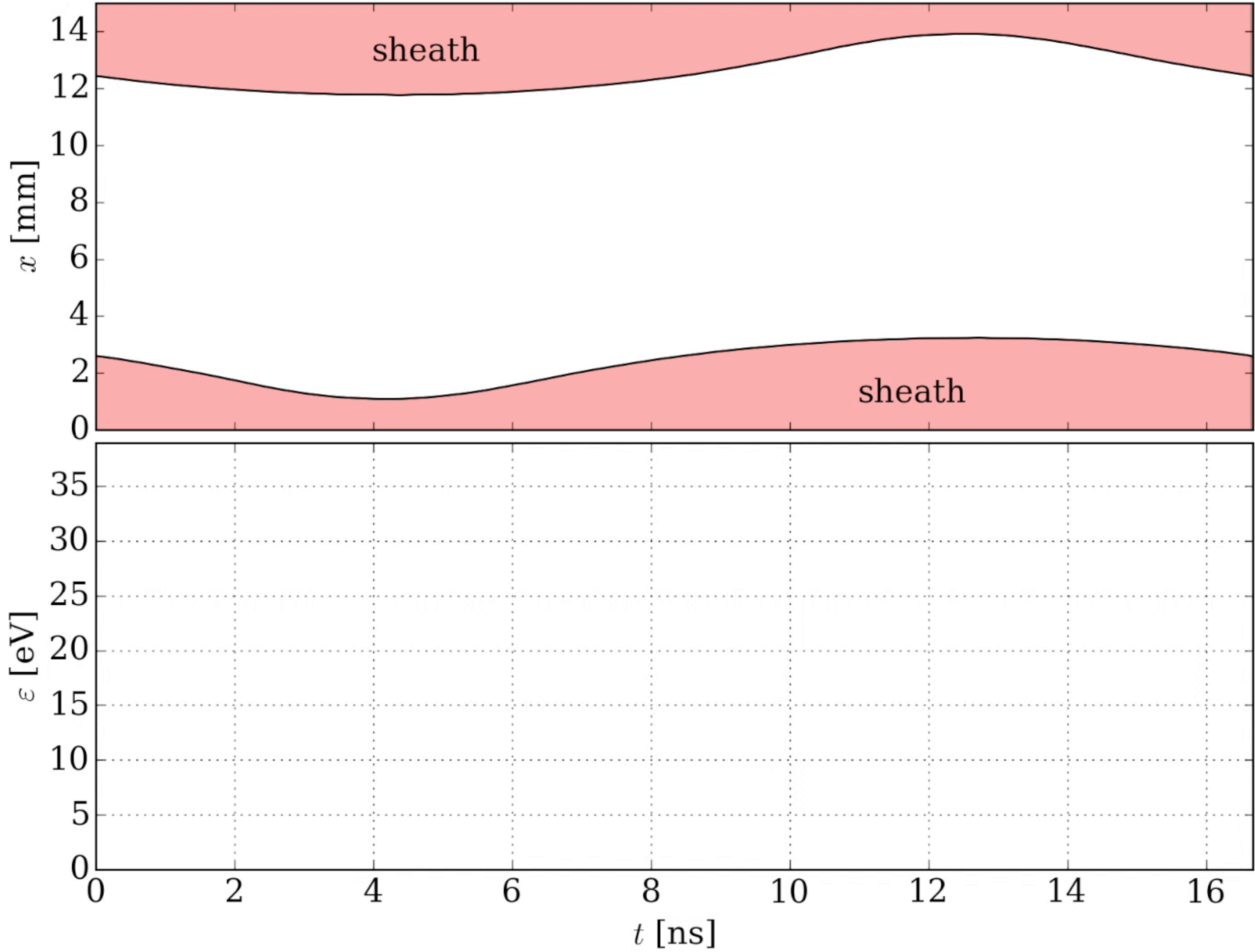
- classical CCRF discharge for etching and sputtering processes
- optimization and control of the plasma density and the ion flux to the wall
- fundamental phenomena are still not fully understood
- electron power gain and loss at low pressures on a kinetic level

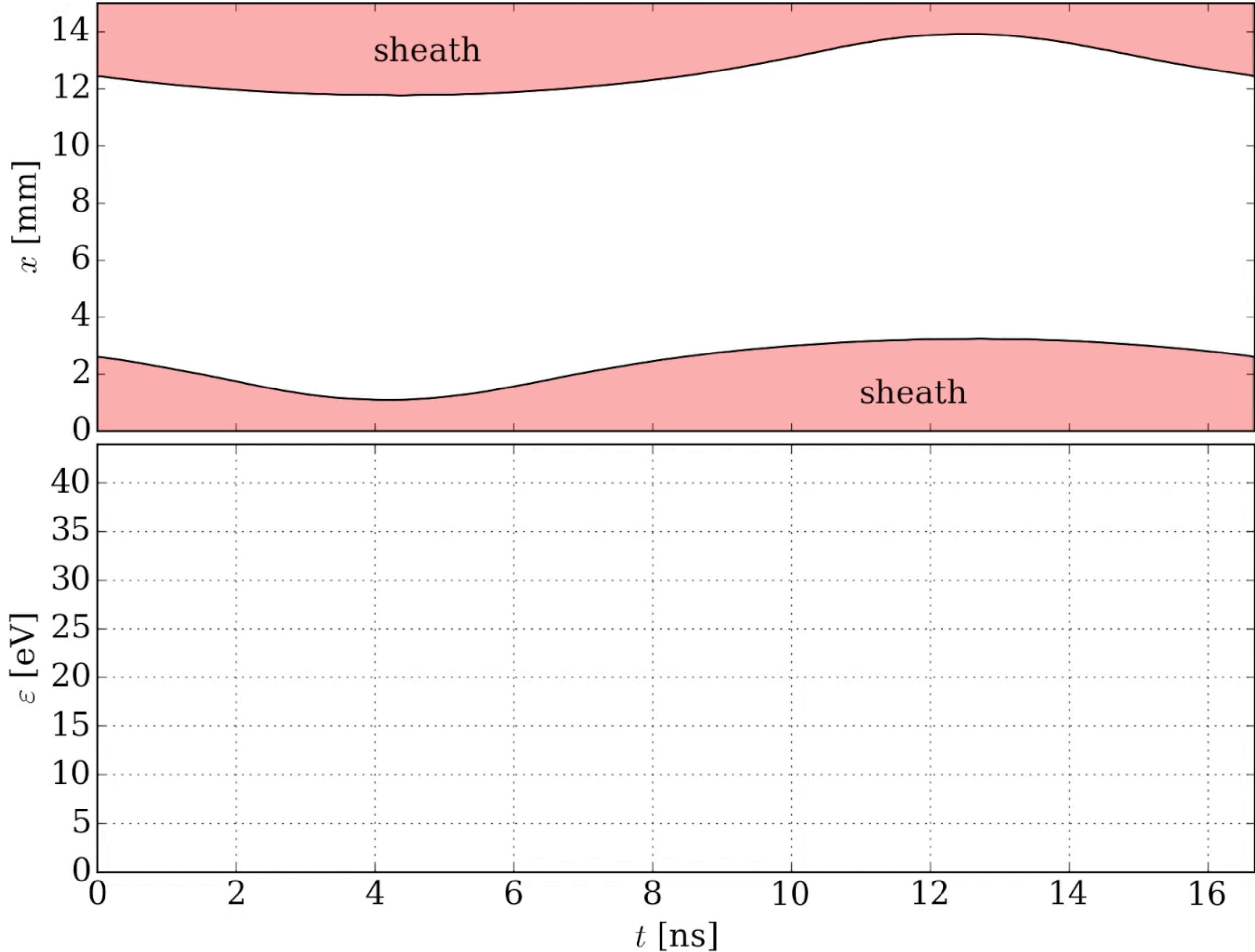
Nonlocal Transport

Motivation: Nonlocal Transport

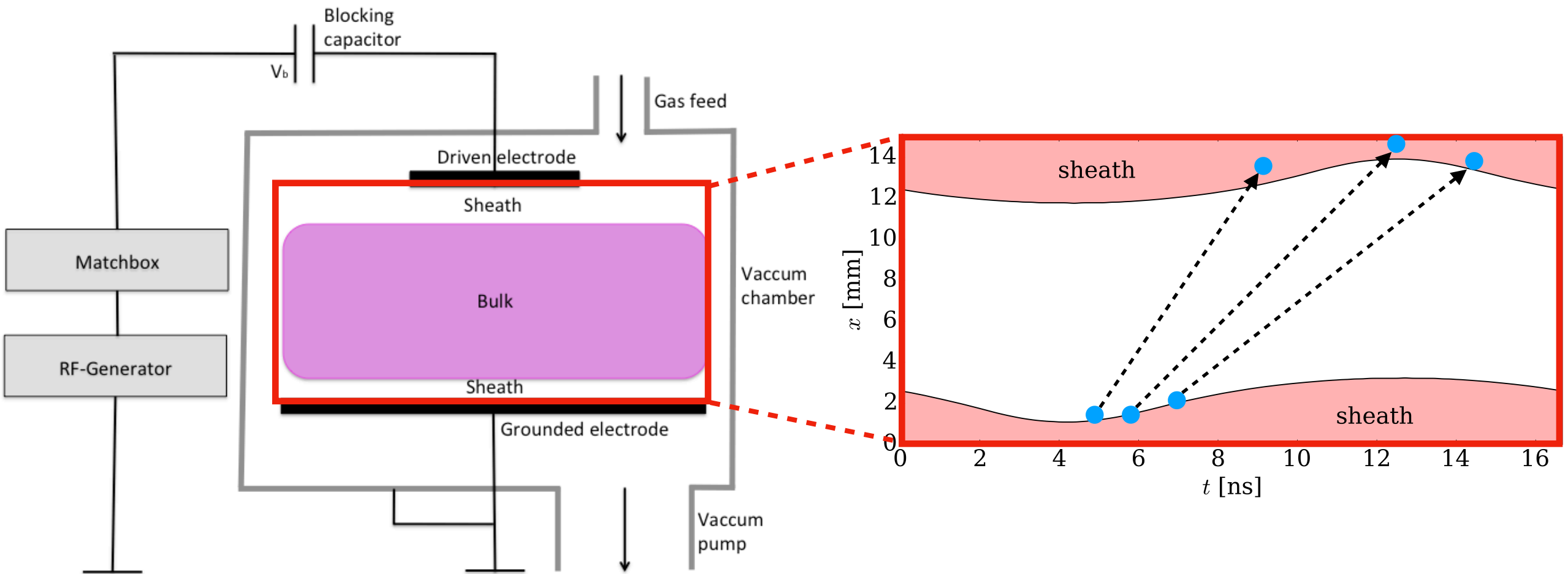


- low pressure (1 Pa), electron mean free path $>$ electrode gap length
- electrons transport their energy almost collisionslessly through the discharge
- confinement and reflection of beam electrons become crucial



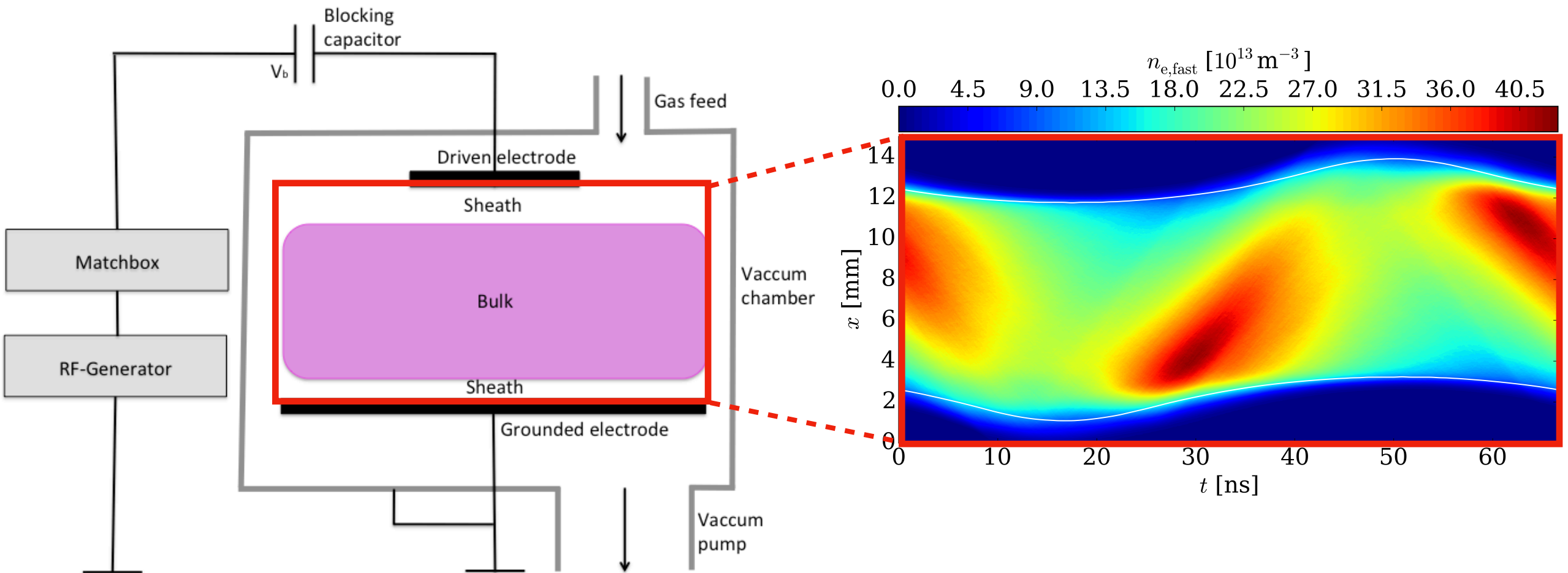


Goal of this work: Nonlocal Transport



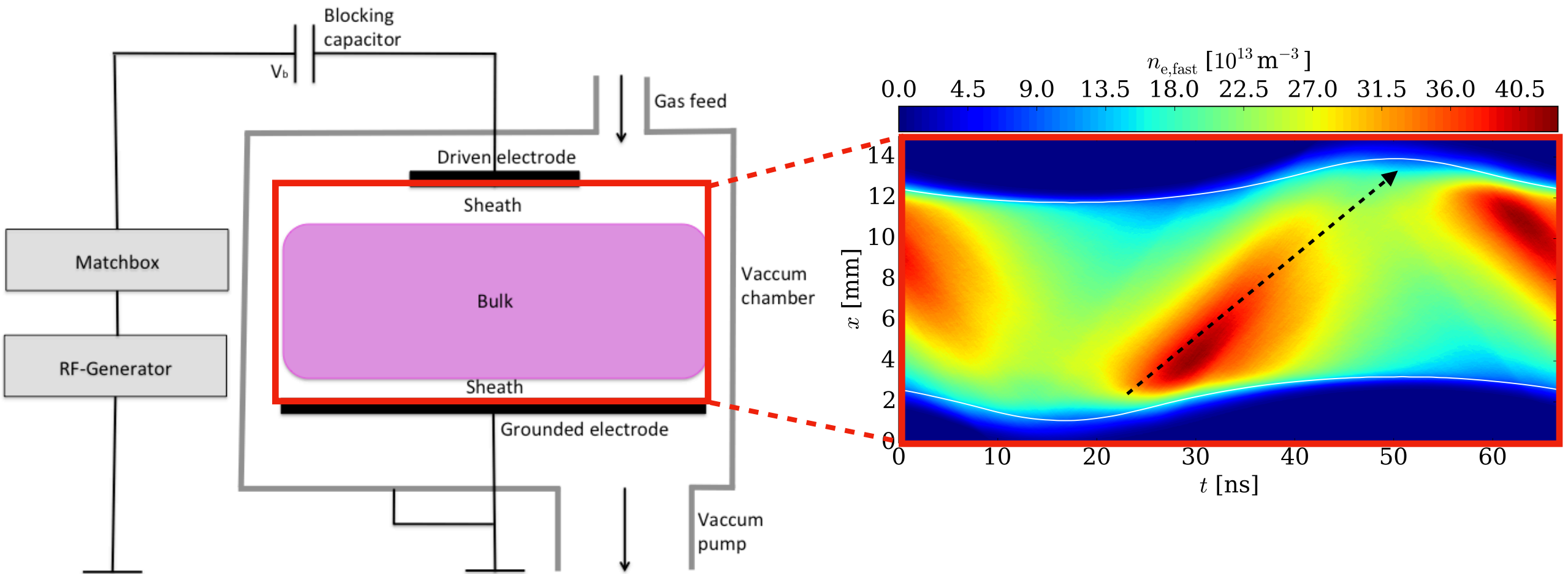
- can we control the trajectory and the impingement phase of energetic electrons?

Goal of this work: Nonlocal Transport



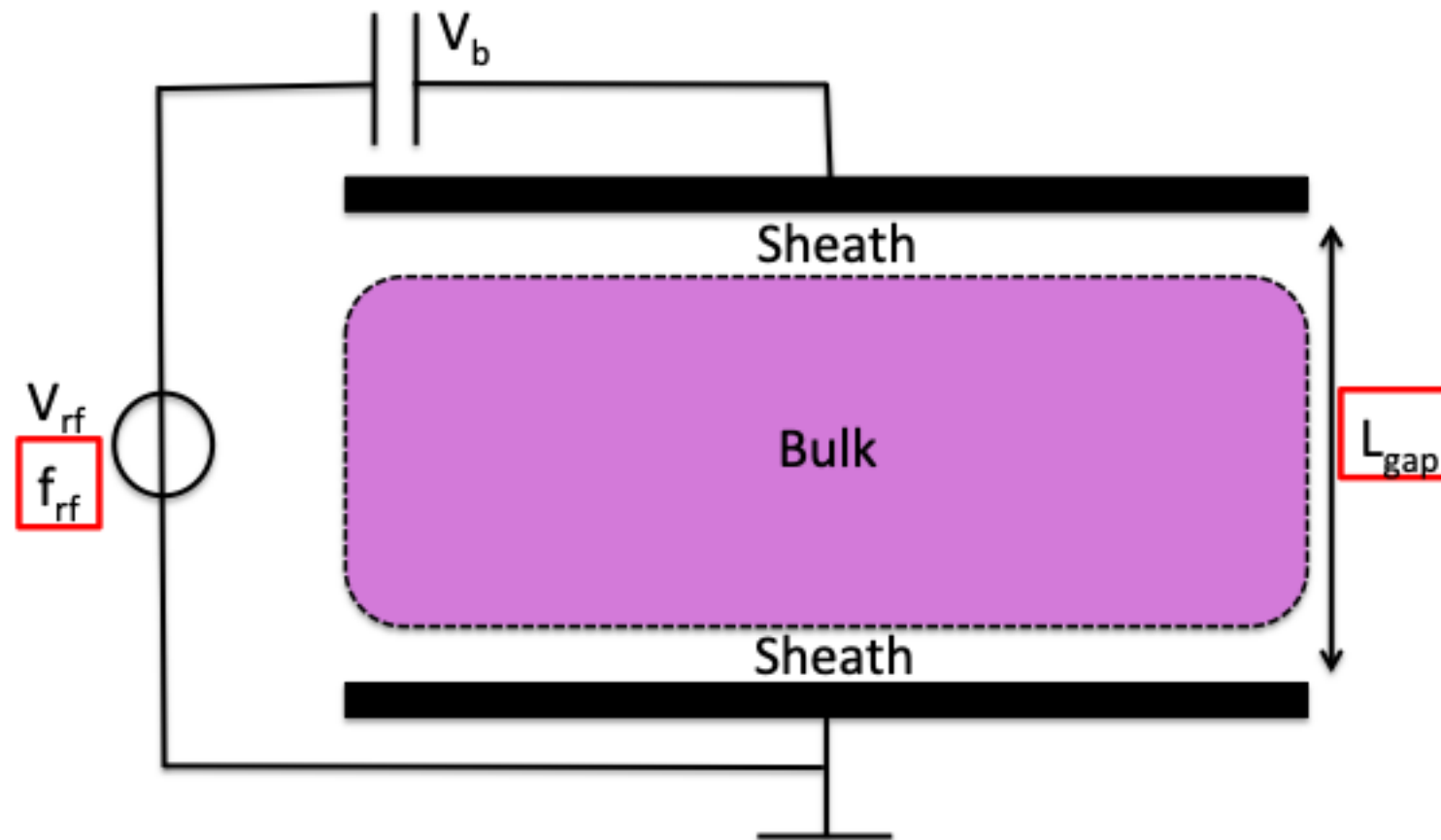
- can we control the trajectory and the impingement phase of energetic electrons?
- electron density of fast electrons above 15.7 eV (comparable to PROES)

Goal of this work: Nonlocal Transport



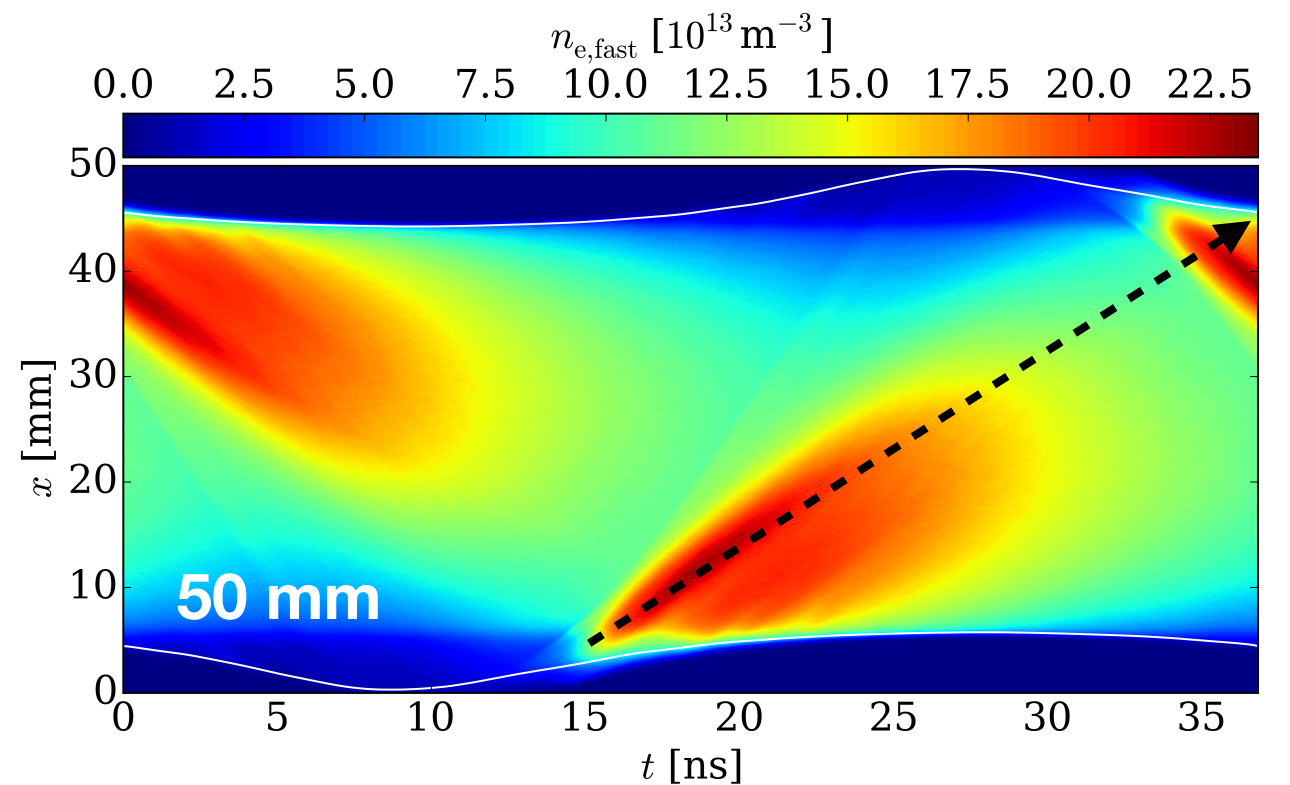
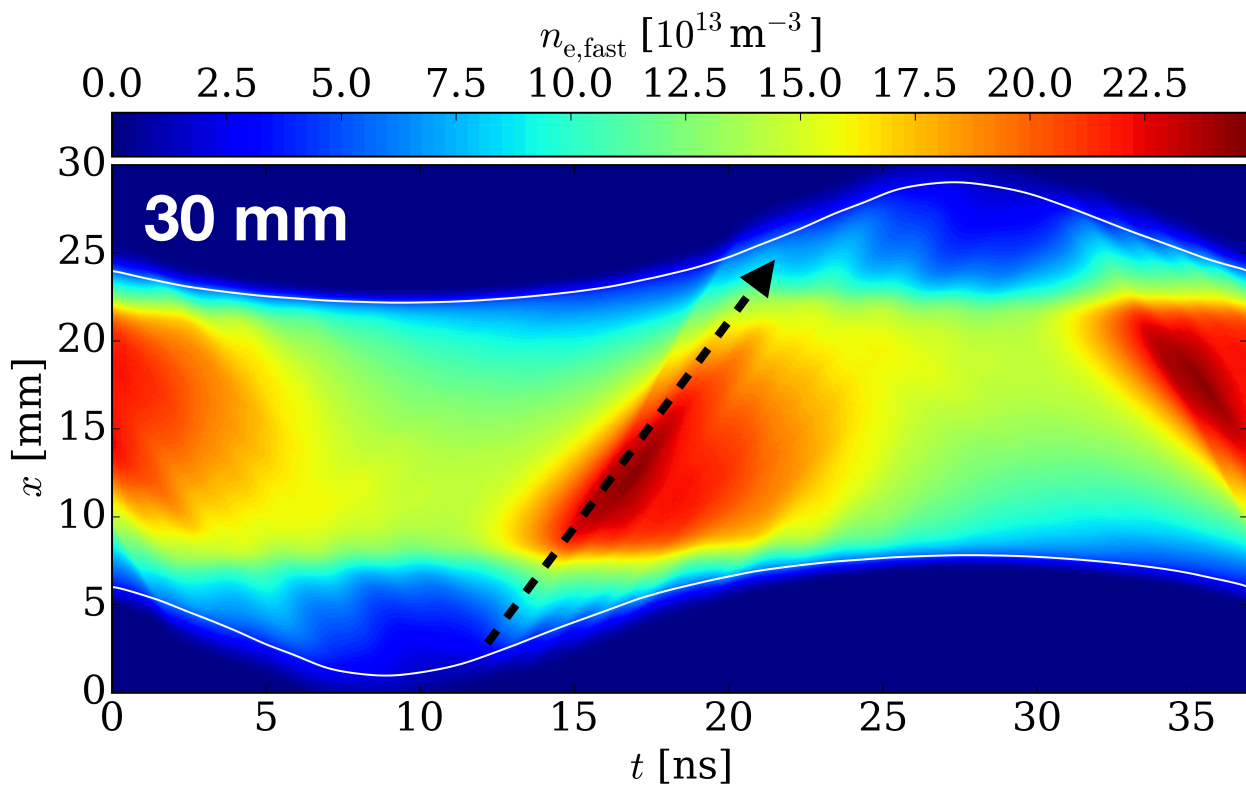
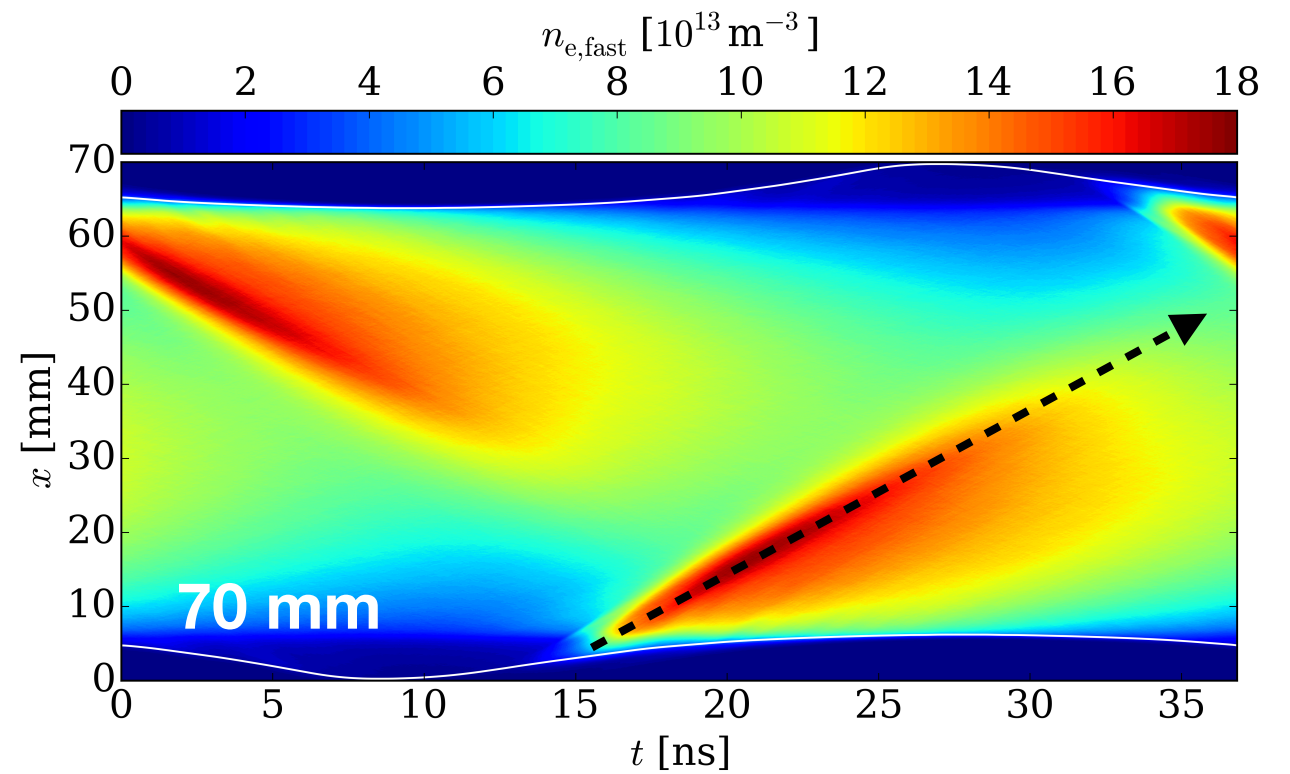
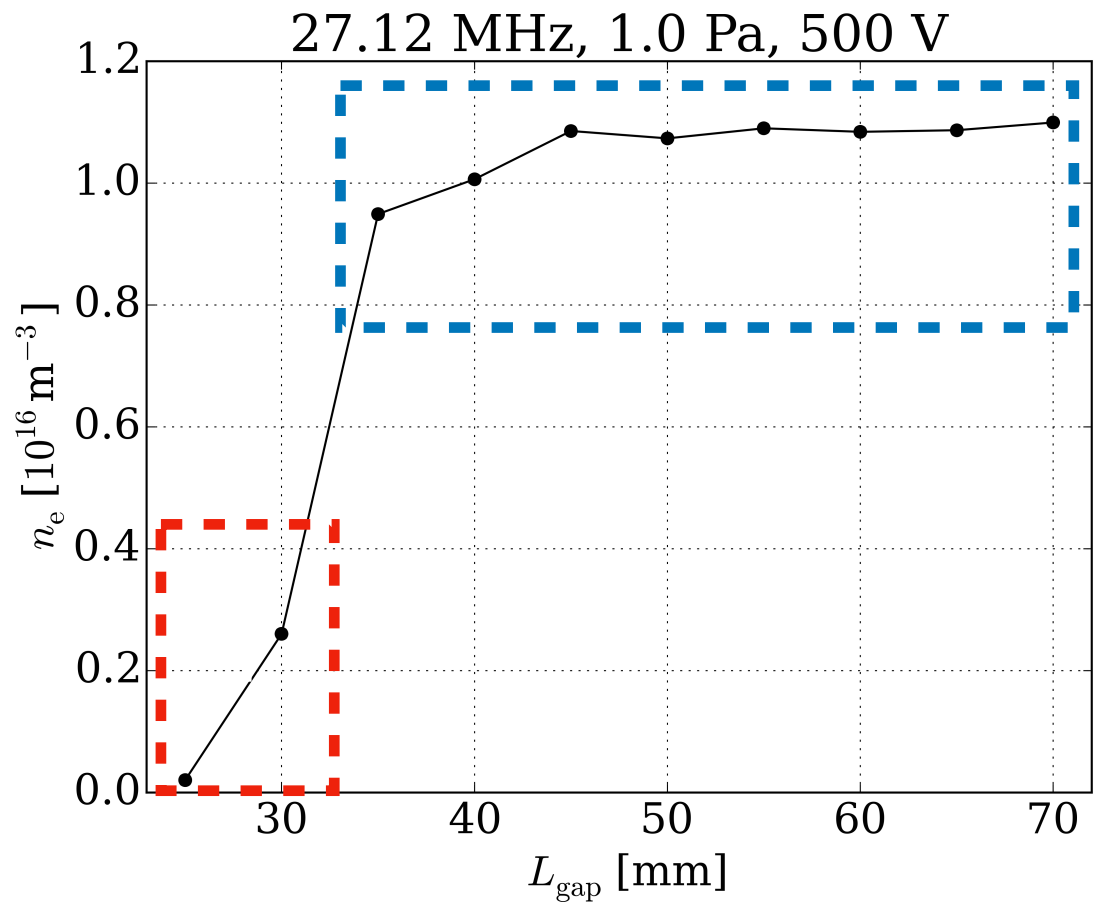
- can we control the trajectory and the impingement phase of energetic electrons?
- electron density of fast electrons above 15.7 eV (comparable to PROES)
- how does input parameter affect the impingement phase?
 - gap size variation
 - driving frequency variation

Gap variation



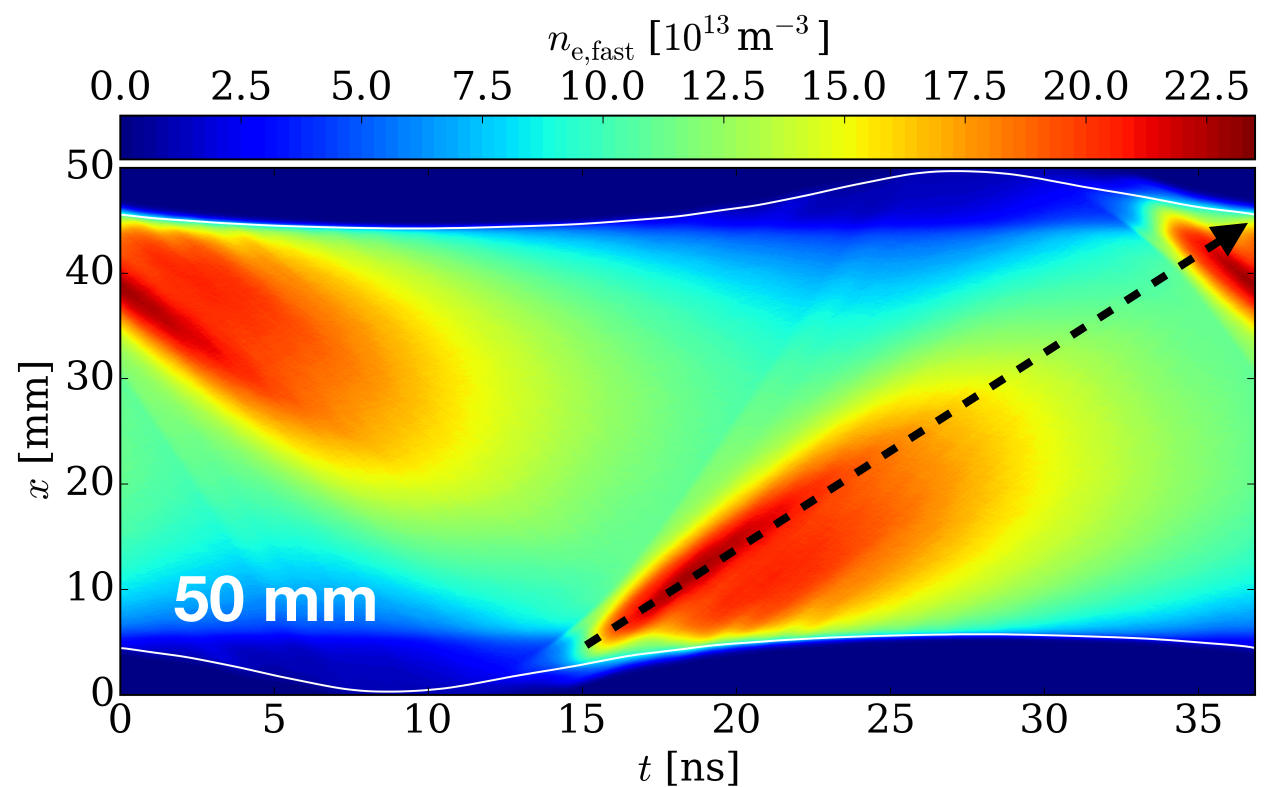
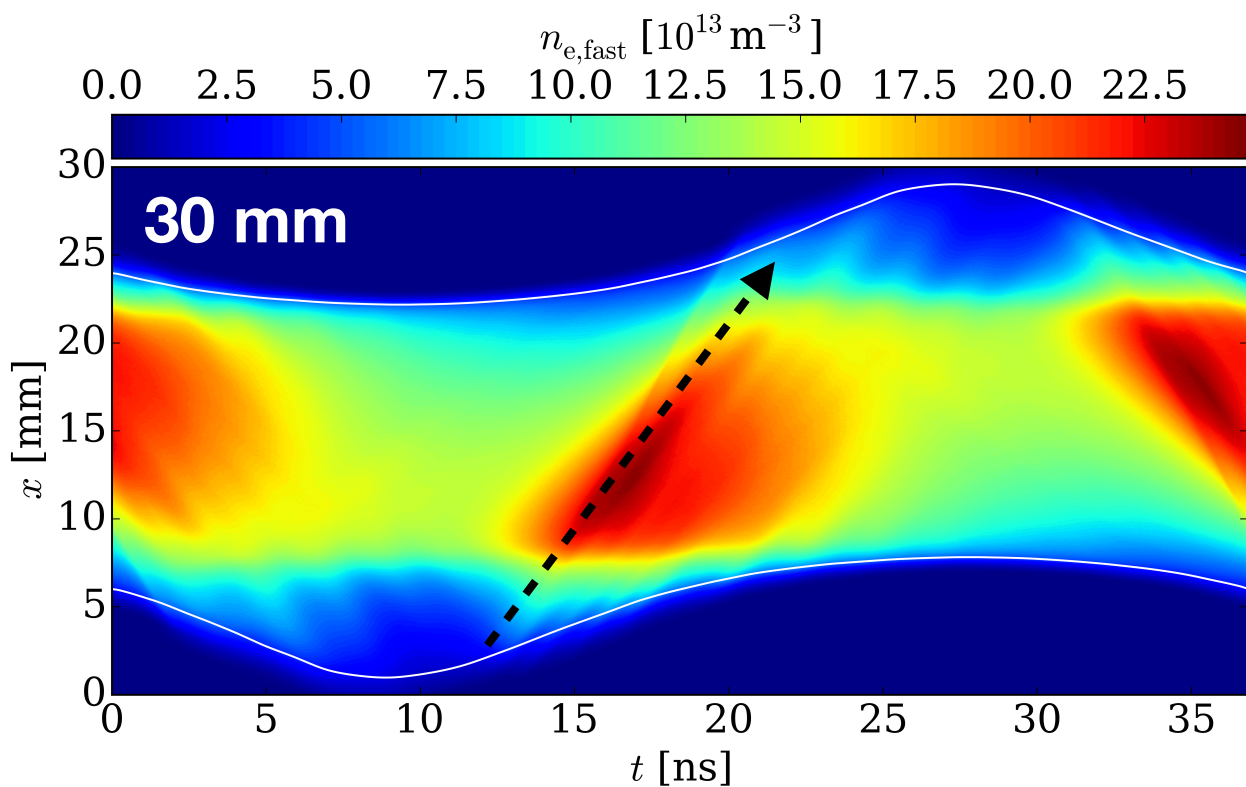
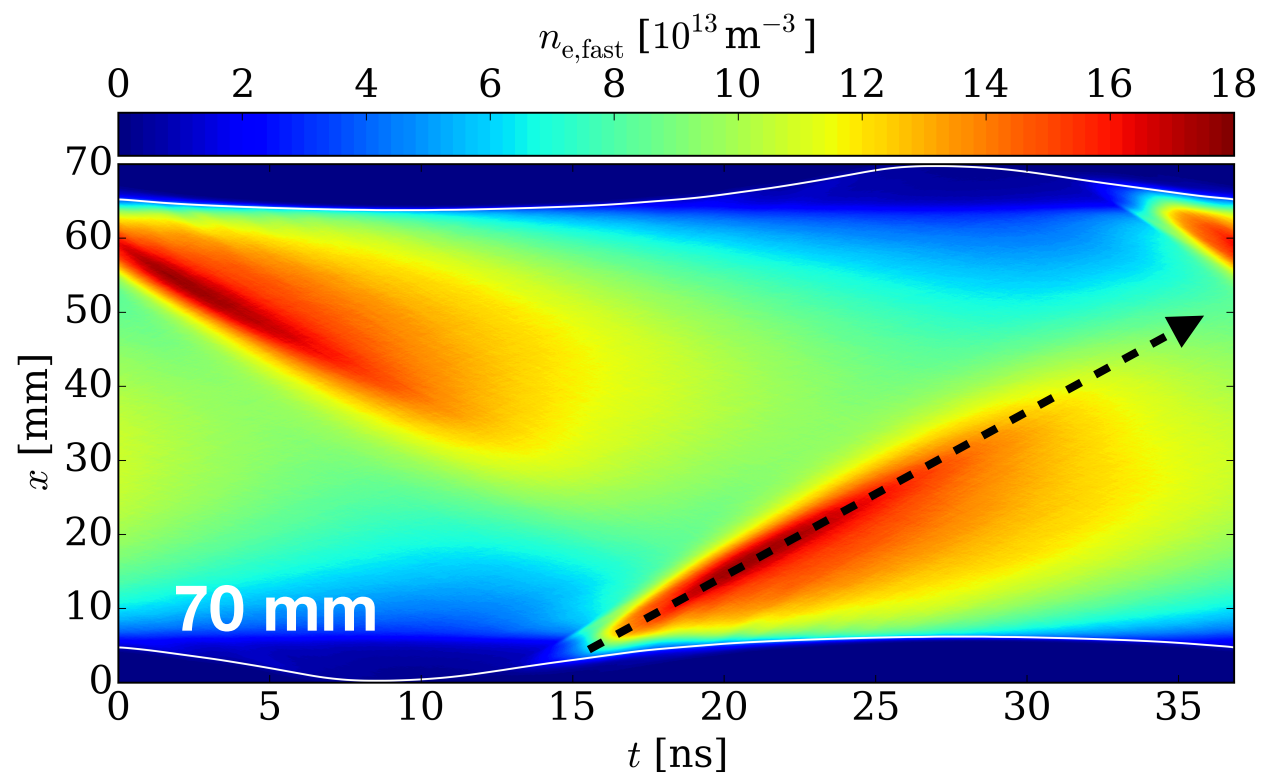
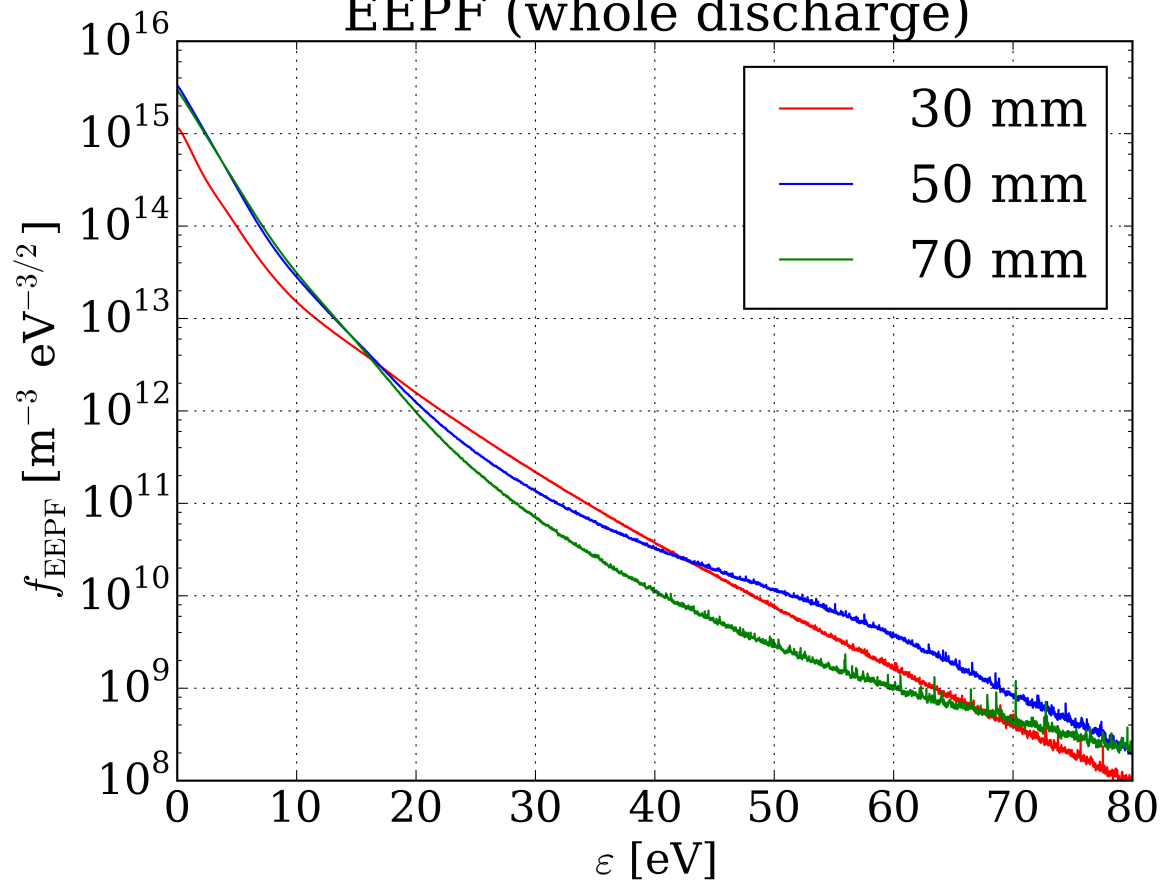
- 1d3v Particle-In-Cell simulation, Cartesian grid (perfectly symmetric), argon
- no secondary electron emission and no reflection at the boundary
- frequency: 27.12 MHz
- gas pressure: 1 Pa
- driving voltage: 500 V
- **gap size: 25 - 70 mm**

Gap variation: Fast electrons ($E > 15.7$ eV)



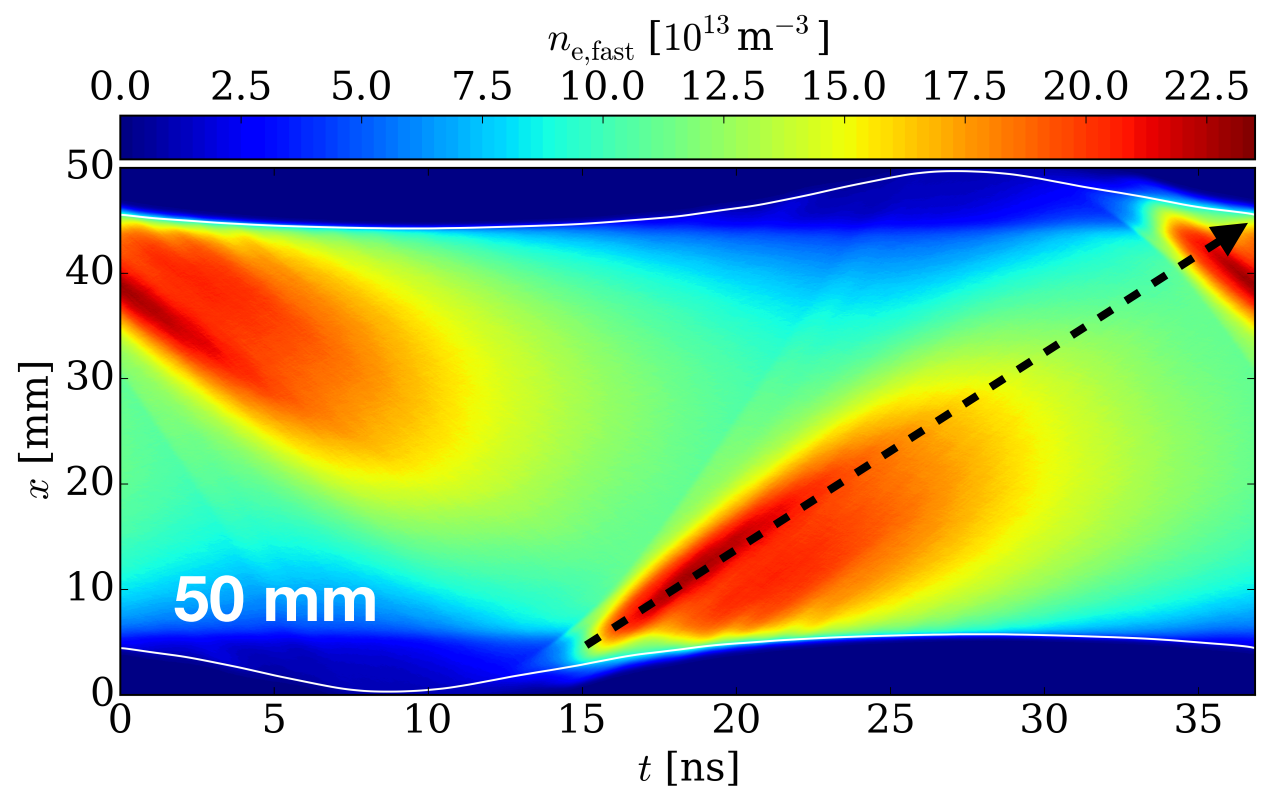
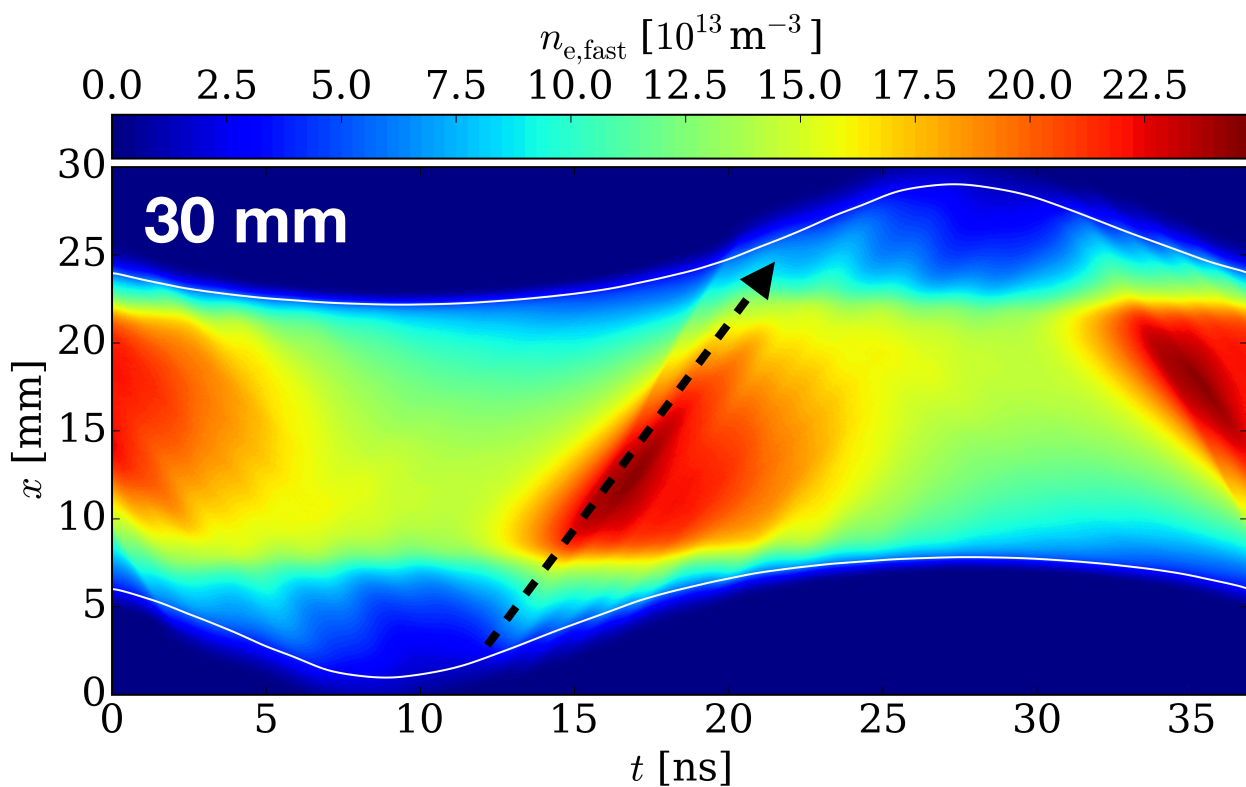
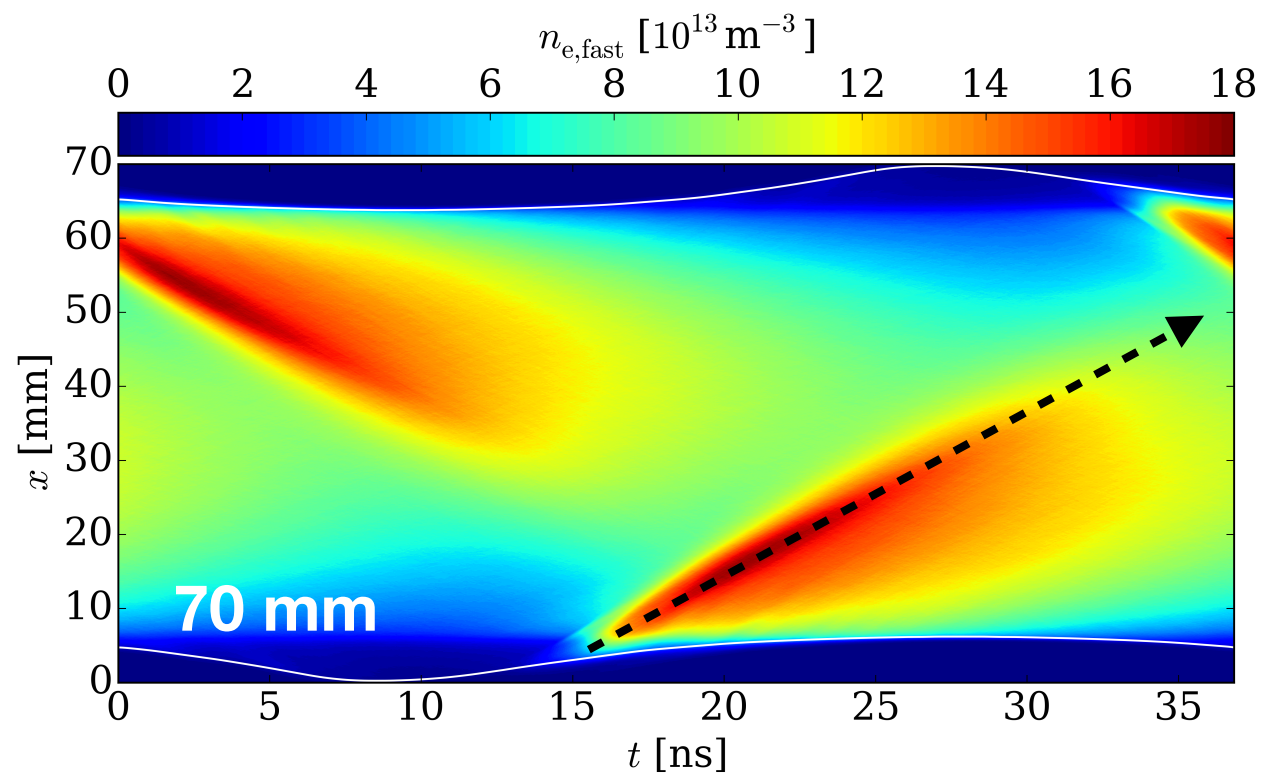
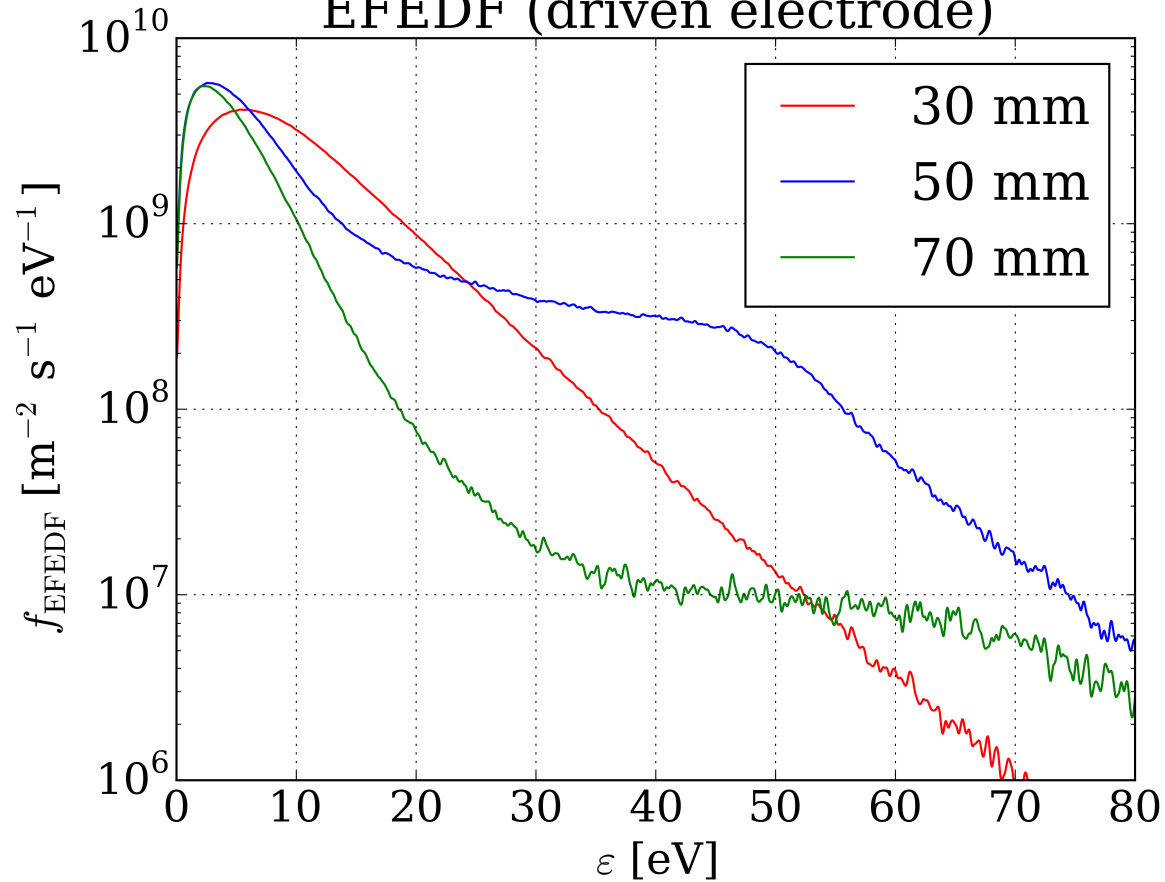
Gap variation: Distribution function

EEPF (whole discharge)

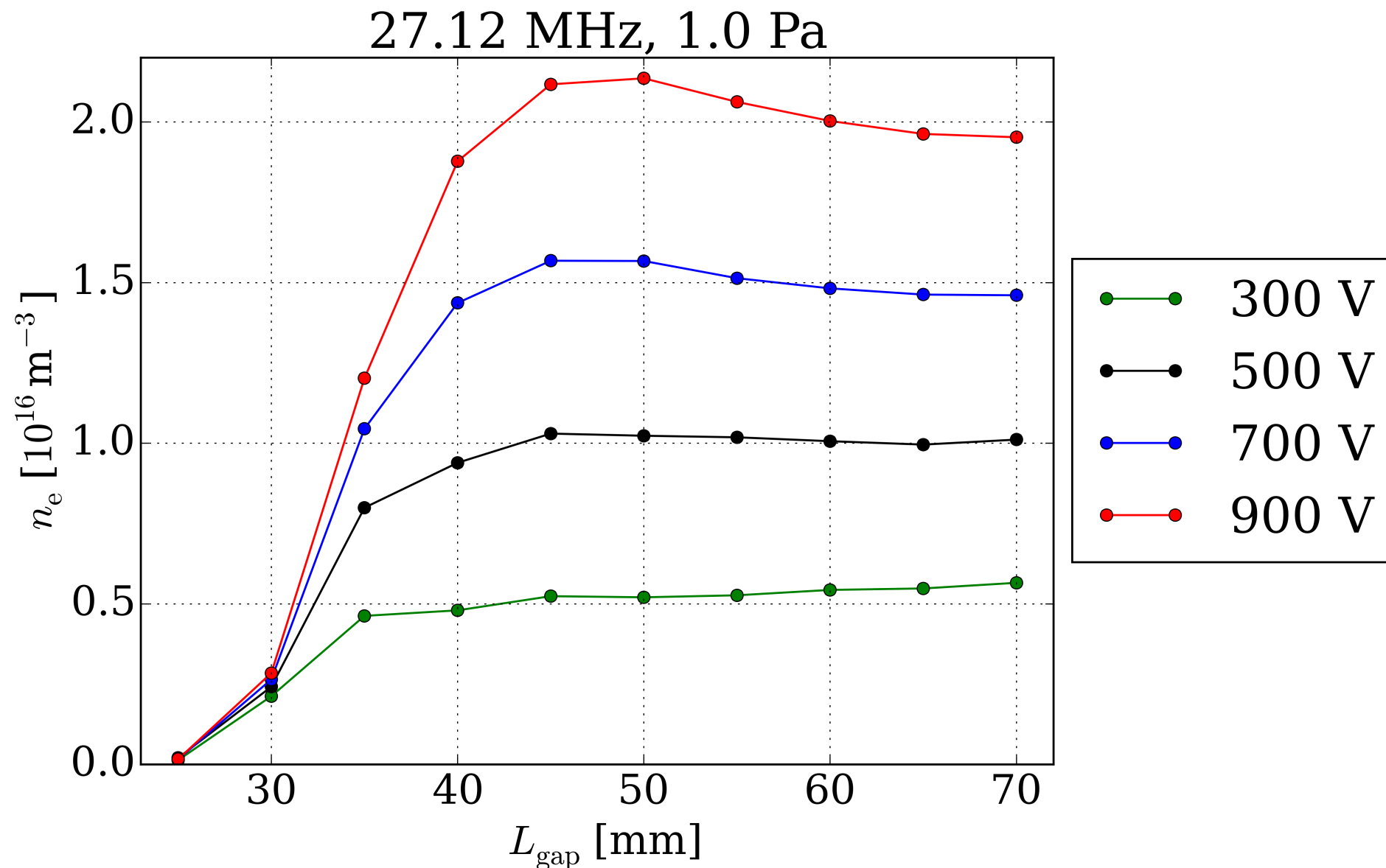


Gap variation: Distribution function

EFEDF (driven electrode)

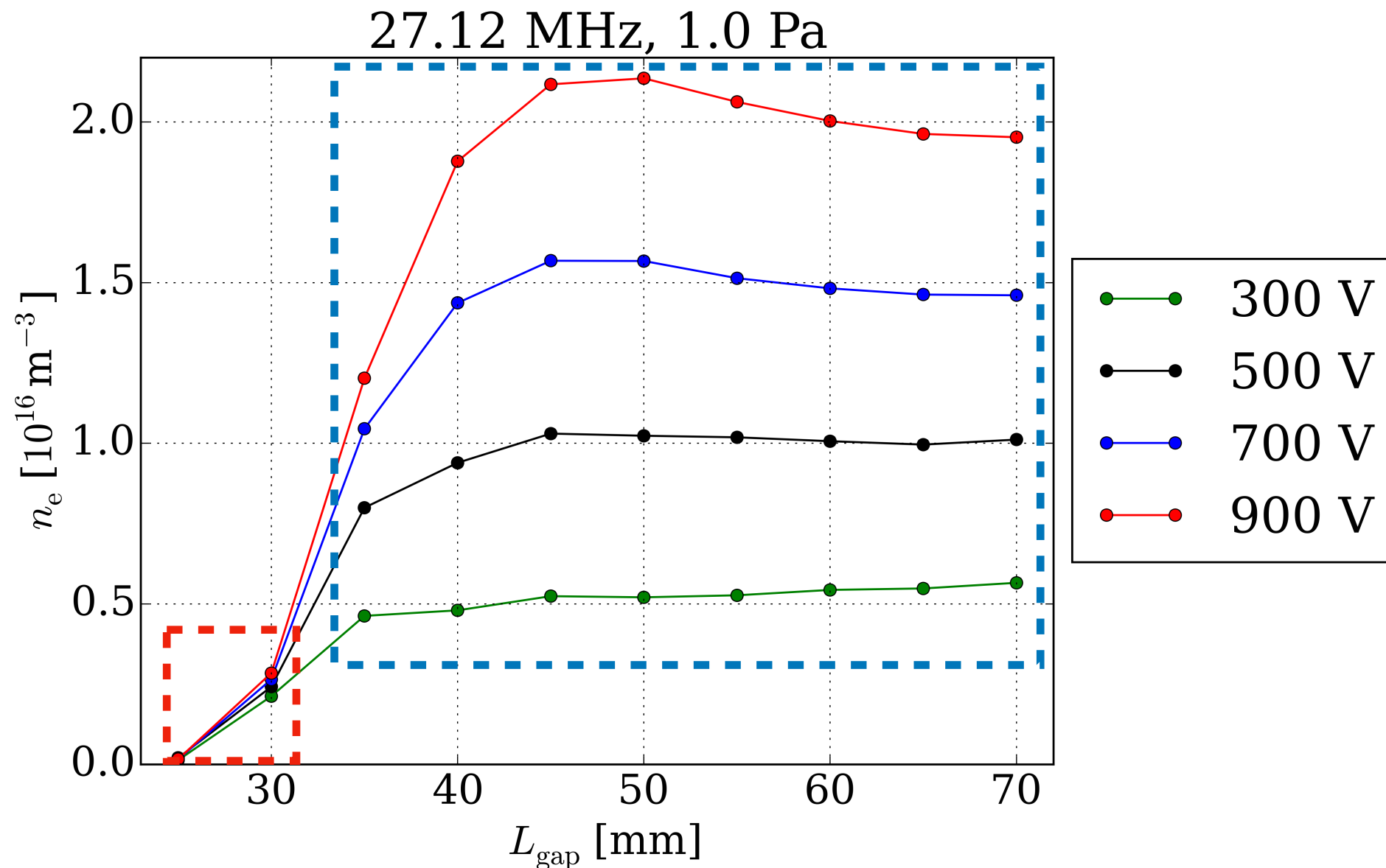


Gap variation: Operation regime



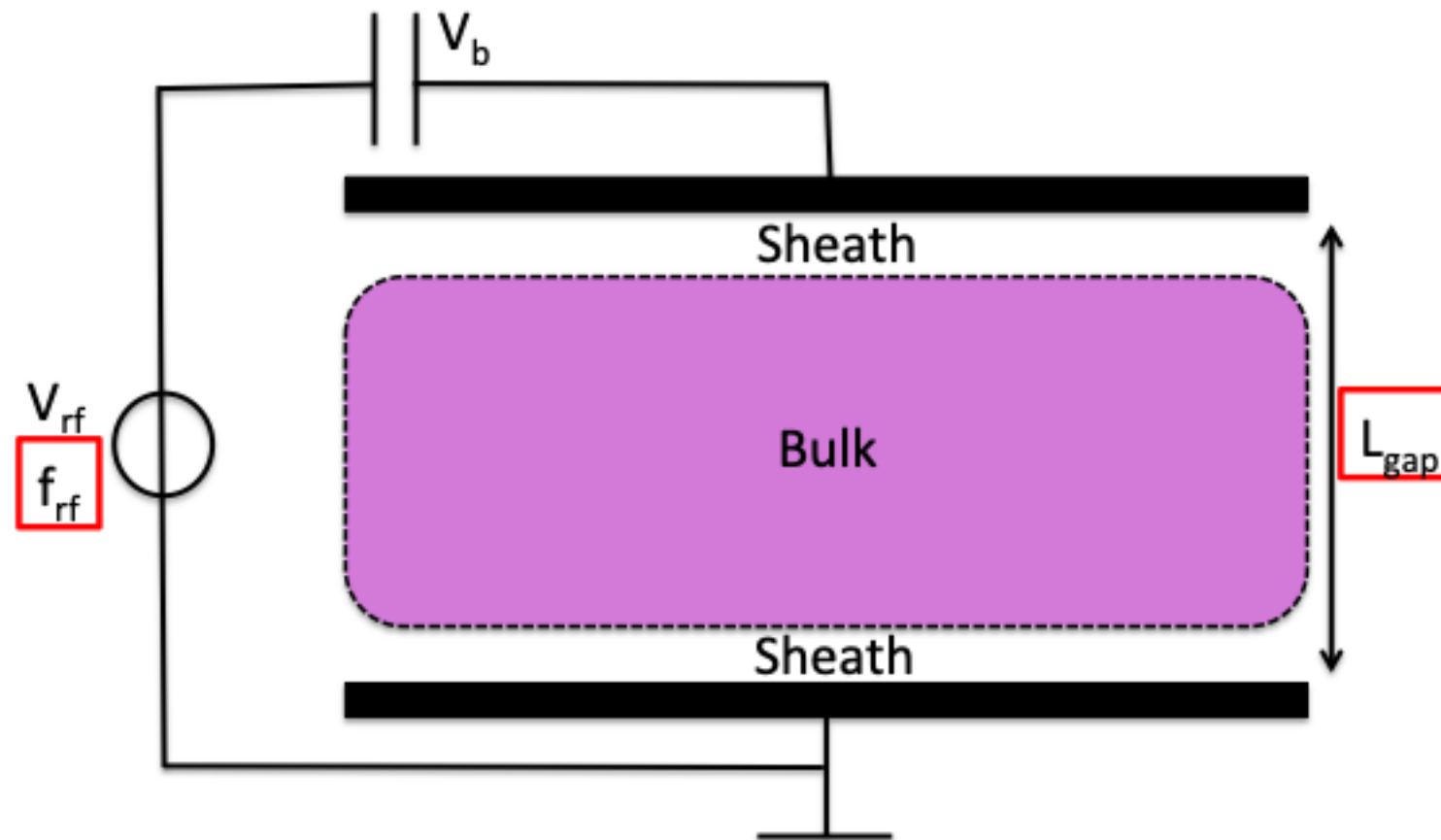
- gap size variation for different rf-voltages
- increasing the rf-voltage leads to higher densities (only at larger gap sizes)

Gap variation: Operation regime



- gap size variation for different rf-voltages
- increasing the rf-voltage leads to higher densities (only at larger gap sizes)
- smaller gap sizes: no chance to obtain a significant increase
- limitation of process control

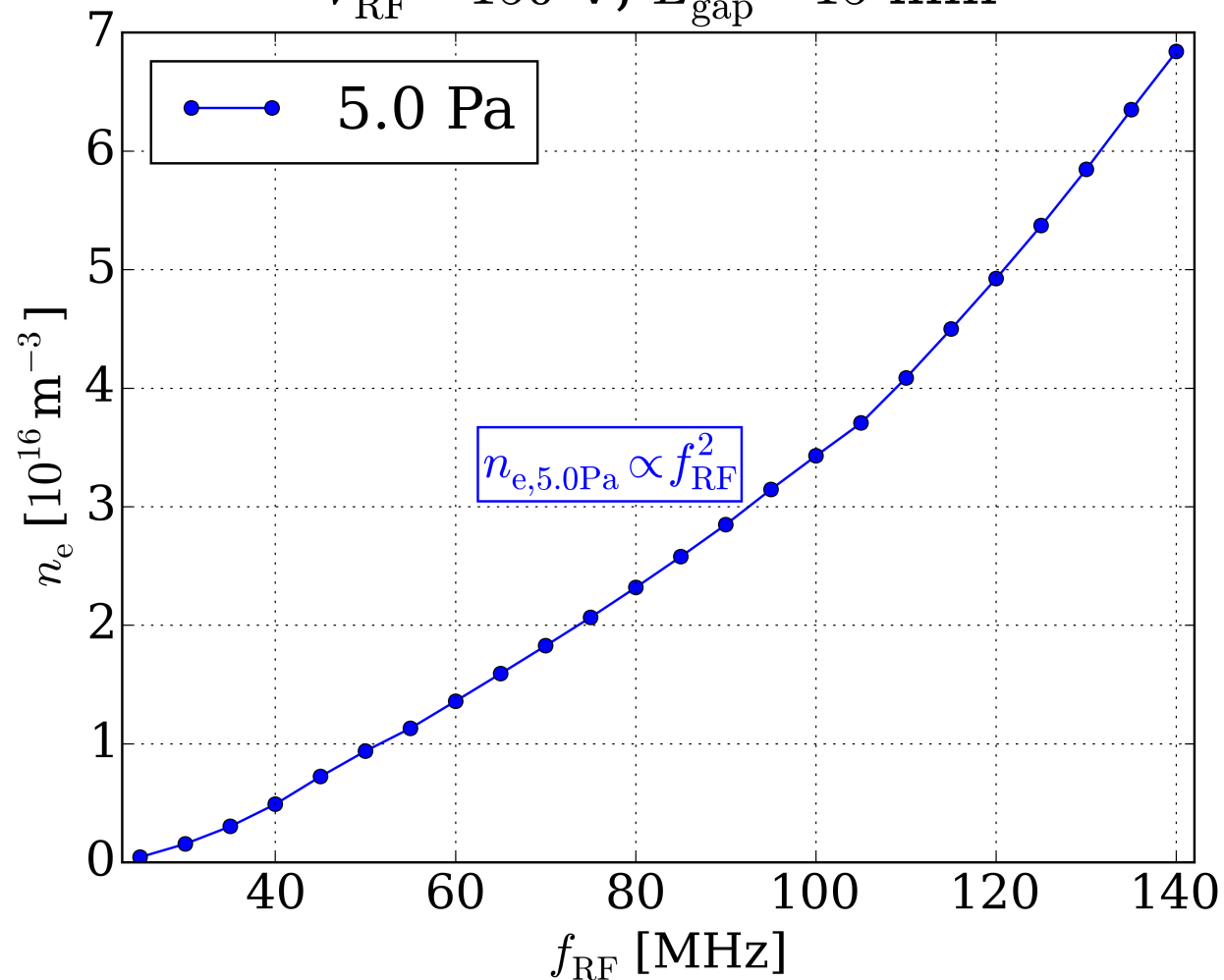
Driving frequency variation



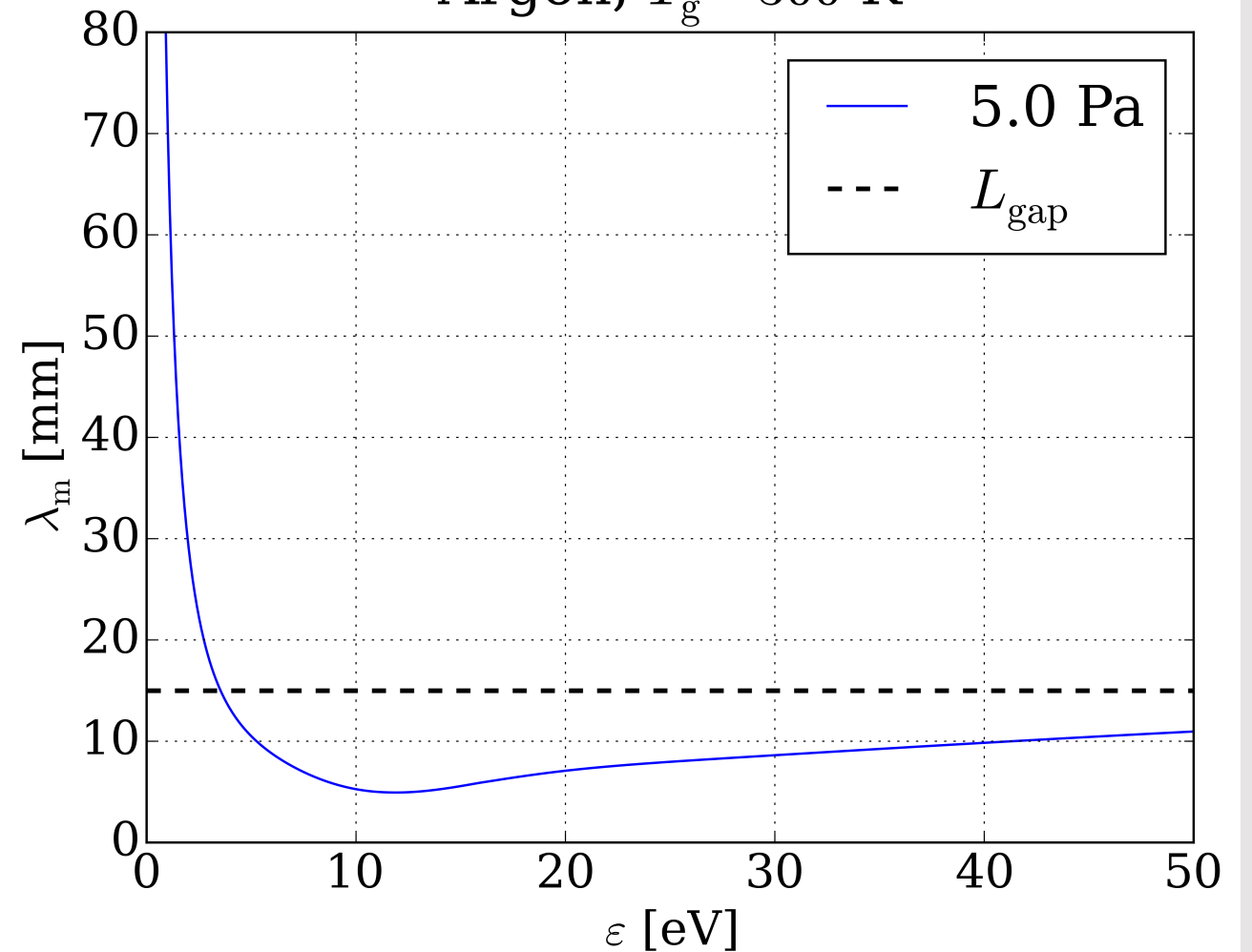
- 1d3v Particle-In-Cell simulation, Cartesian grid (perfectly symmetric), argon
- no secondary electron emission and no reflection at the boundary
- gap size: 15 mm
- gas pressure: 1 - 5 Pa
- driving voltage: 150 V
- **driving frequency: 25 - 100 MHz**

Driving frequency variation

$V_{\text{RF}} = 150 \text{ V}$, $L_{\text{gap}} = 15 \text{ mm}$



Argon, $T_g = 300 \text{ K}$

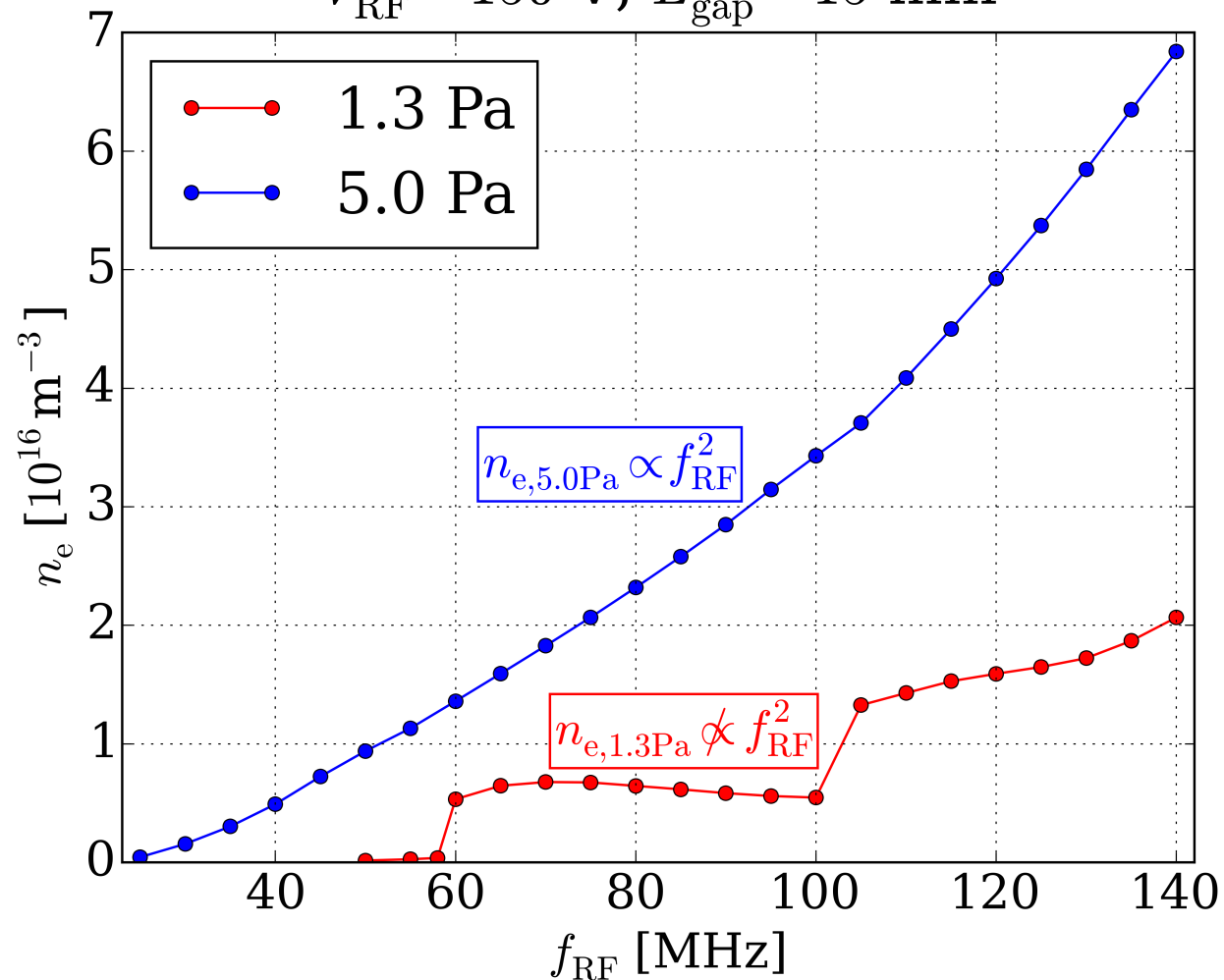


- expected quadratic dependence of the driving frequency¹
- 5 Pa \rightarrow collisional regime for energetic electrons

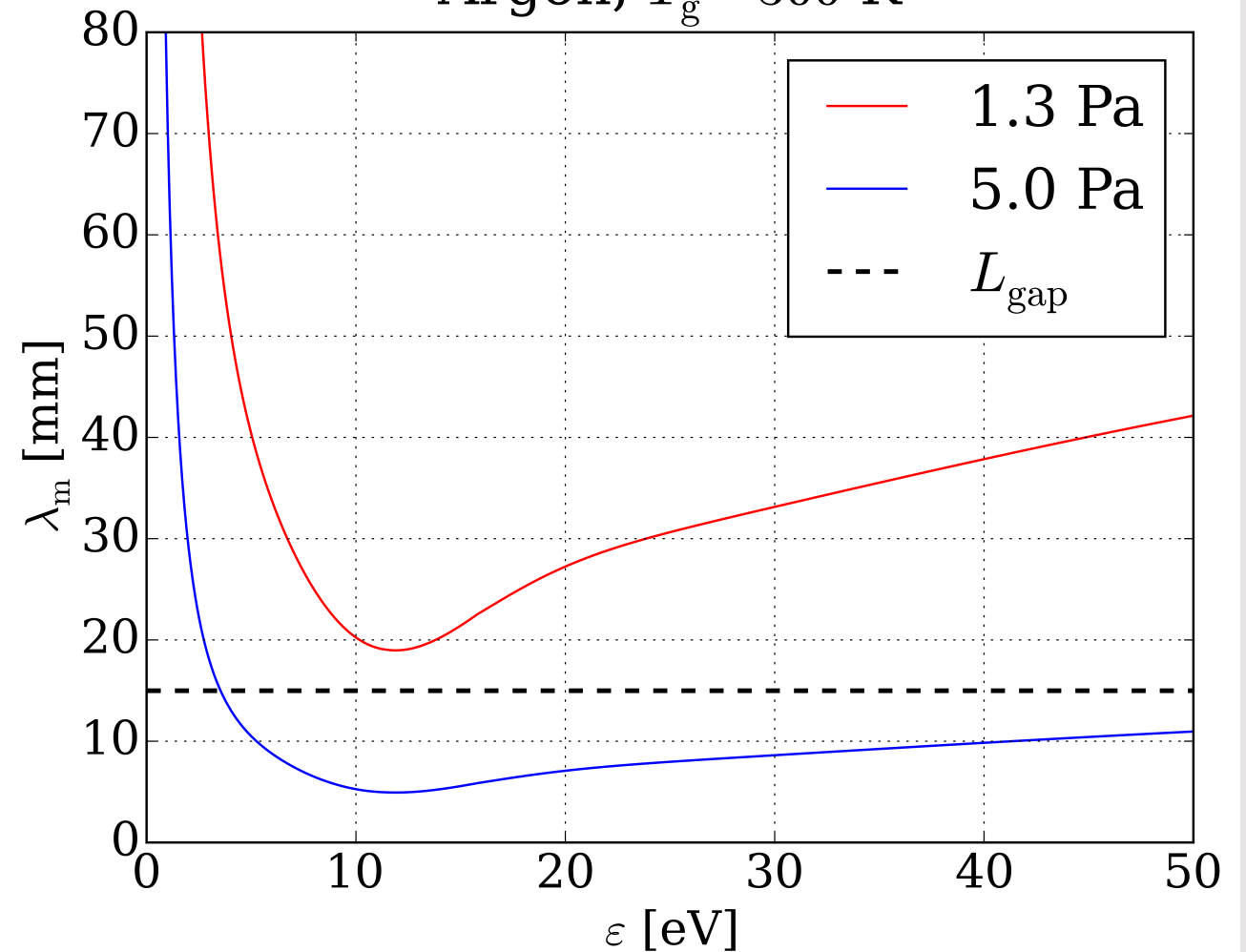
[1] M. A. Liebermann and A. J. Lichtenberg, *Principles of Plasma Discharges and Materials Processing* (2005)

Driving frequency variation

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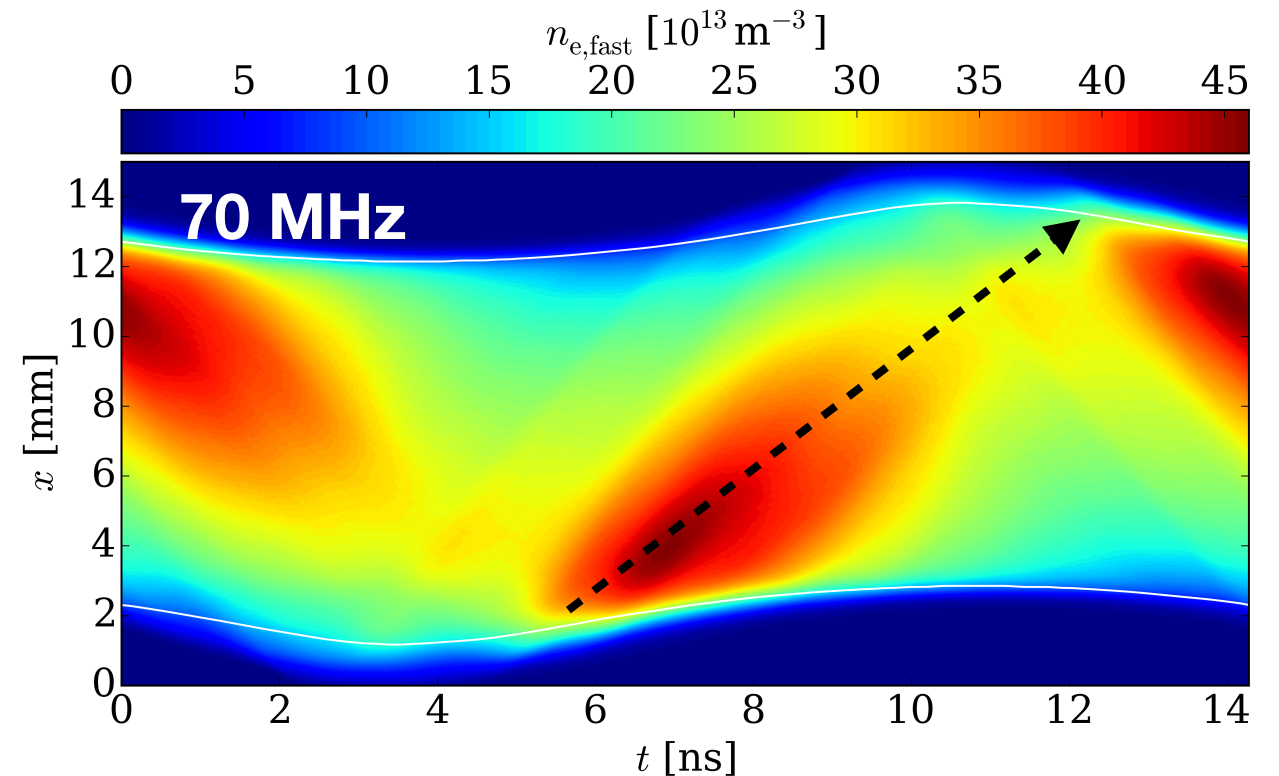
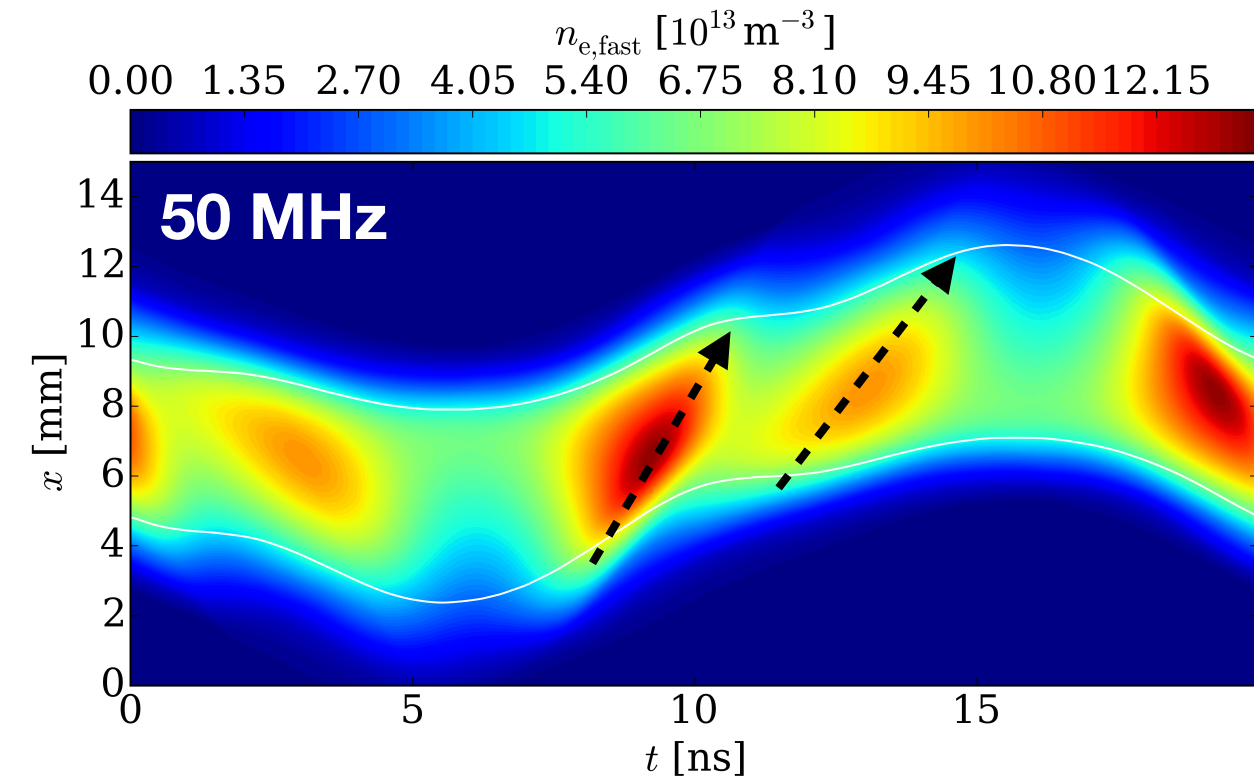
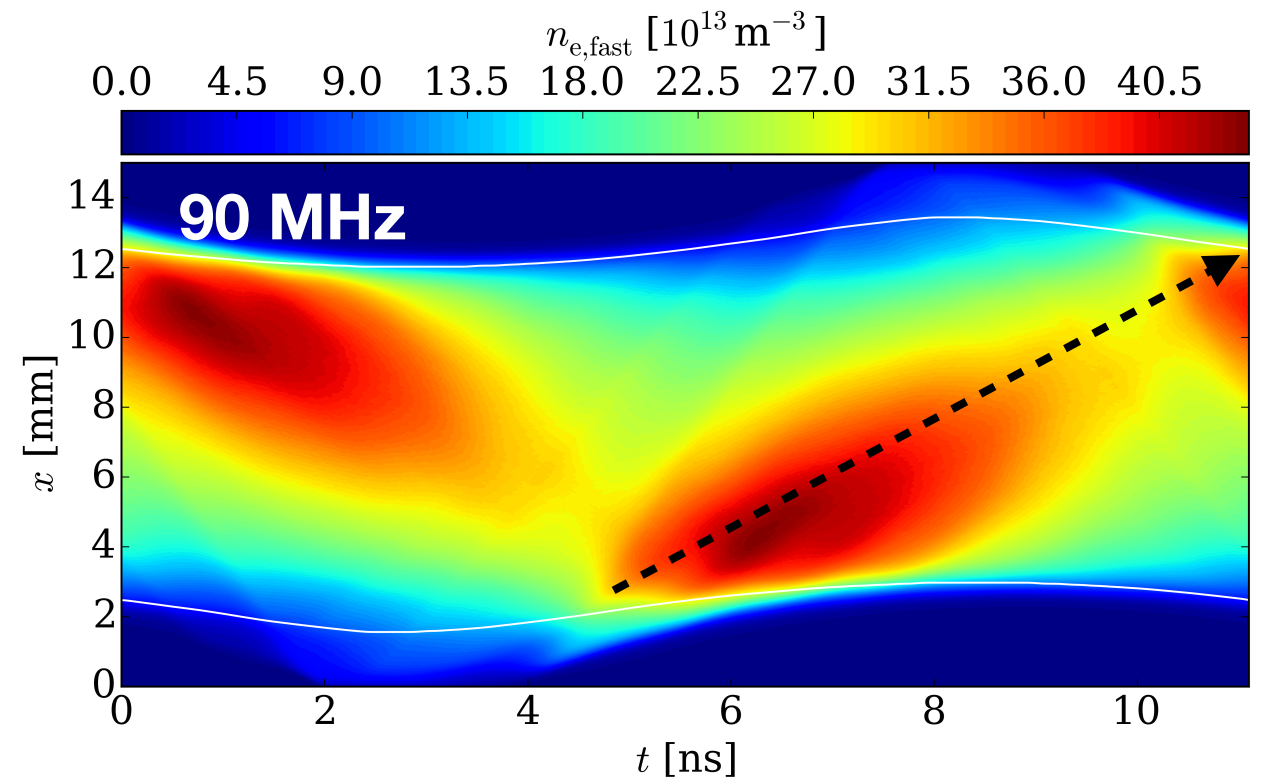
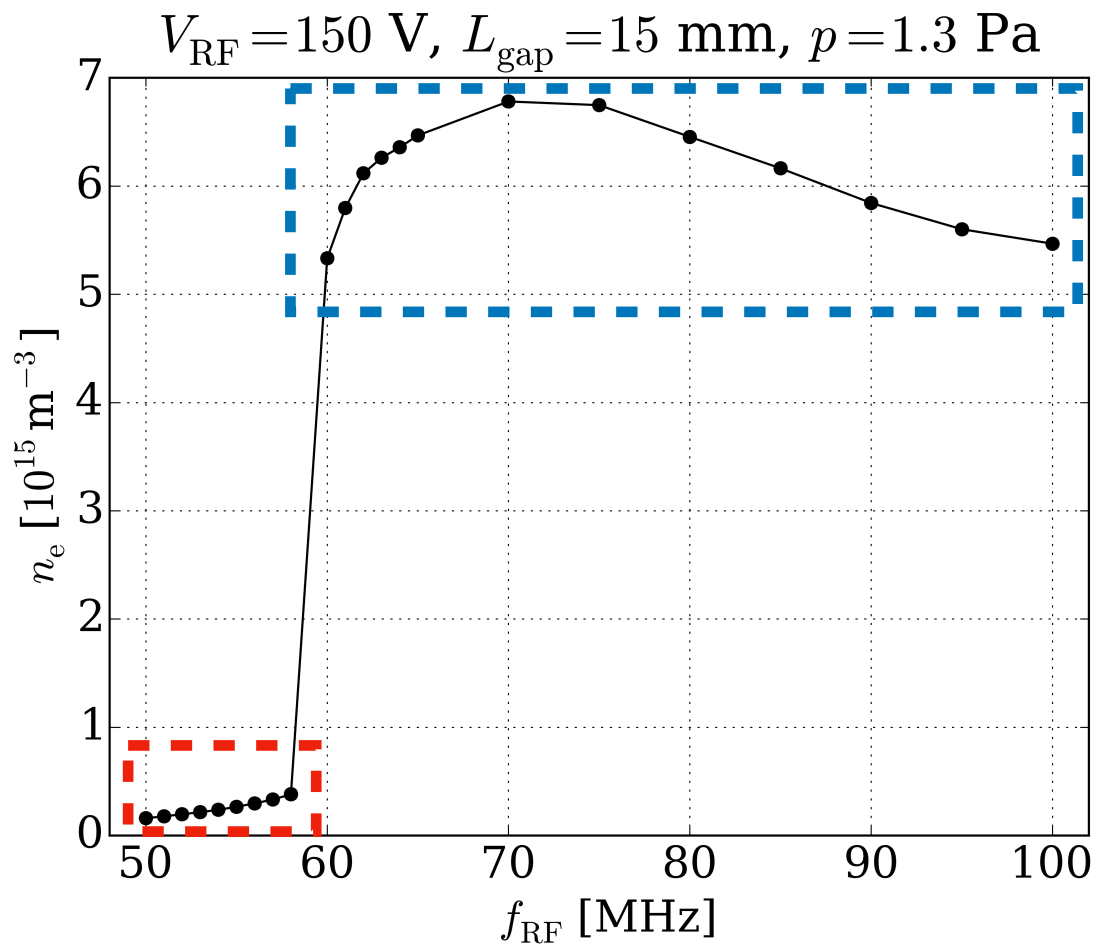


- expected quadratic dependence of the driving frequency¹
- 5 Pa → collisional regime for energetic electrons
- decreasing the pressure, no quadratic dependence of the driving frequency²
- 1.3 Pa → energetic electrons traverse collisionlessly through the discharge

[1] M. A. Liebermann and A. J. Lichtenberg, *Principles of Plasma Discharges and Materials Processing* (2005)

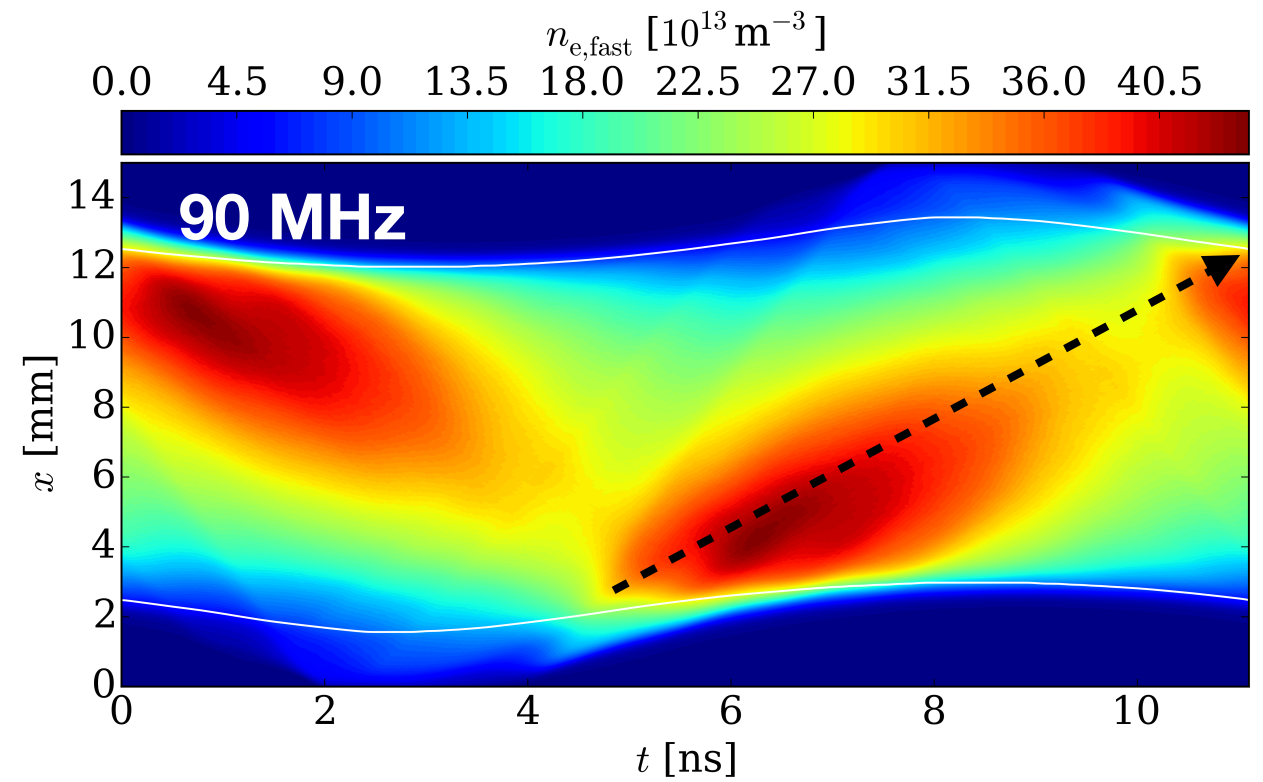
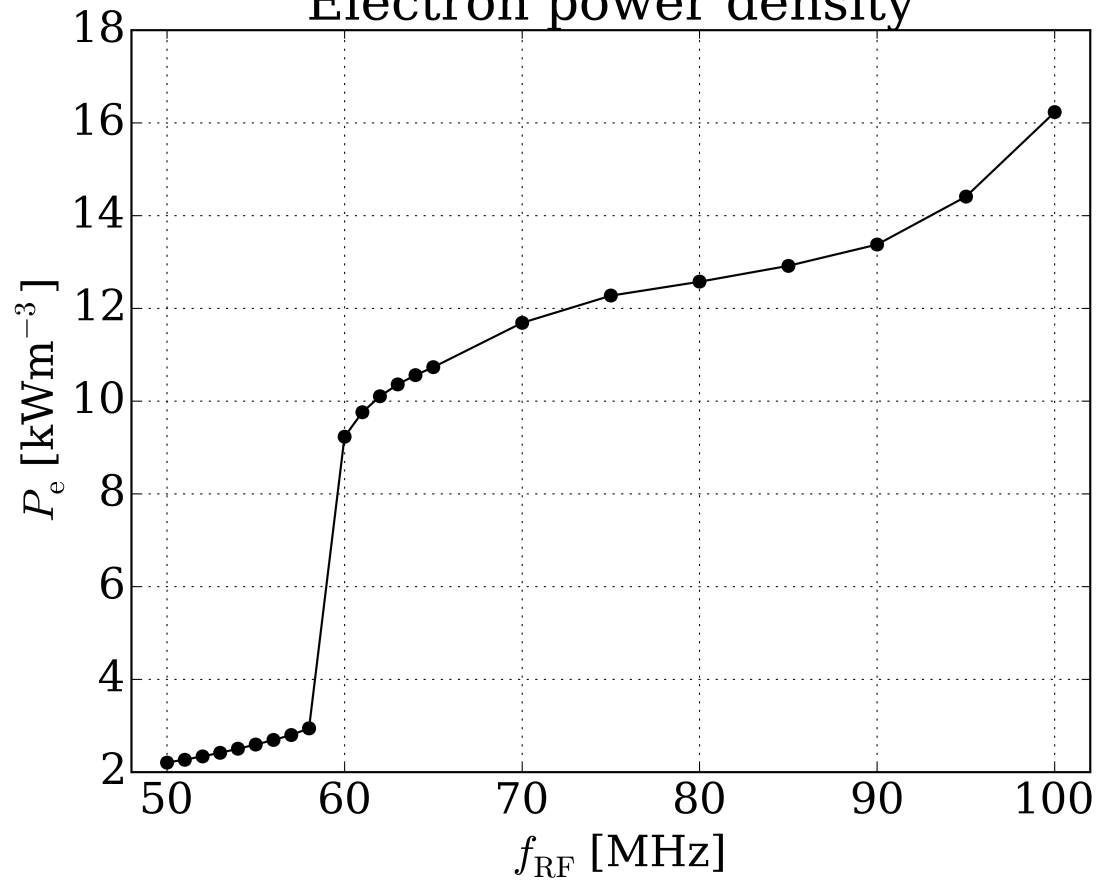
[2] S. Wilczek et al., *Plasma Sources Sci. Technol.*, 24, 024002 (2015)

RF variation: Fast electrons ($E > 15.7$ eV)

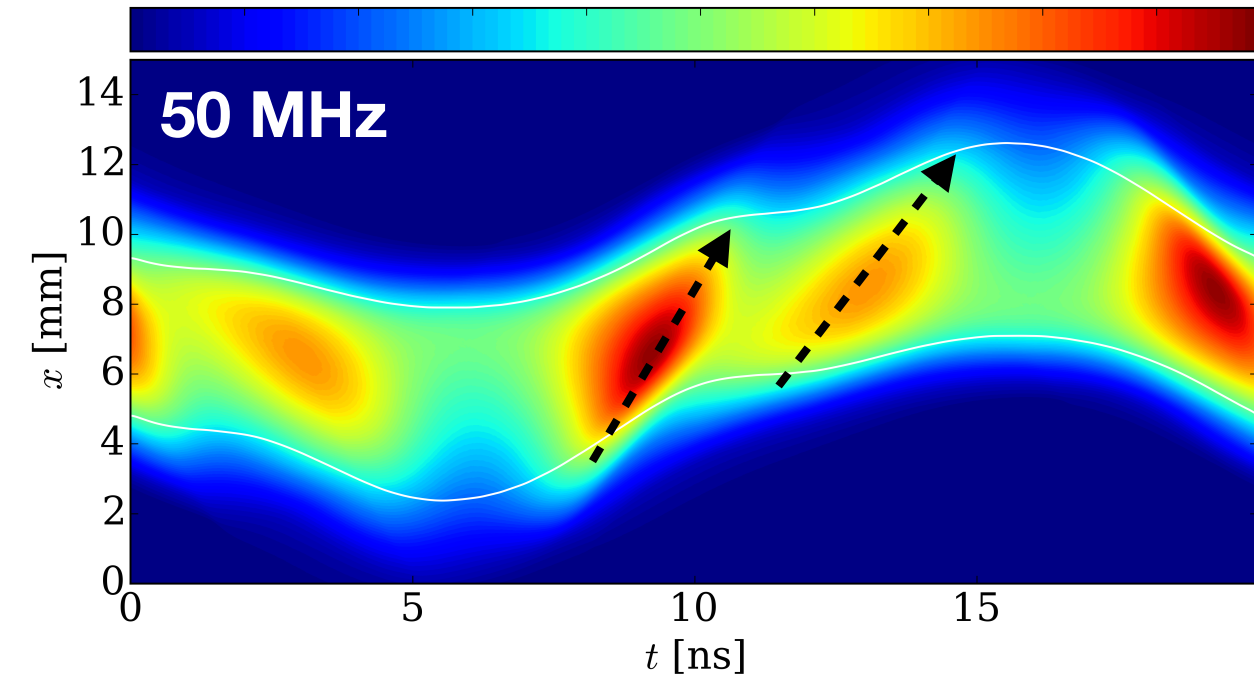


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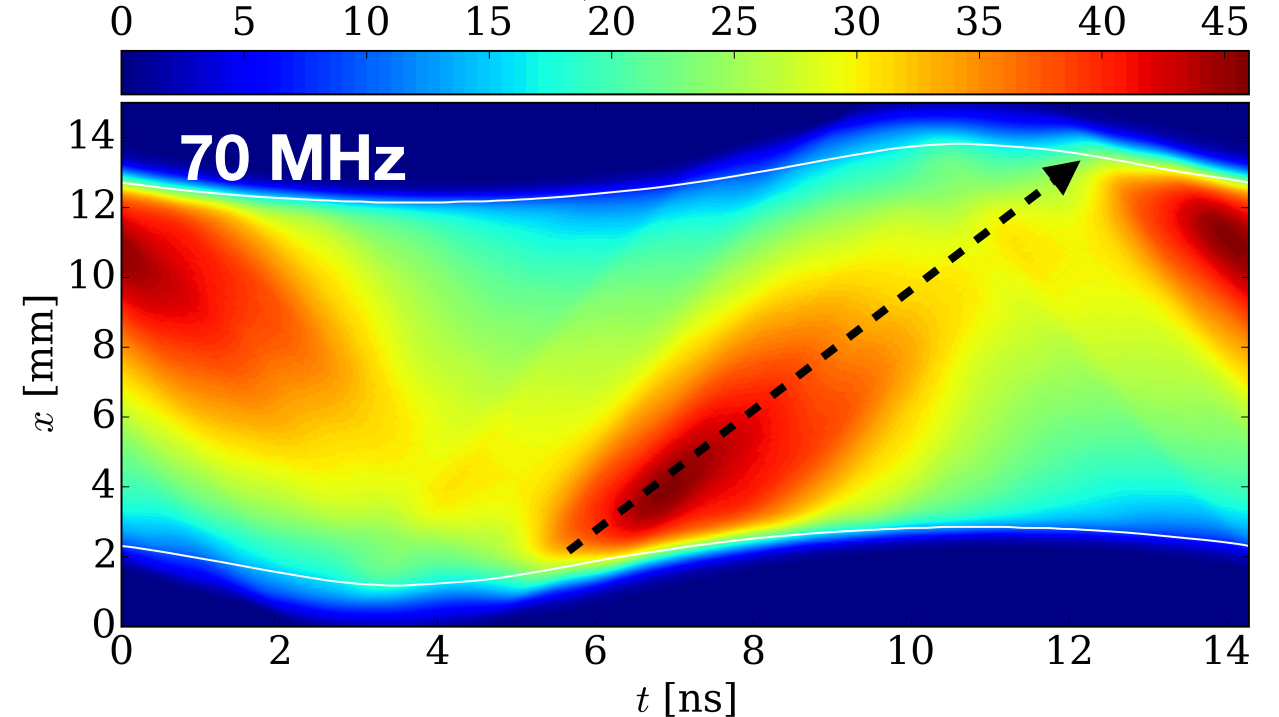
Electron power density



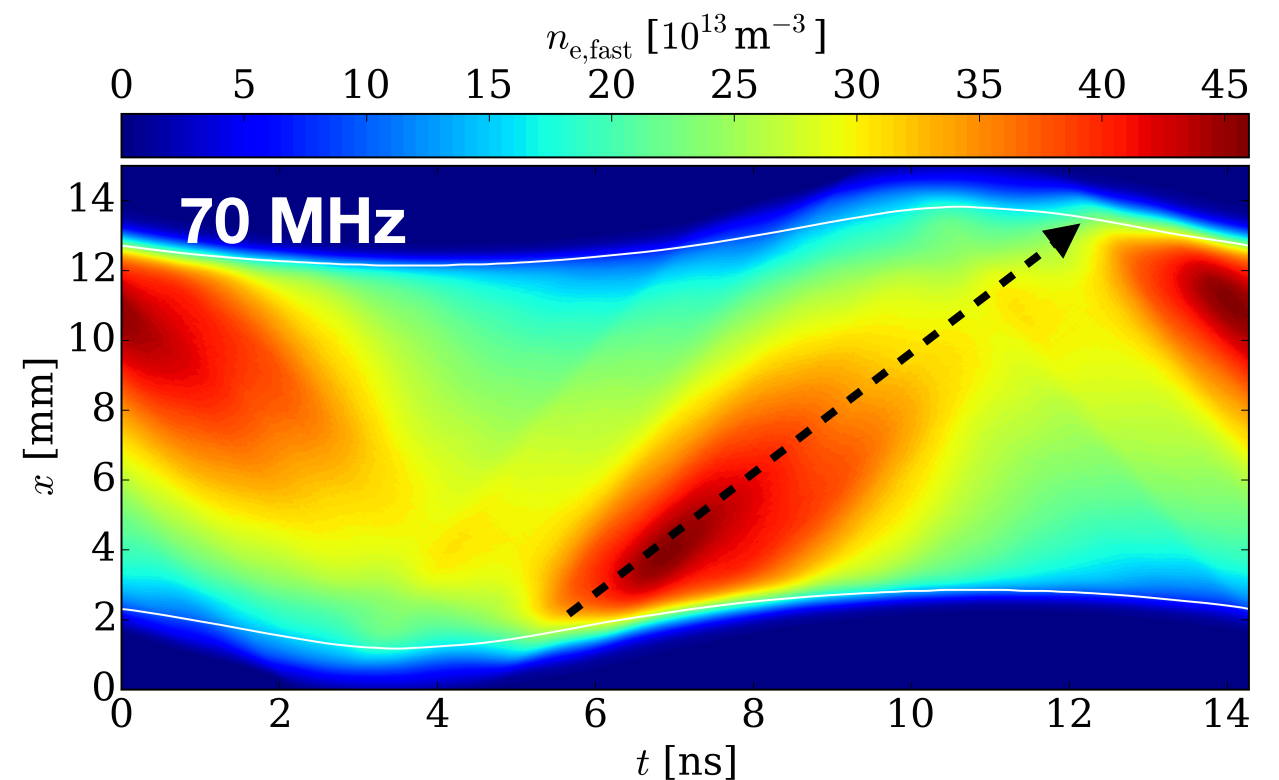
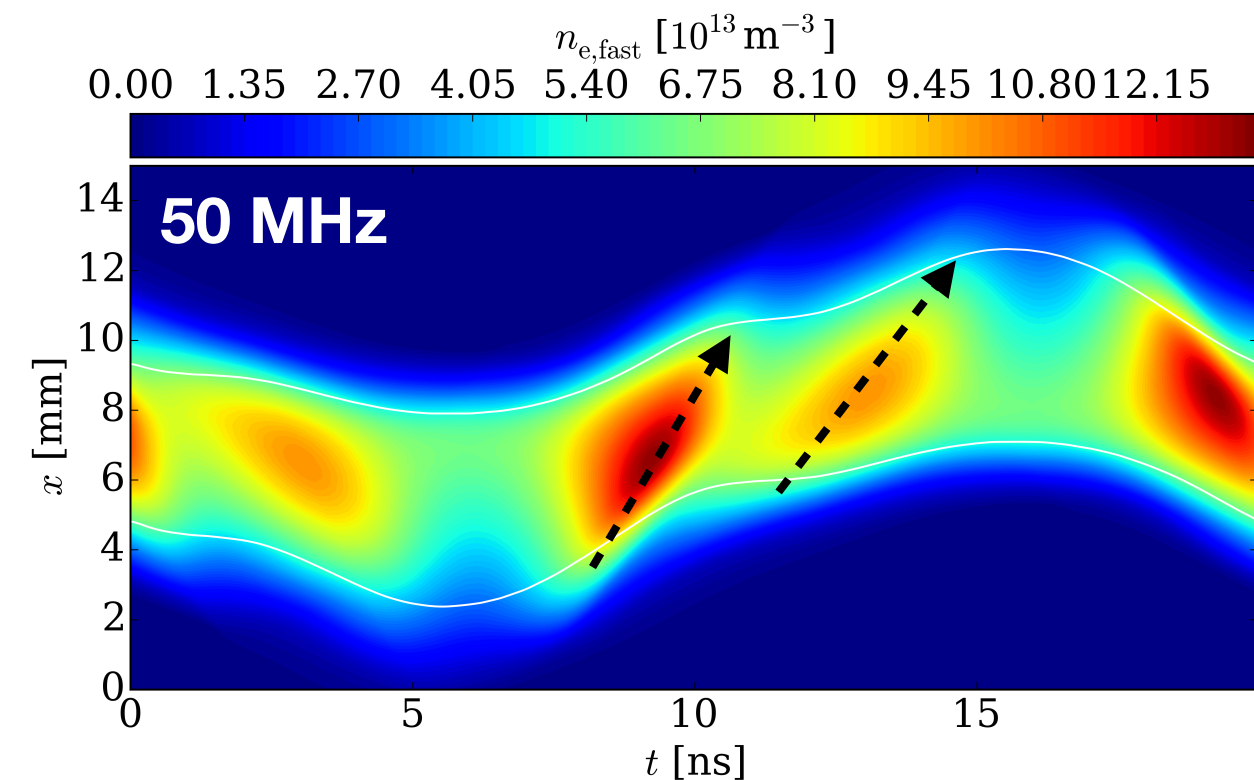
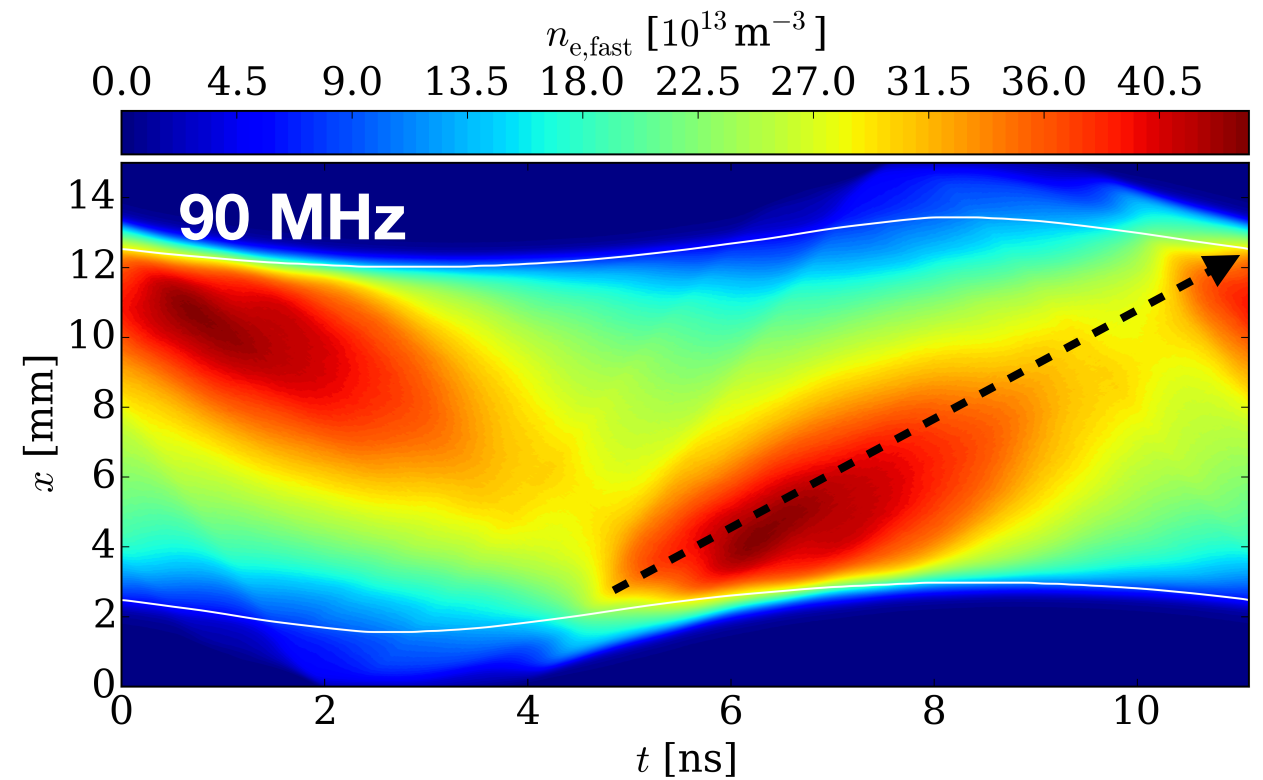
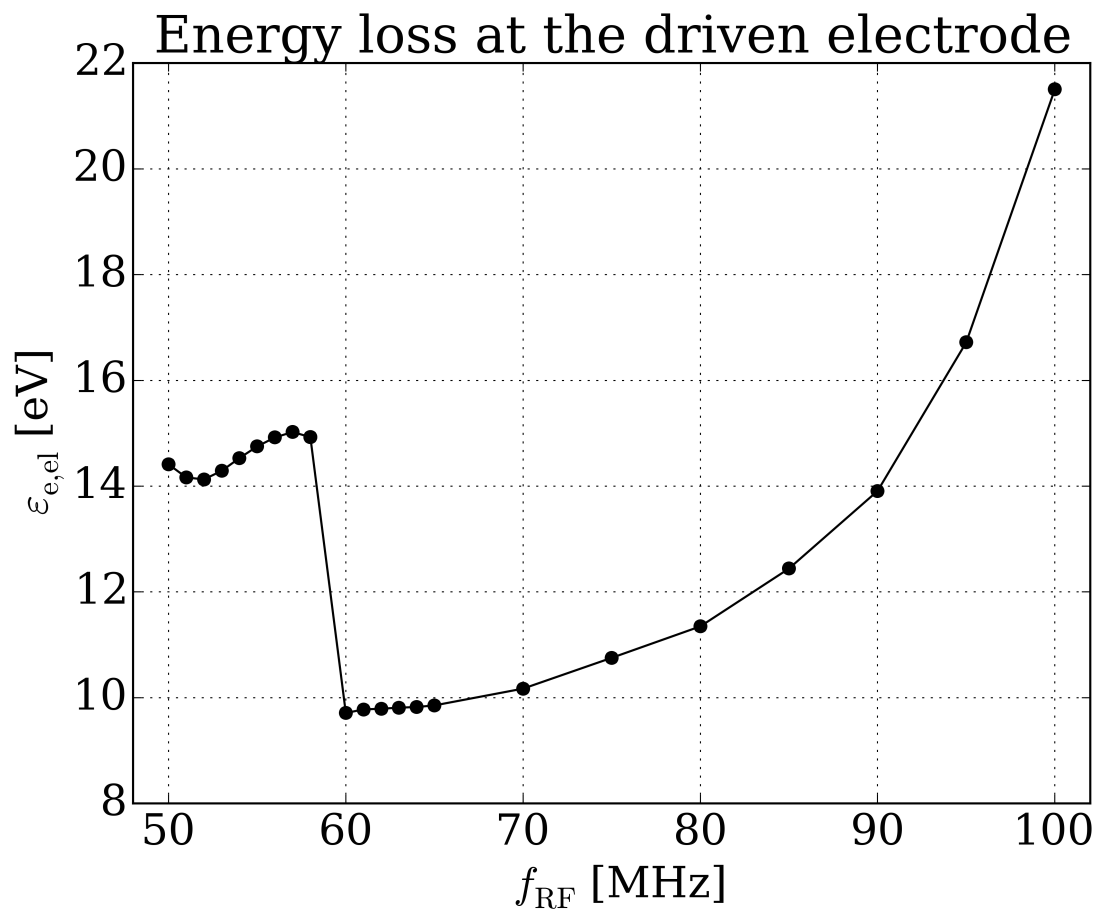
$n_{e,fast}$ [10¹³ m⁻³]



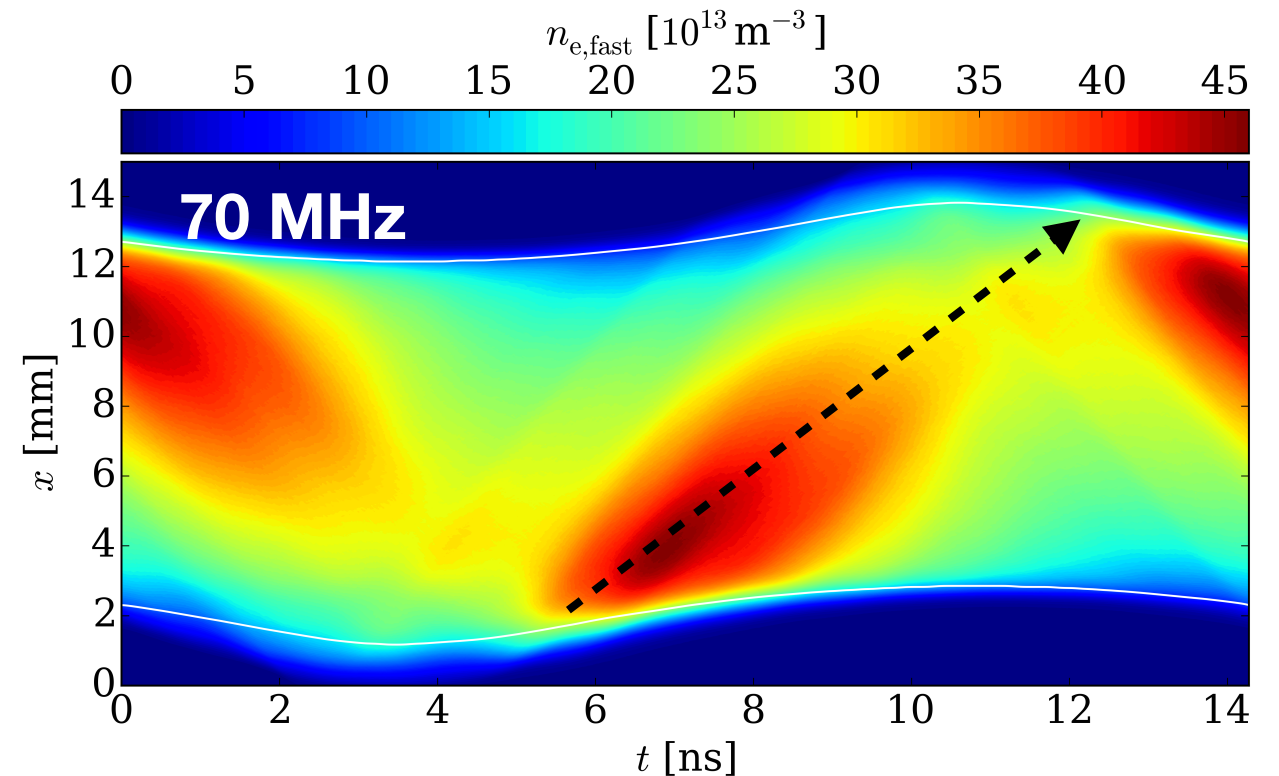
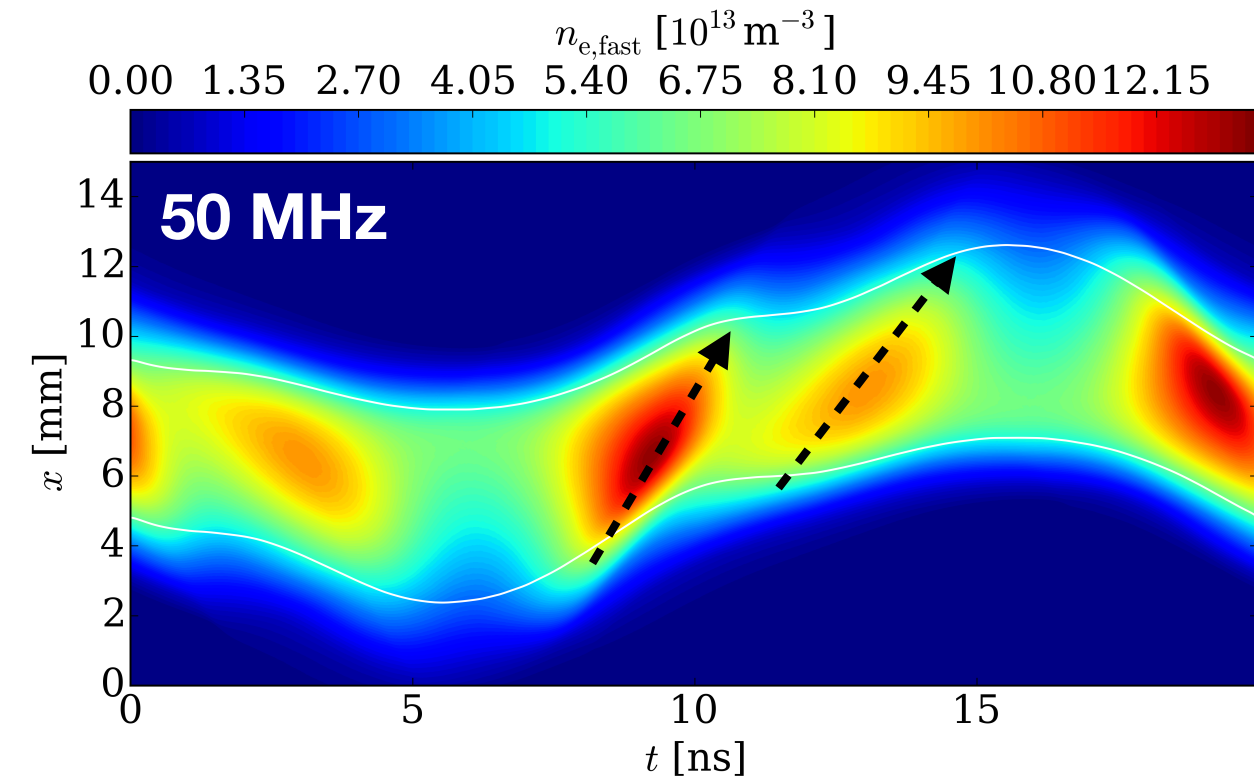
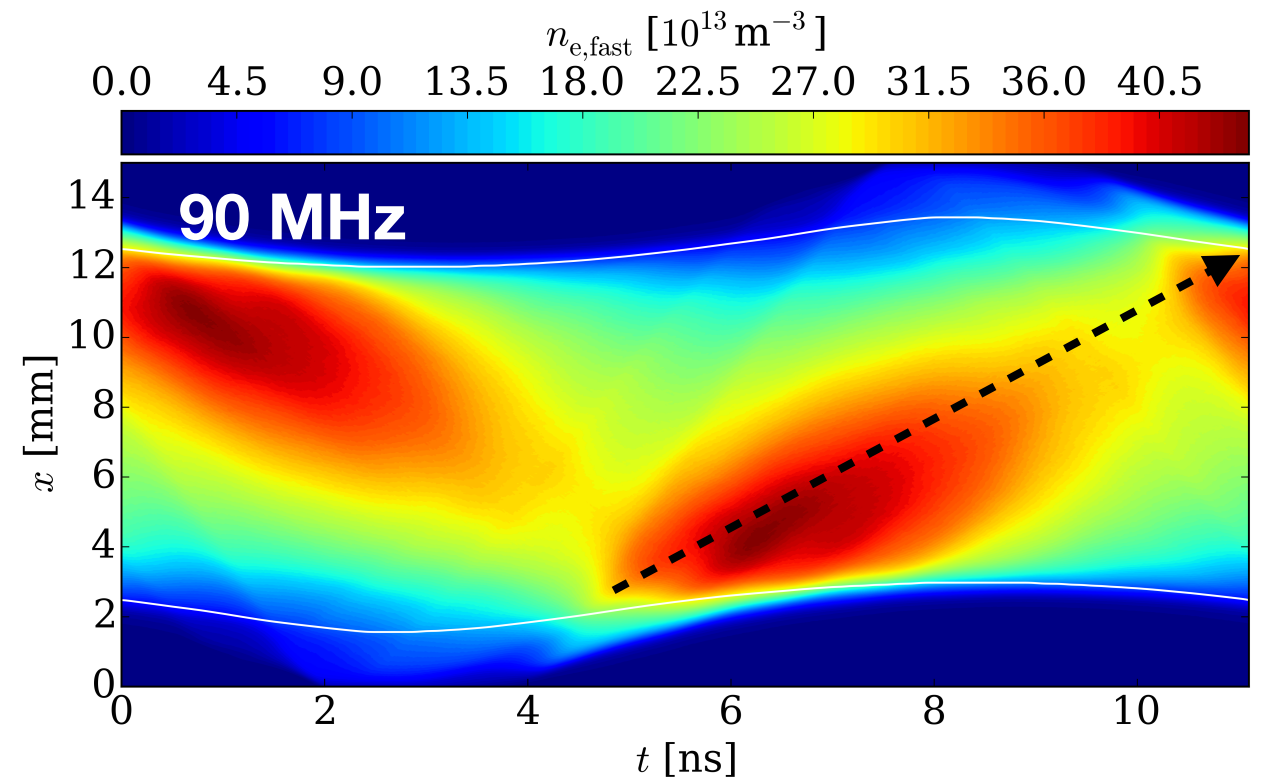
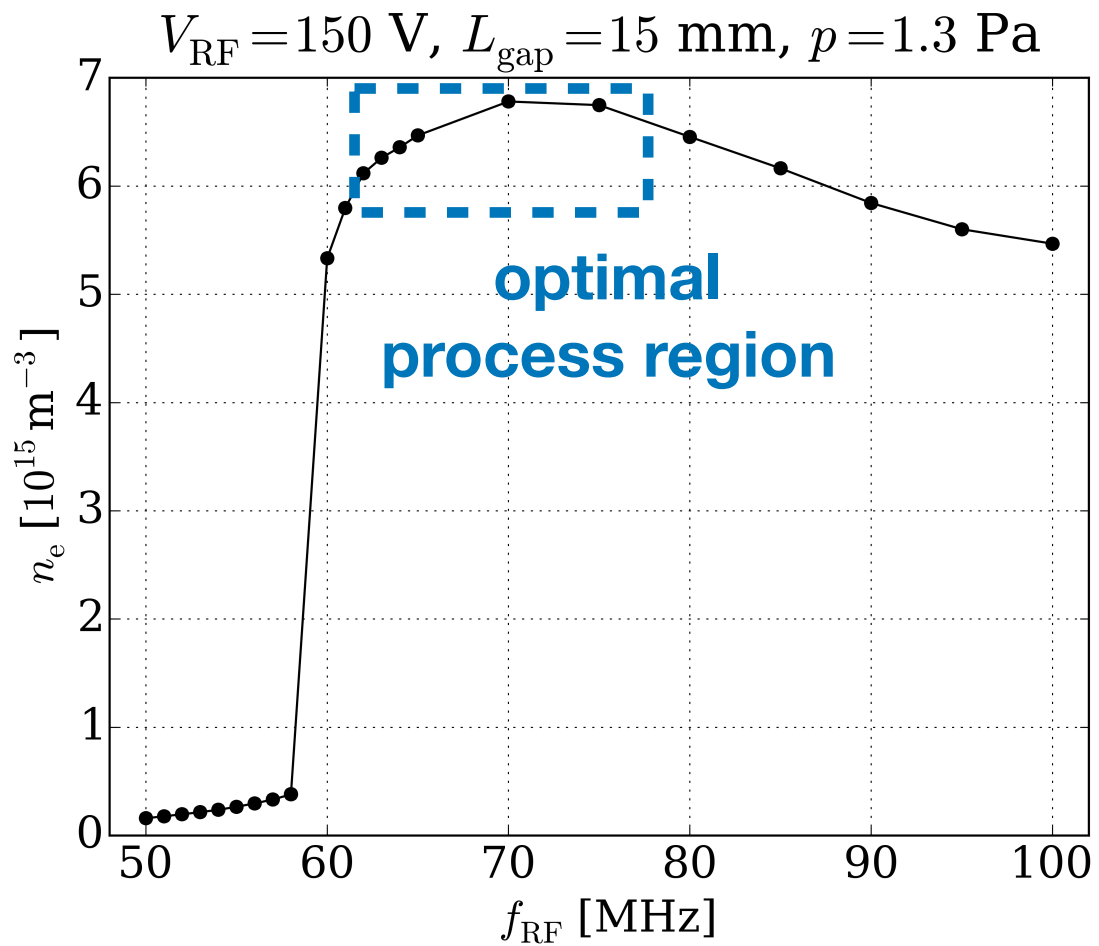
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RF variation: Fast electrons ($E > 15.7$ eV)



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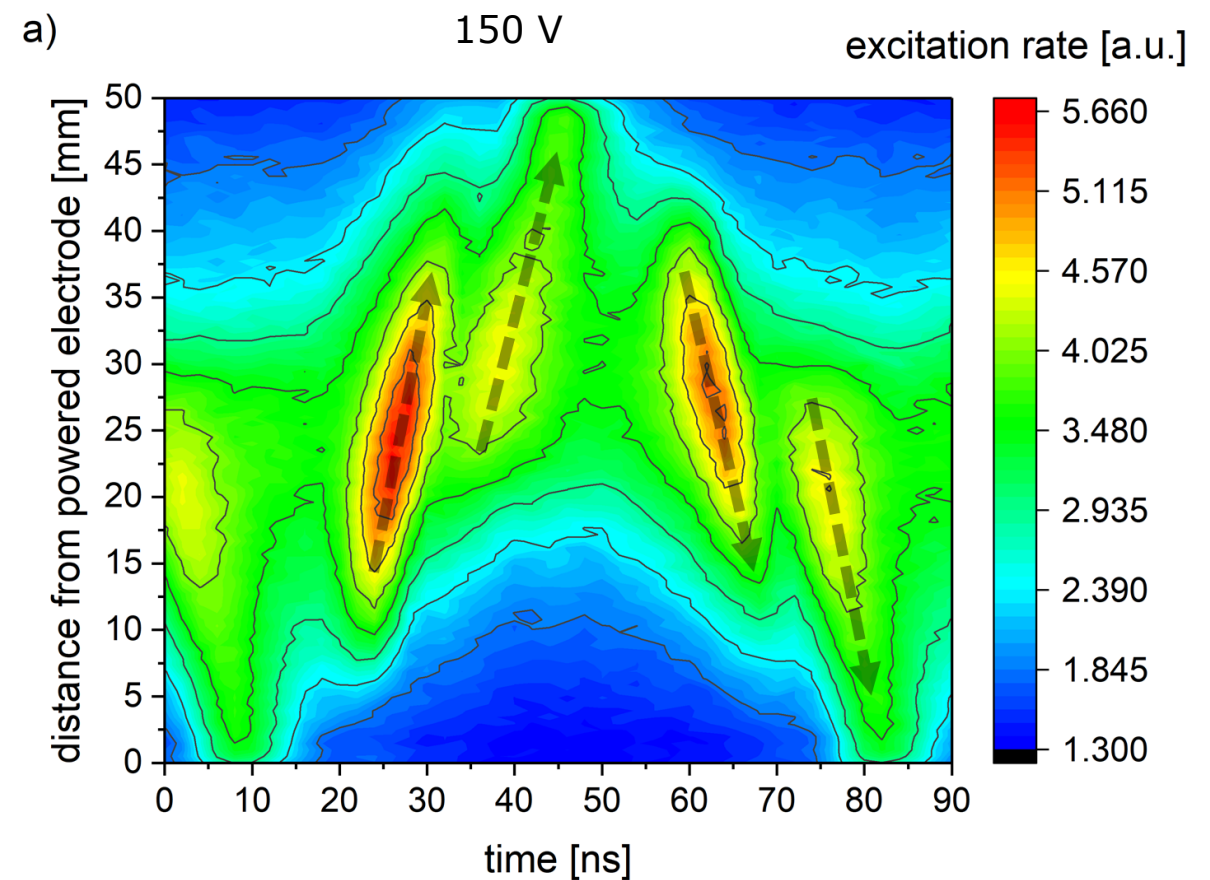
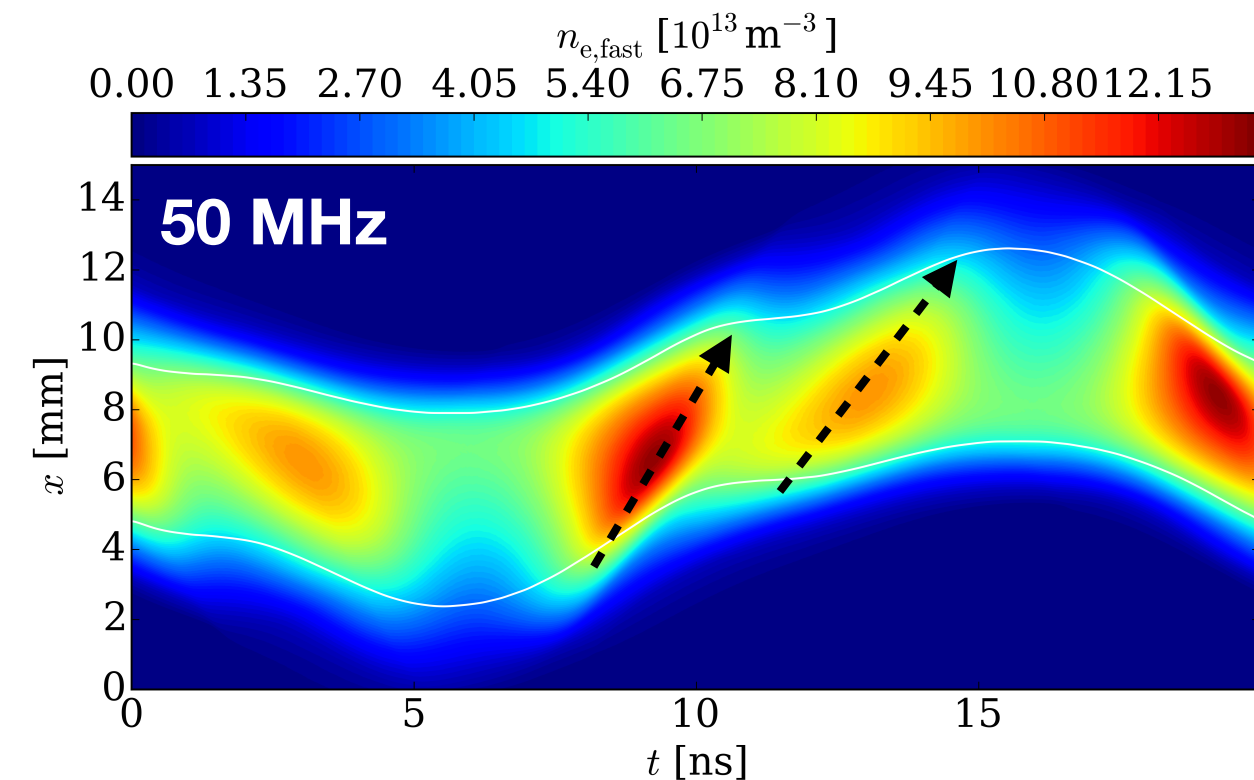


Next Step:

Why do we see sometimes multiple electron beams?

PIC³

Experiment, PROES⁴

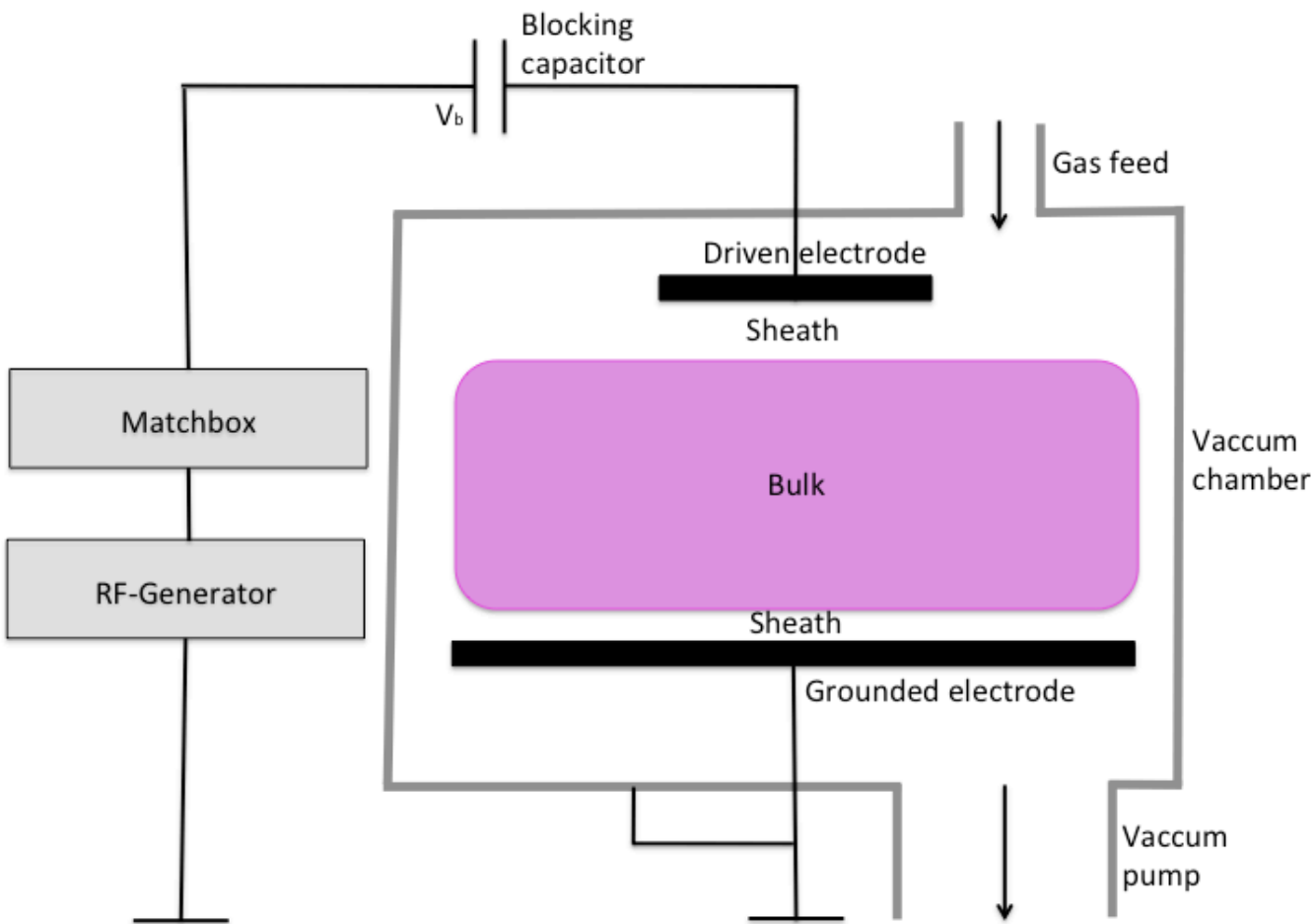


[3] S. Wilczek et al., Phys. Plasma 23, 063514 (2016)

[4] B. Berger et al., Plasma Sources Sci. Technol. 27, 12LT02 (2018)

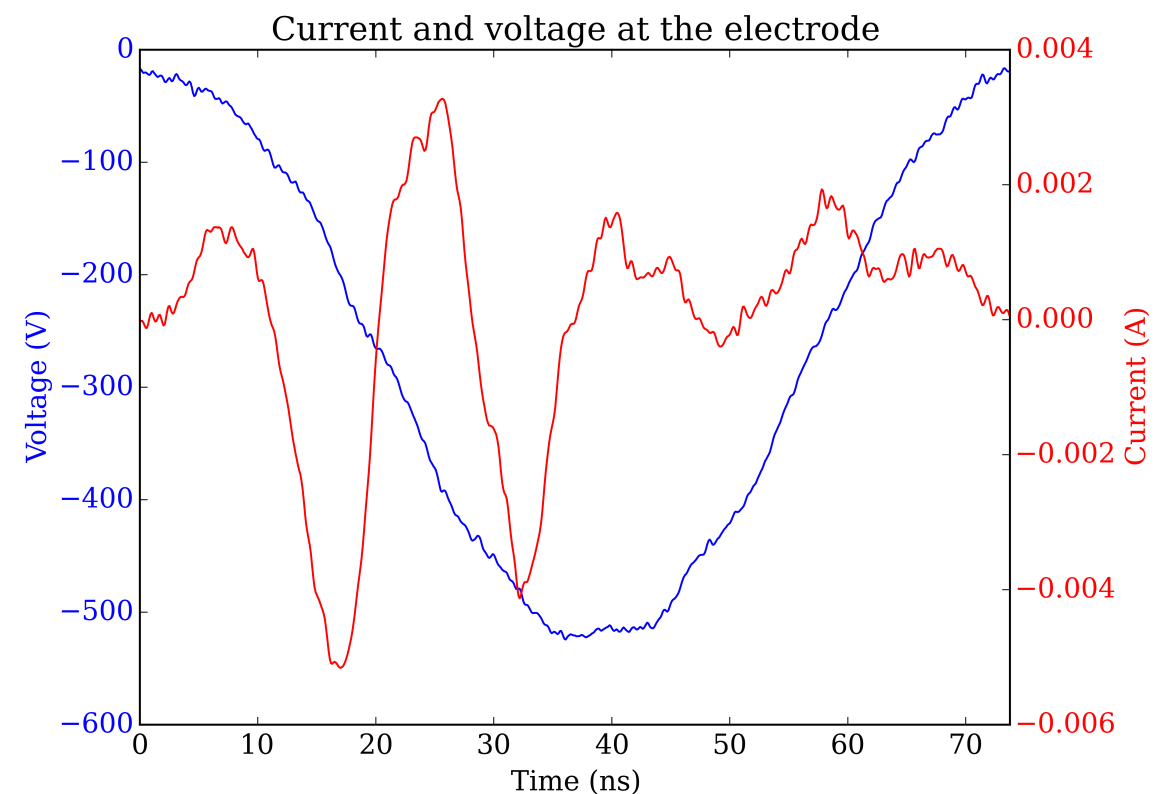
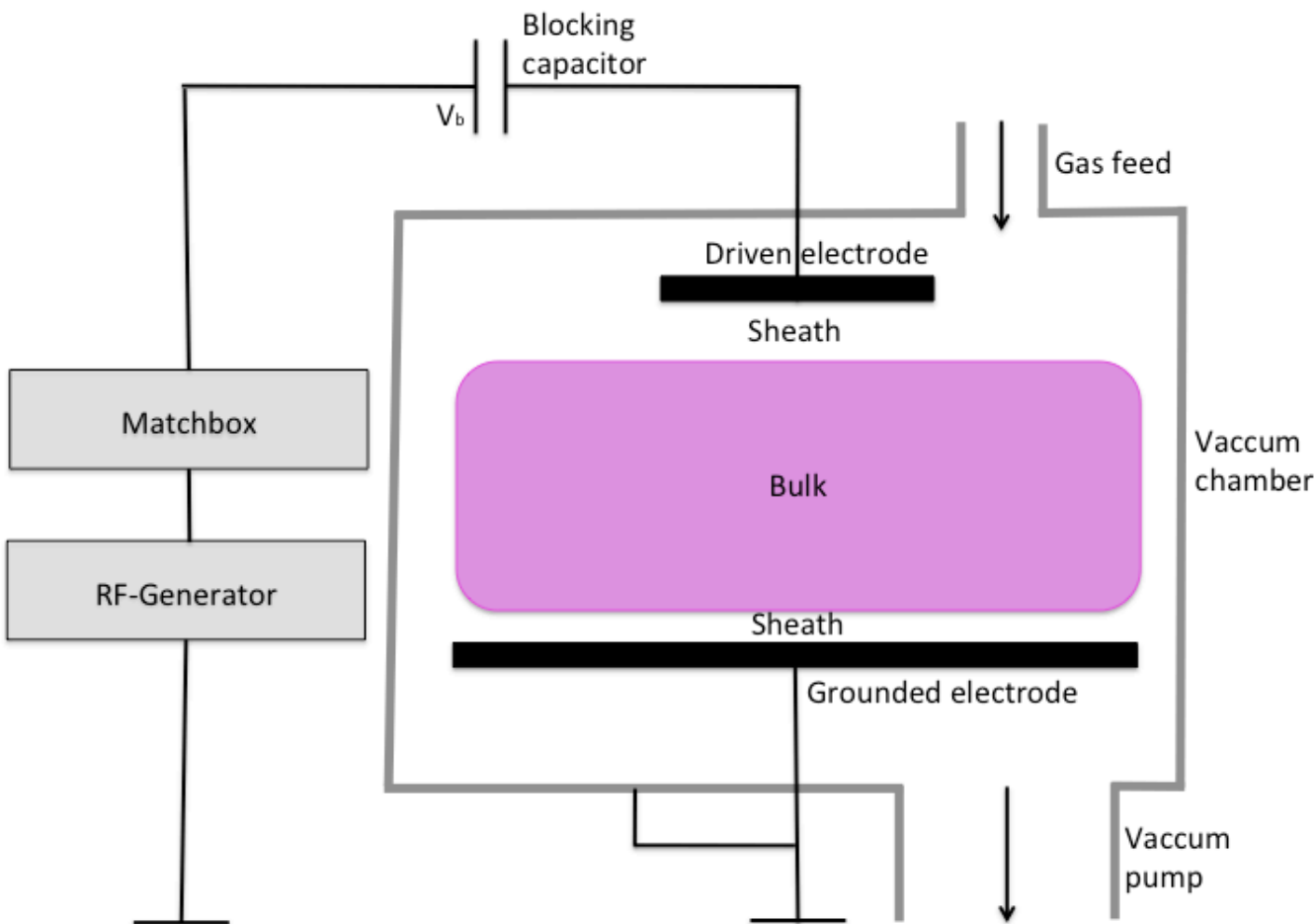
Nonlinearity

Motivation: Non-linearity



- CCRF discharges are mostly asymmetric, surfaces are naturally grounded
- measuring global parameters, like the voltage and the current at the electrode

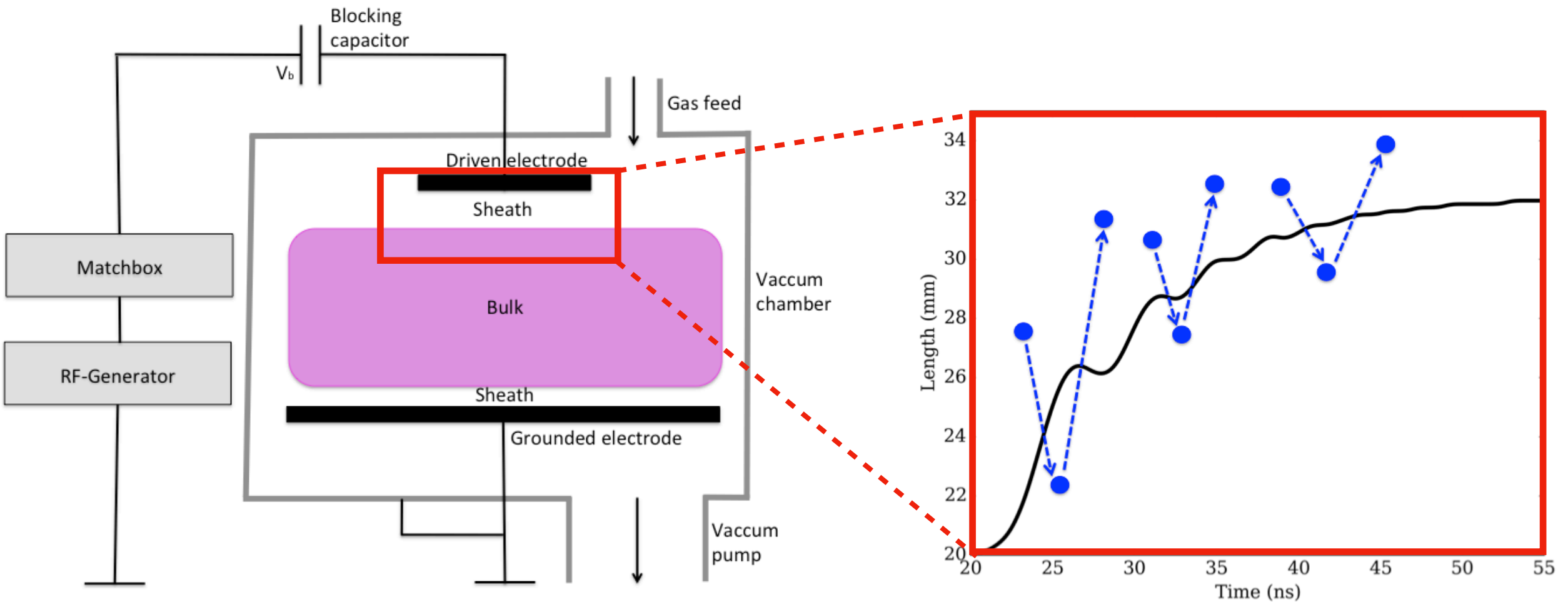
Motivation: Non-linearity



- CCRF discharges are mostly asymmetric, surfaces are naturally grounded
- measuring global parameters, like the voltage and the current at the electrode
- voltage is almost sinusoidal but current indicates harmonic oscillations⁵
- response of the nonlinear dynamics of the system

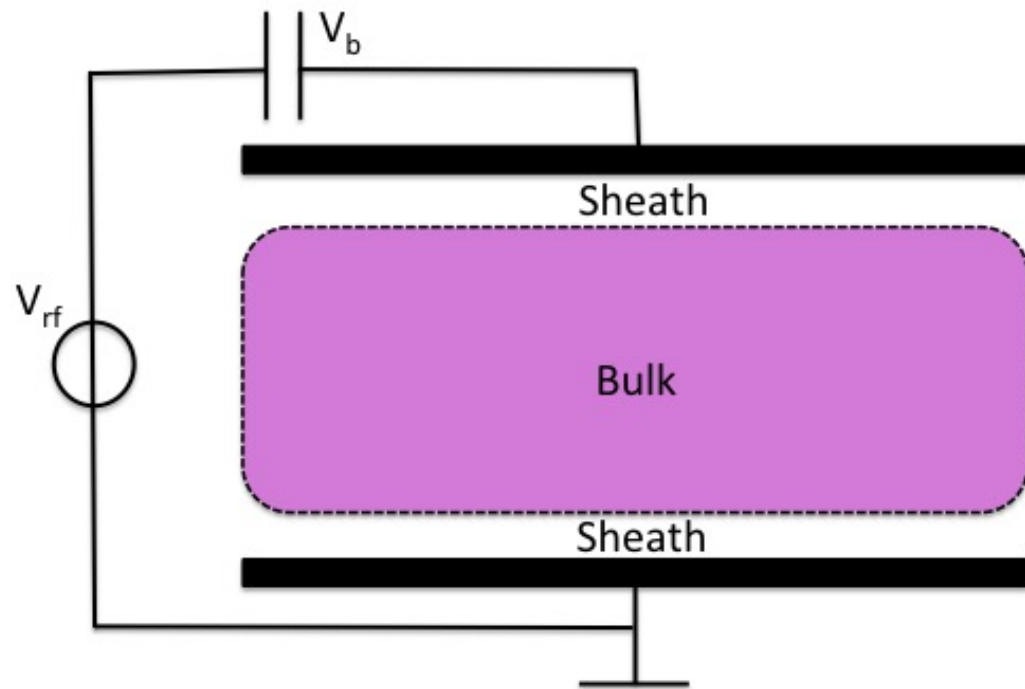
[5] J. Schulze et al., *J. Phys. D: Appl. Phys.*, **41**, 195212 (2008)

Goal of this work: Non-linearity

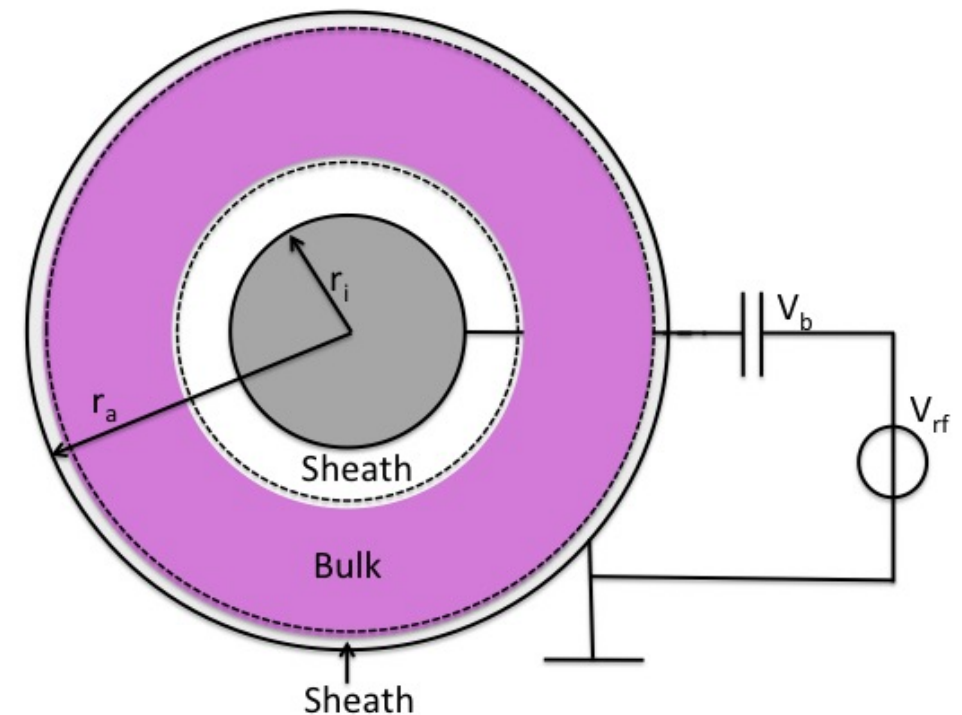


- focus on the nonlinear interaction between sheath and bulk on a ns timescale
- what is the physical origin of the generation of harmonics?
- how does the nonlinearity influence the electron power dynamics?

symmetric (Cartesian)



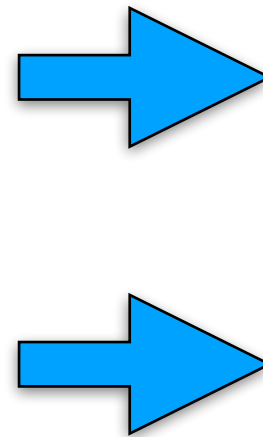
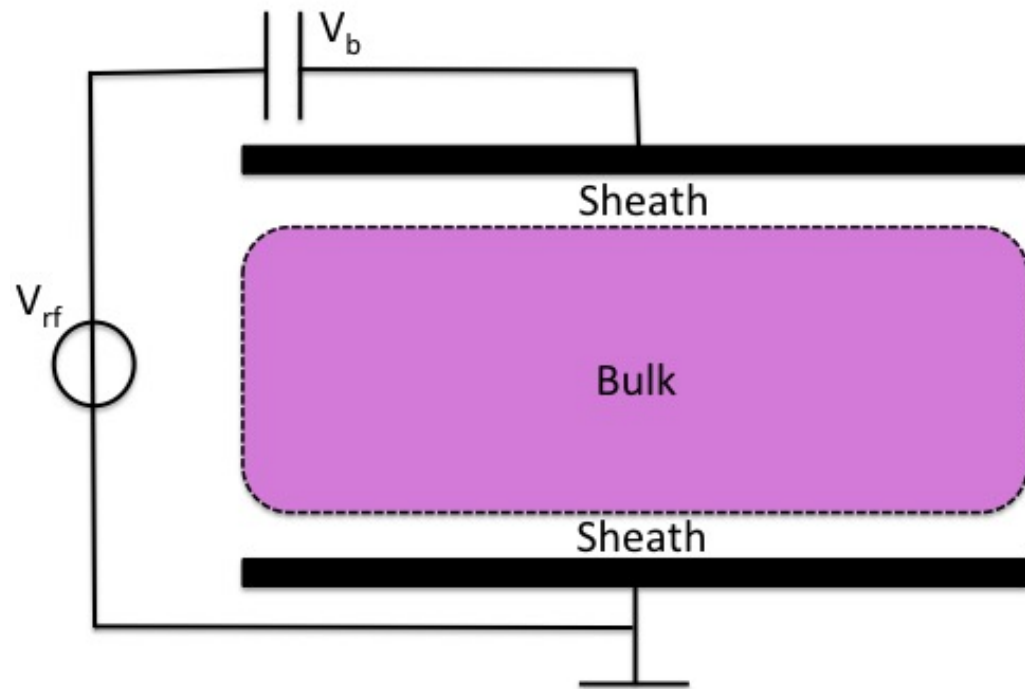
asymmetric (cylindrical)



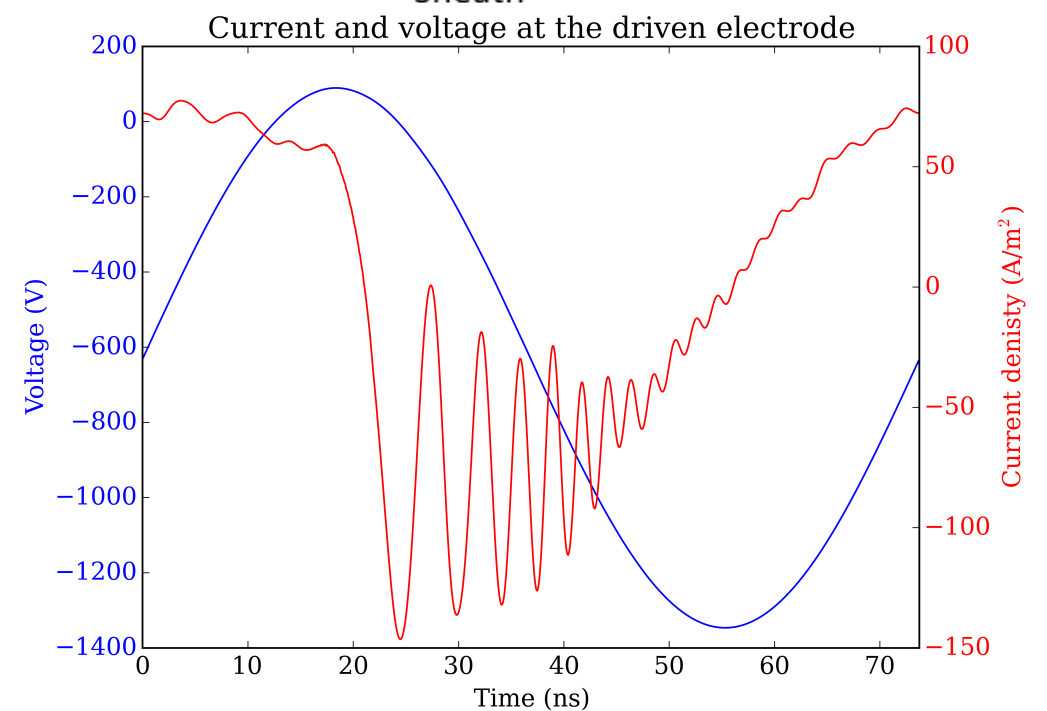
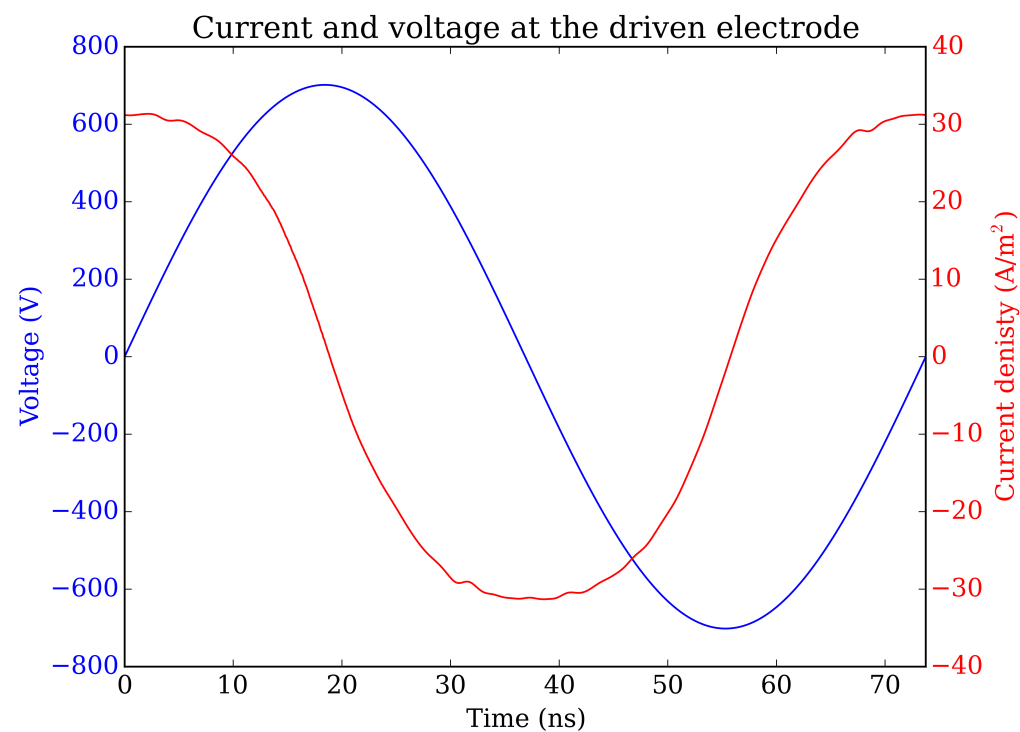
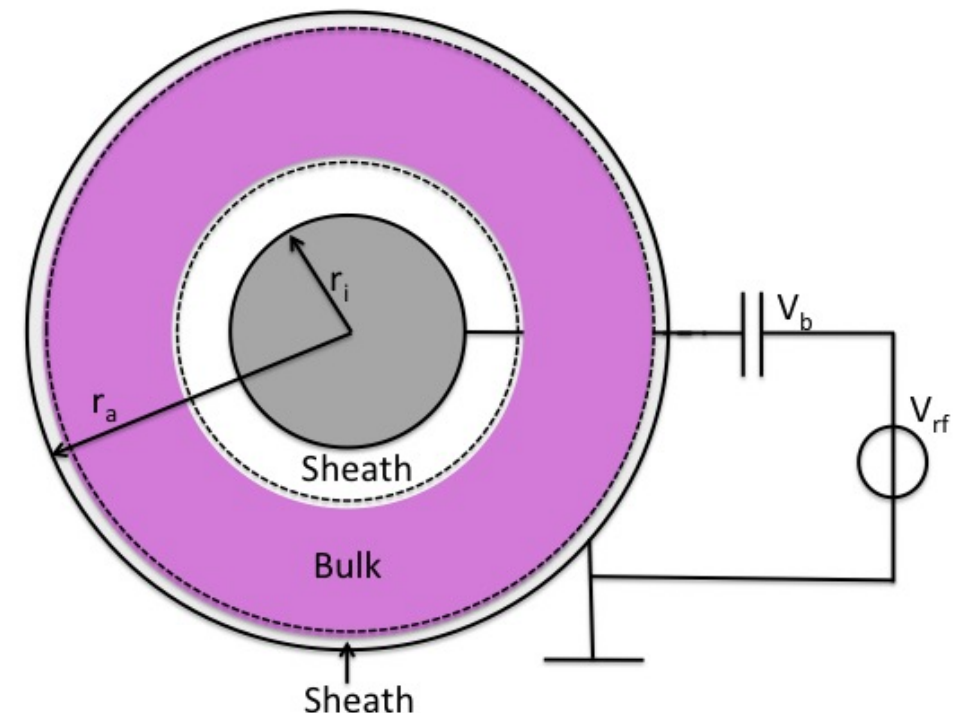
- cylindrical shells, purely 1d along the radial coordinate
- obtain a geometrical asymmetry and a self-consistent self-bias

Particle-In-Cell Simulation

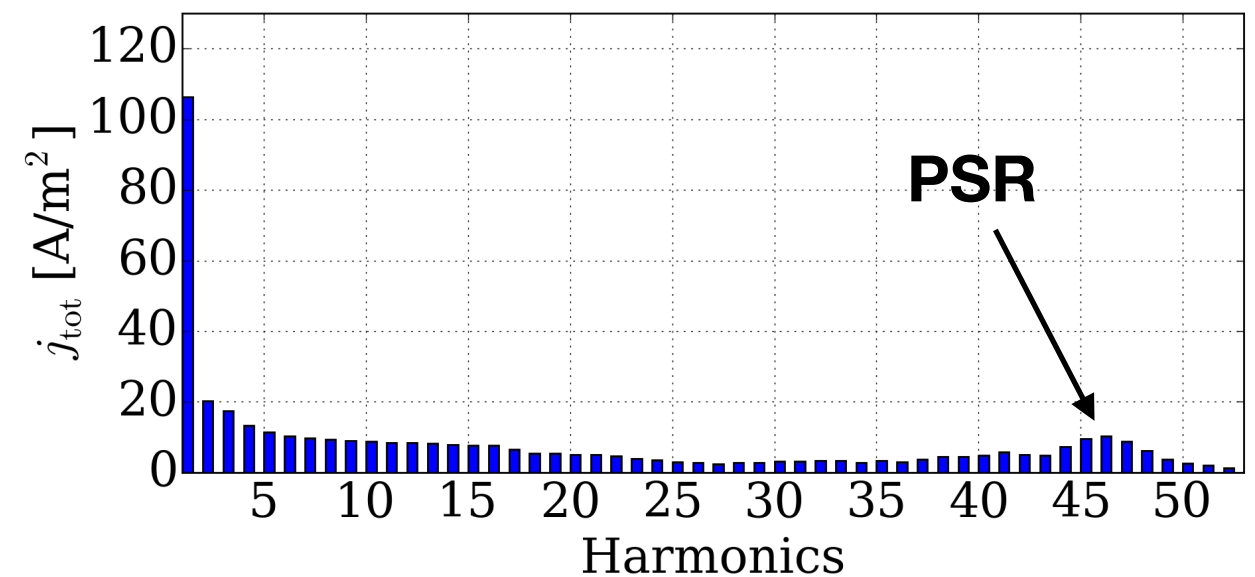
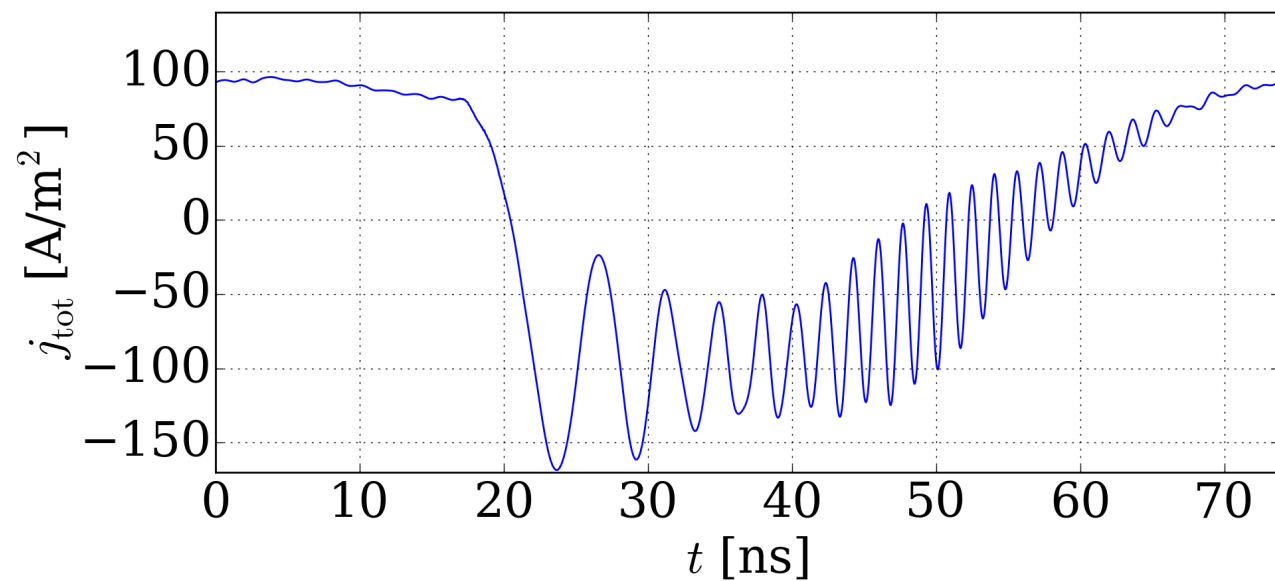
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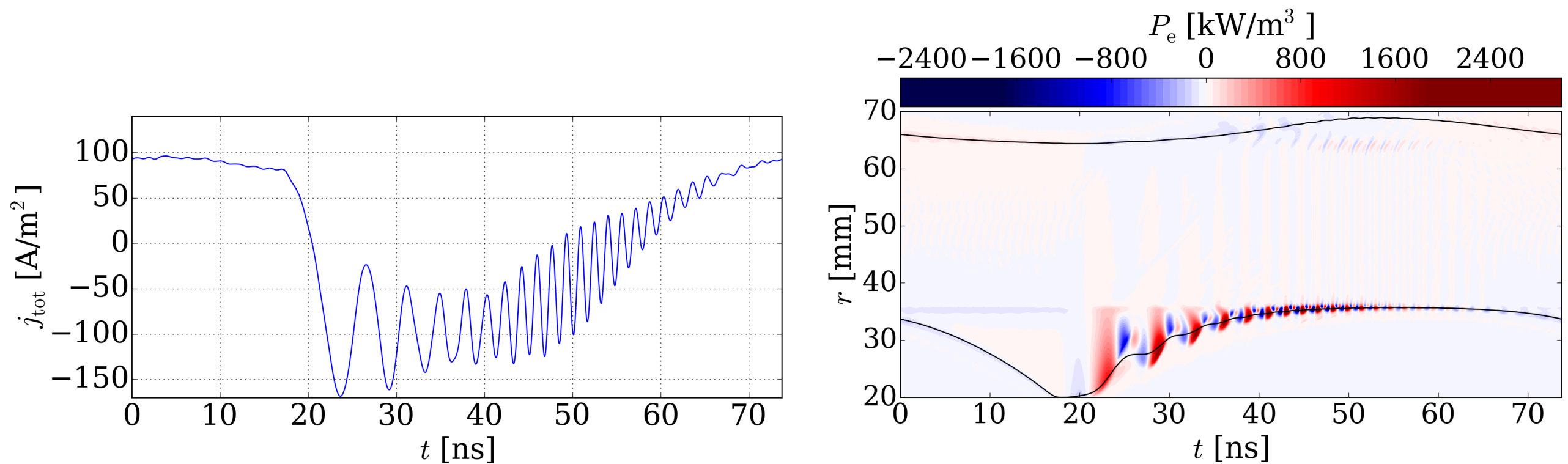


RF-current at the electrode



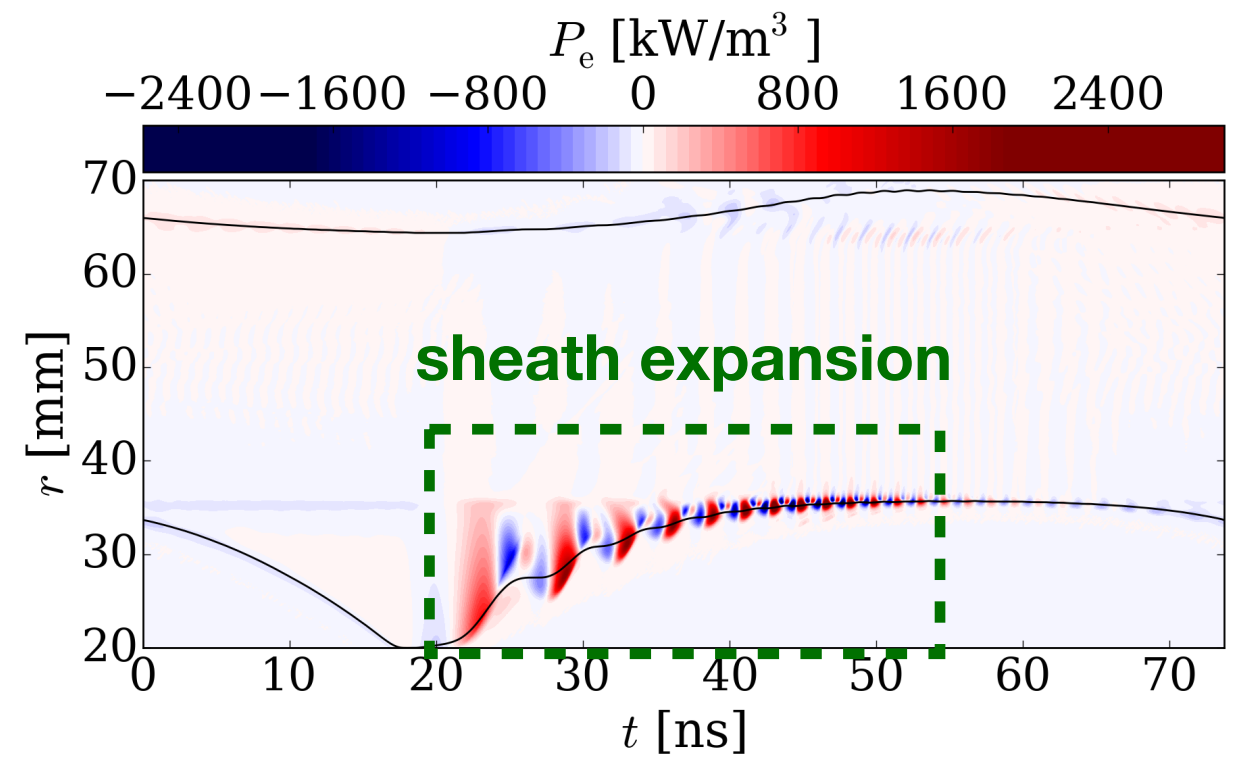
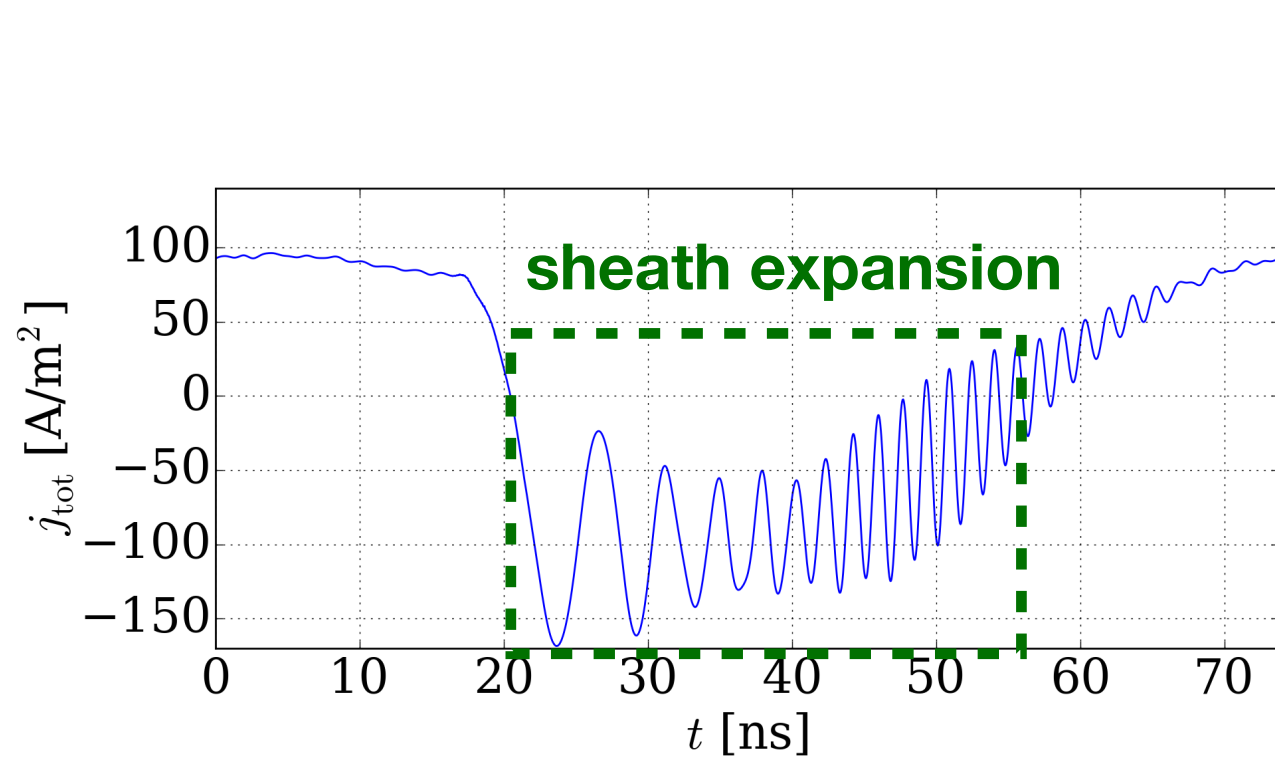
- asymmetric discharge: $f_{\text{RF}} = 13.56 \text{ MHz}$, $p = 1 \text{ Pa}$ argon, $L_{\text{gap}} = 50 \text{ mm}$, $V_{\text{RF}} = 1400 \text{ V}$
- strong nonlinear dynamics in the rf-current
- Fourier spectrum indicates higher harmonics which match the PSR frequency
- final goal: understand the origin of these oscillations

Electron power gain and loss



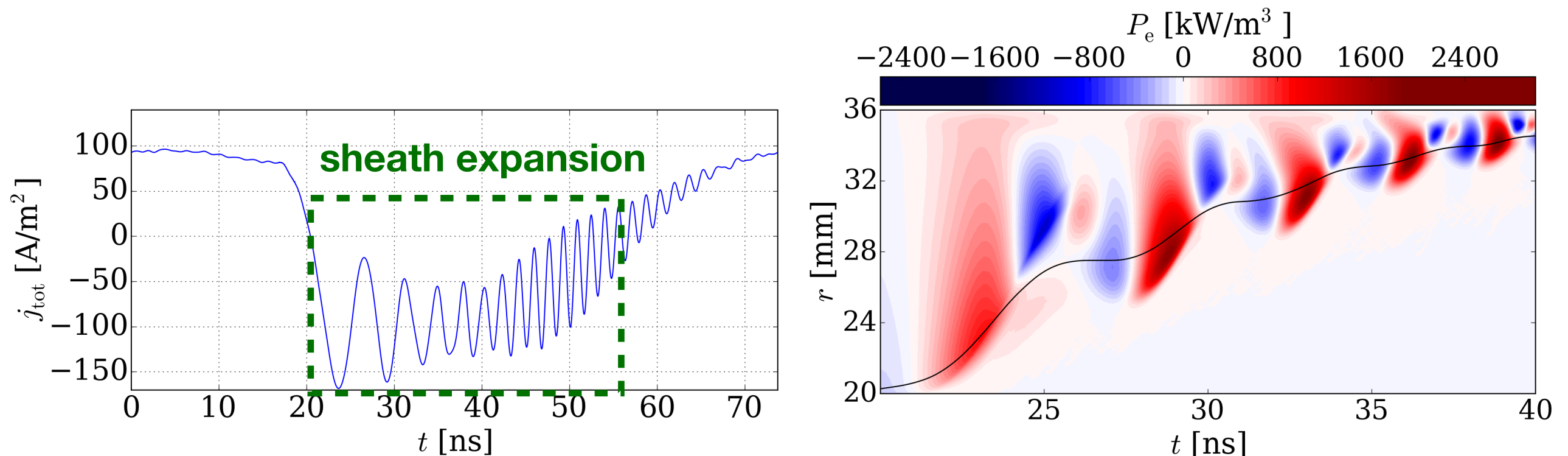
- nonlinear dynamics leads to a complex electron power gain and loss mechanism

Electron power gain and loss



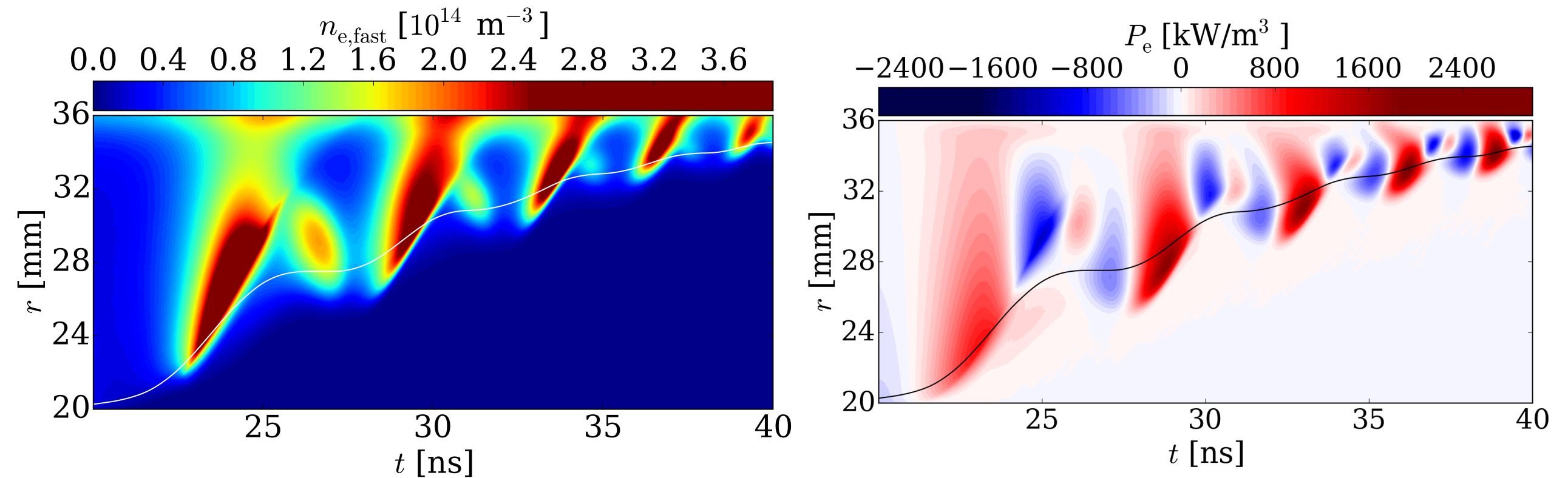
- nonlinear dynamics leads to a complex electron power gain and loss mechanism
- focus on phase of sheath expansion in order to explain the generation of harmonics

Electron power gain and loss



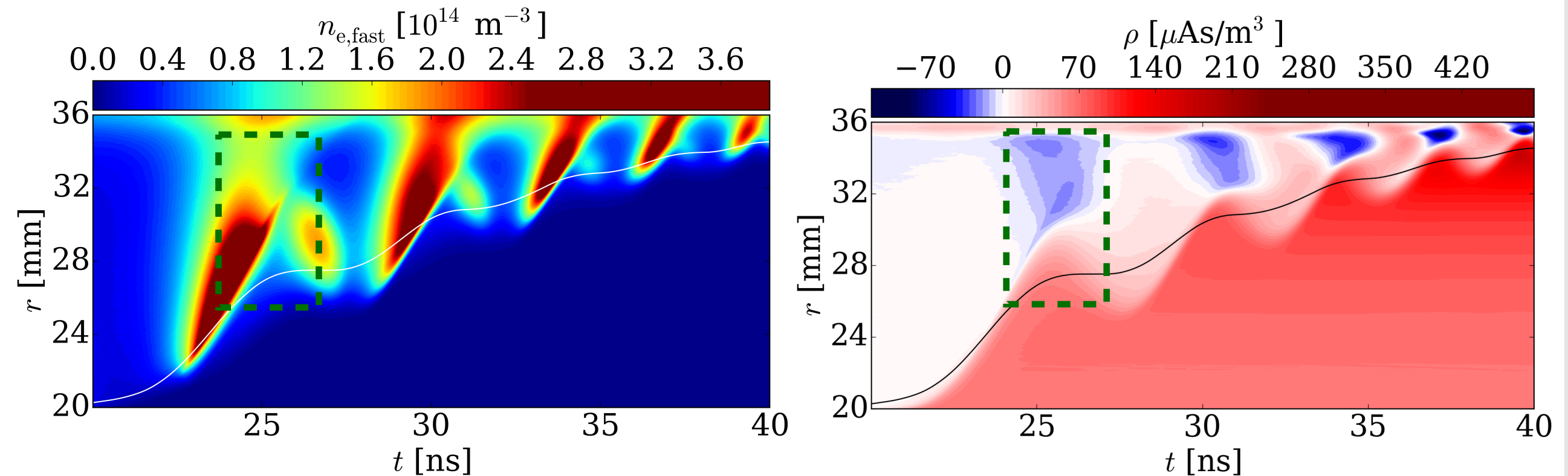
- nonlinear dynamics leads to a complex electron power gain and loss mechanism
- focus on phase of sheath expansion in order to explain the generation of harmonics
- change between electron power gain and loss during expansion
- what does kinetically happen during these power gain and loss phases?

Generation of electron beams



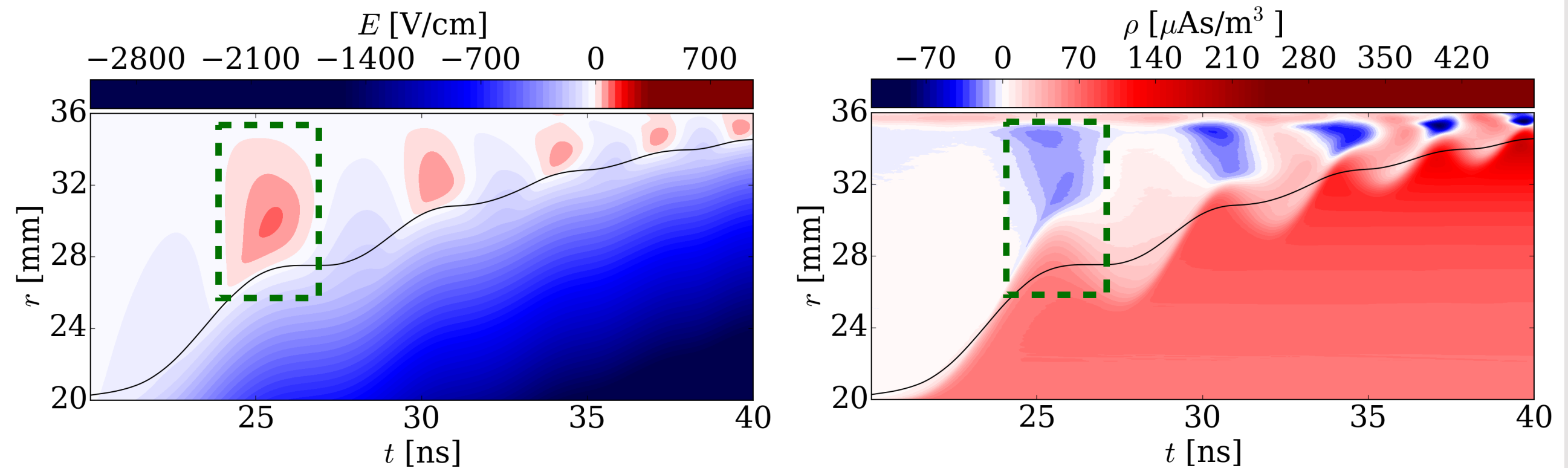
- during electron power gain, energetic beam electrons are generated
- due to the low pressure (1 Pa), they penetrate into the bulk almost collisionlessly
- additionally electrons are flowing back from the bulk into the sheath
- why are electrons attracted back to the sheath?

Formation of local charge densities

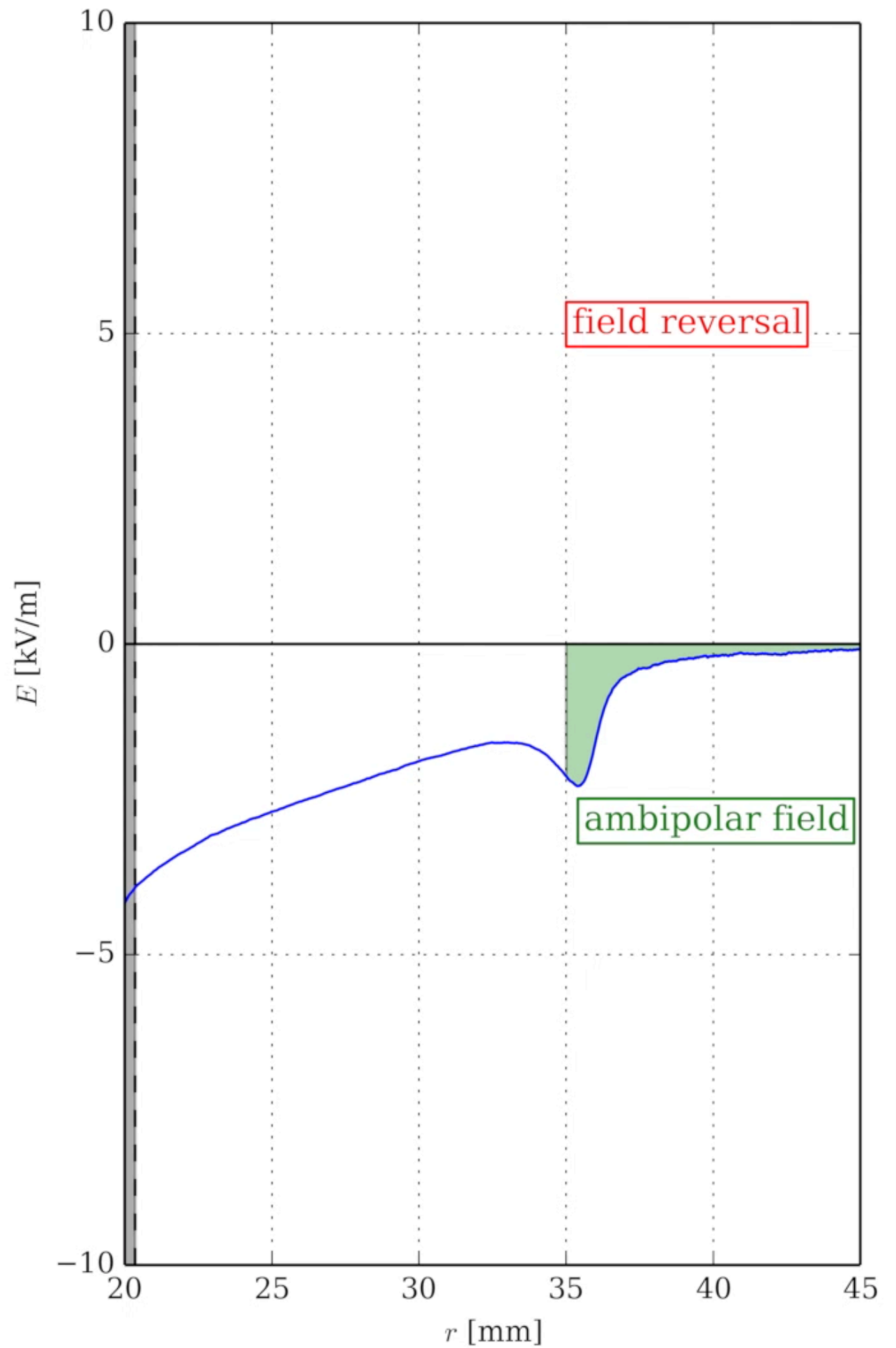
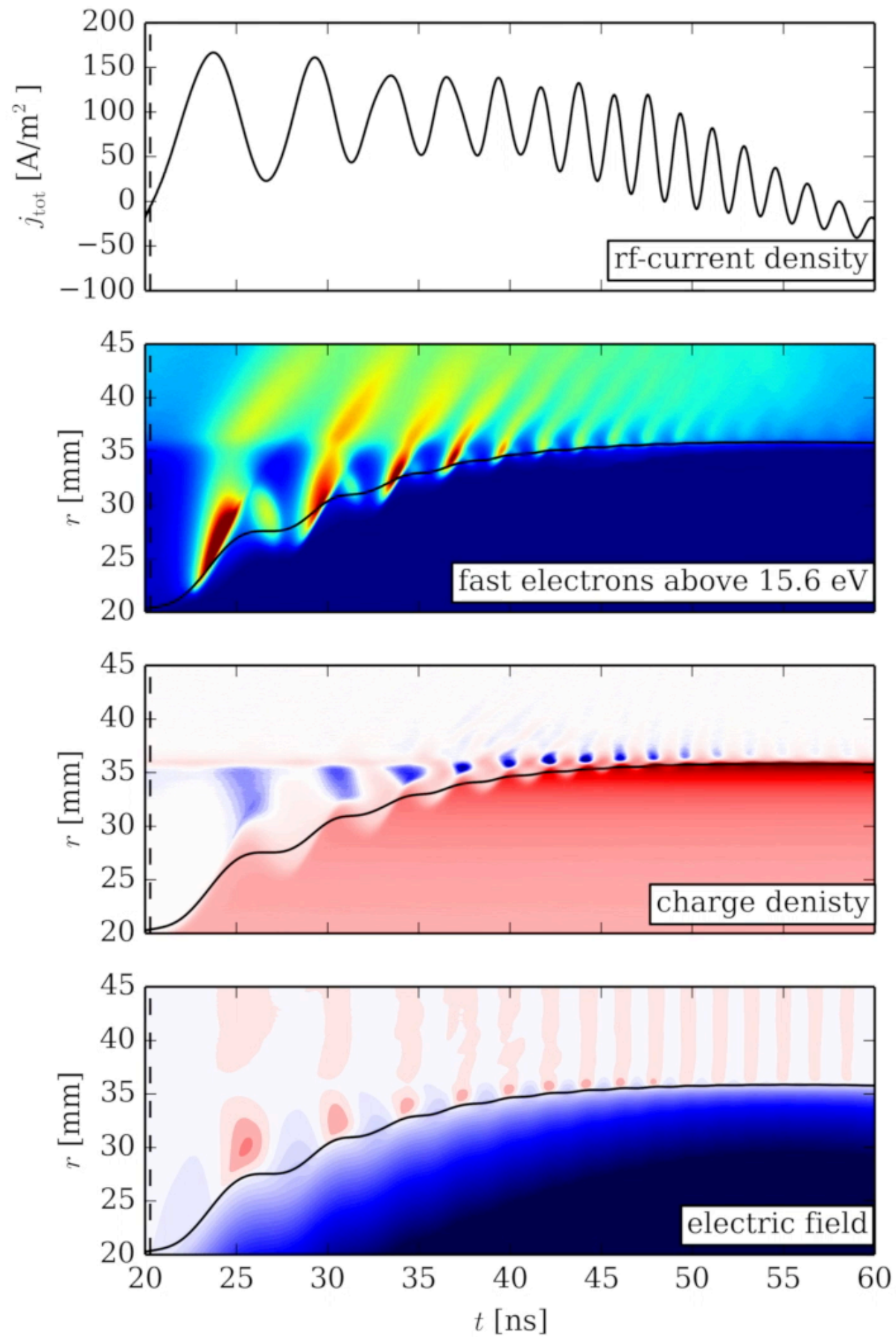


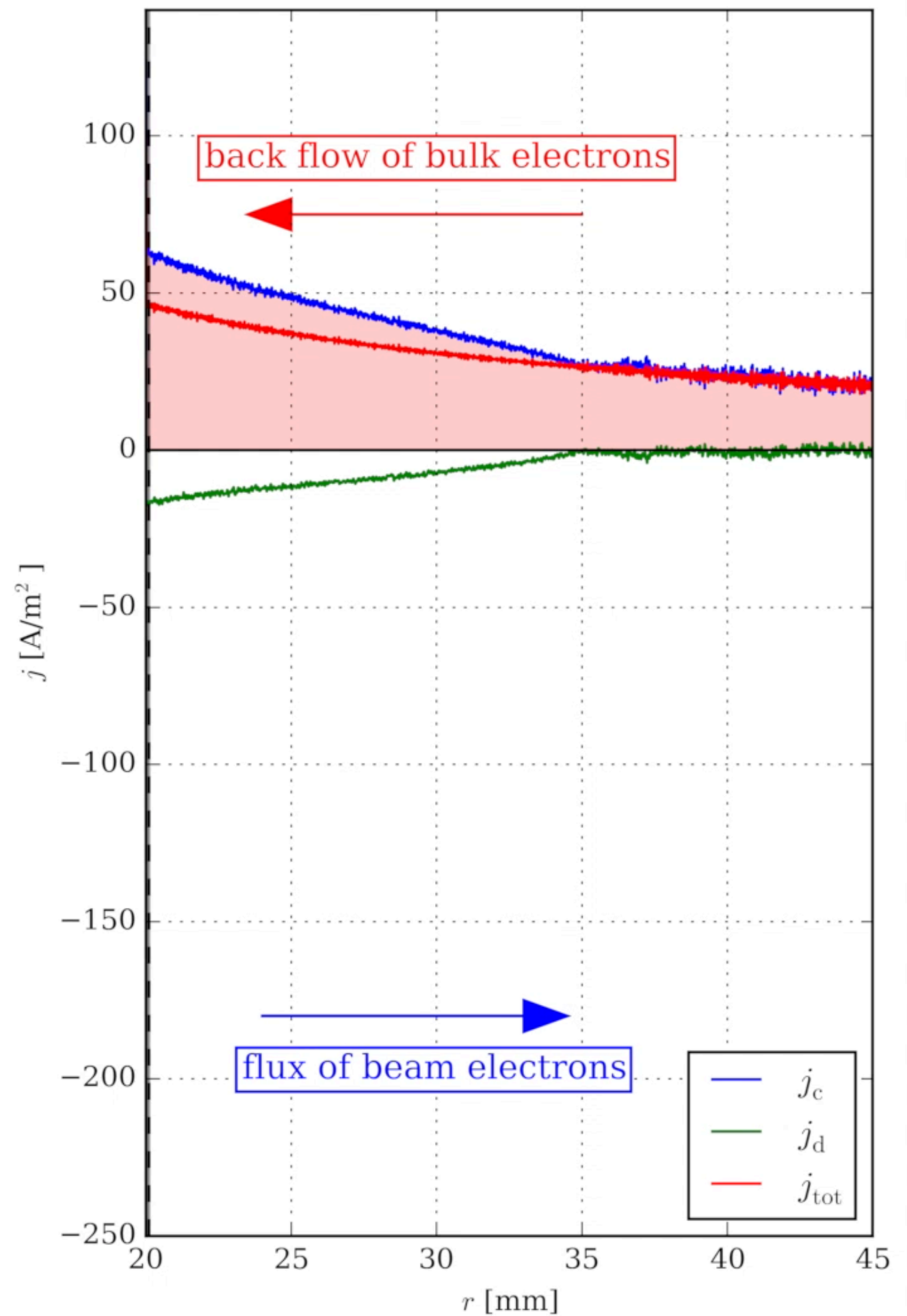
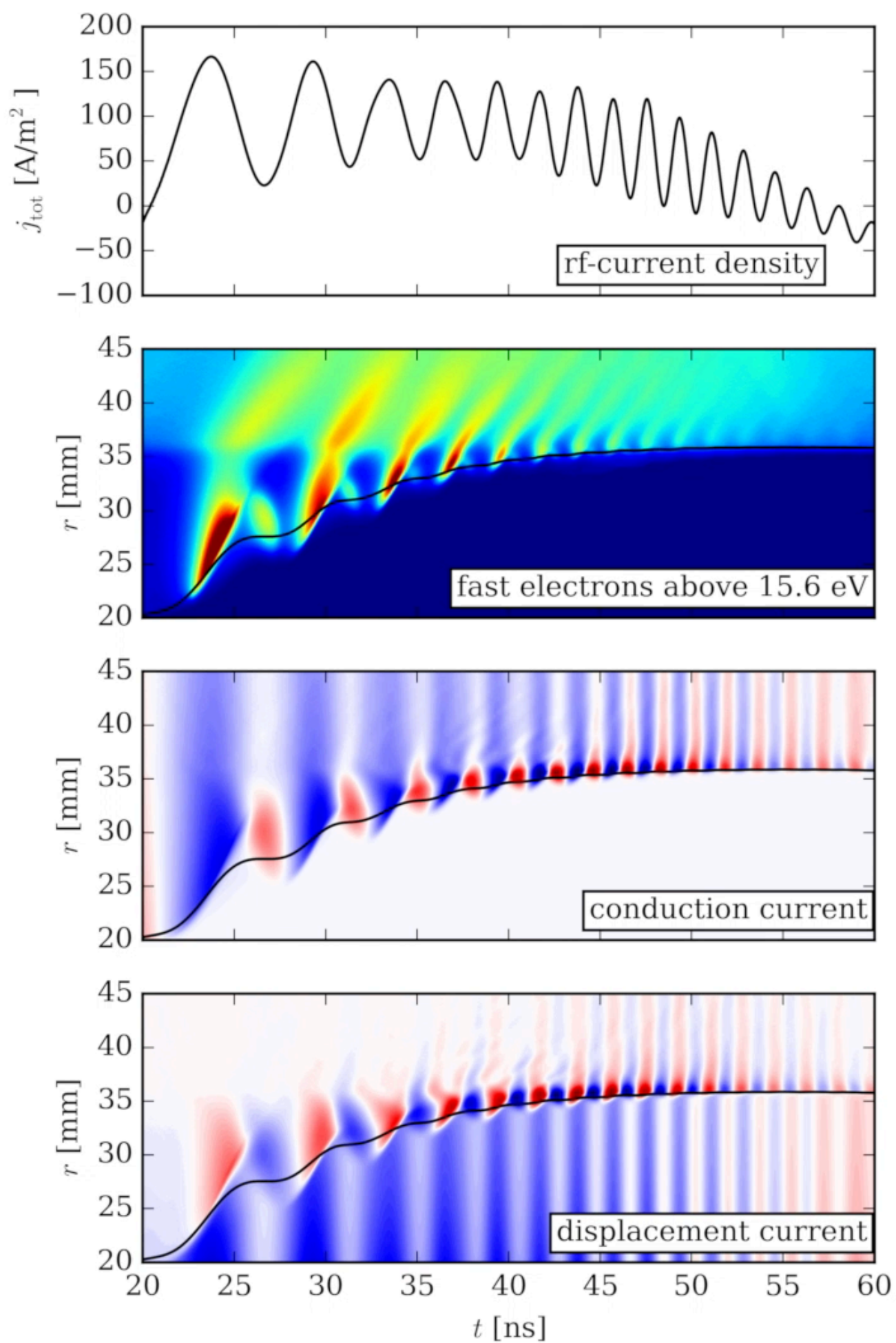
- very non-local regime, beam dynamics generates significant charge densities
- positive charge density is formed close to the sheath edge
- negative charge density is formed in front of the sheath edge
- what is the result of such a charge difference?

Electric field reversal

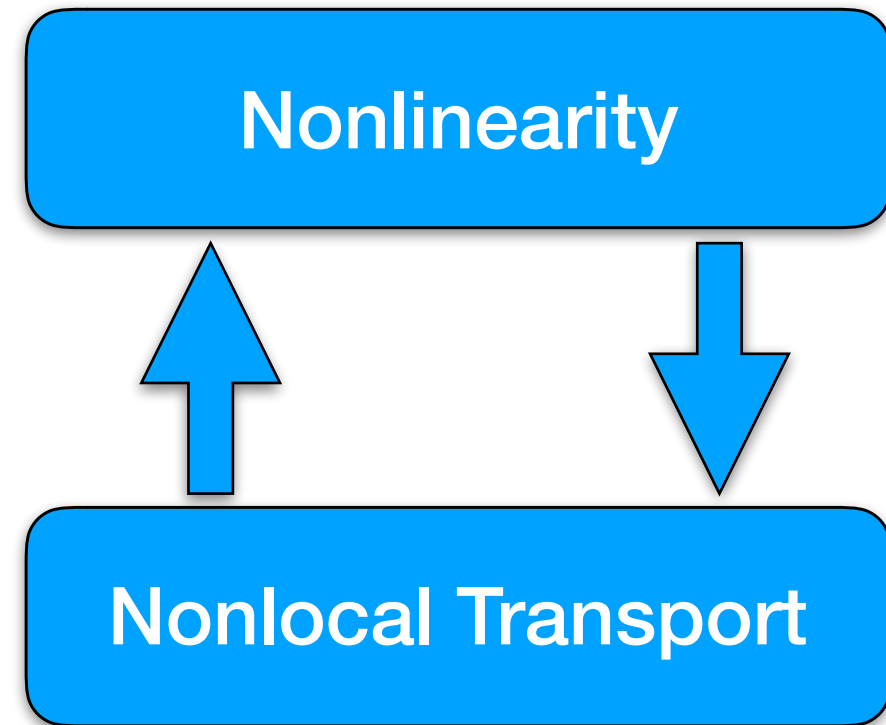
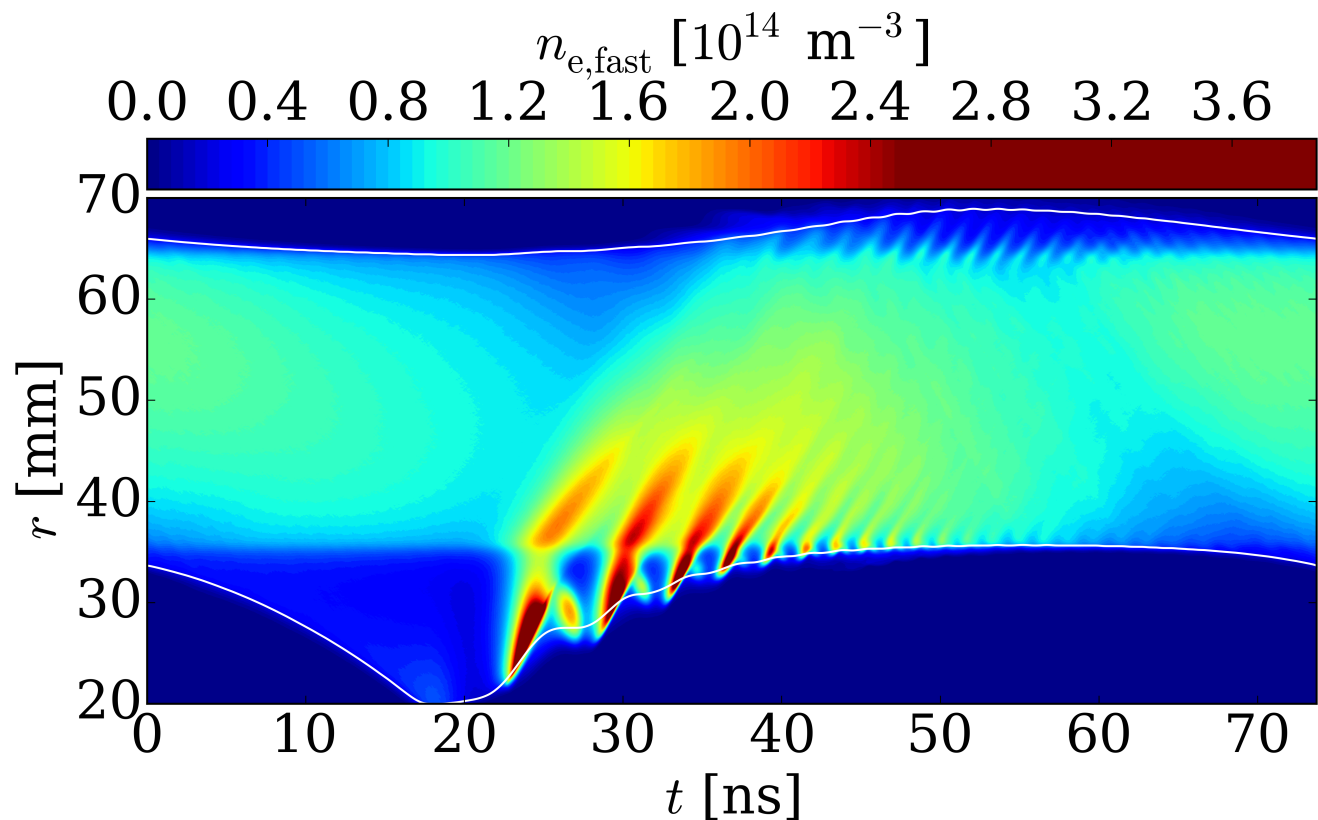


- formation of charge densities generates an electric field reversal in front of the sheath
- this positive electric field accelerates bulk electrons back to the sheath
- these bulk electrons interact with the nonlinear sheath
- this interaction strongly affects the global rf-current and generates harmonics



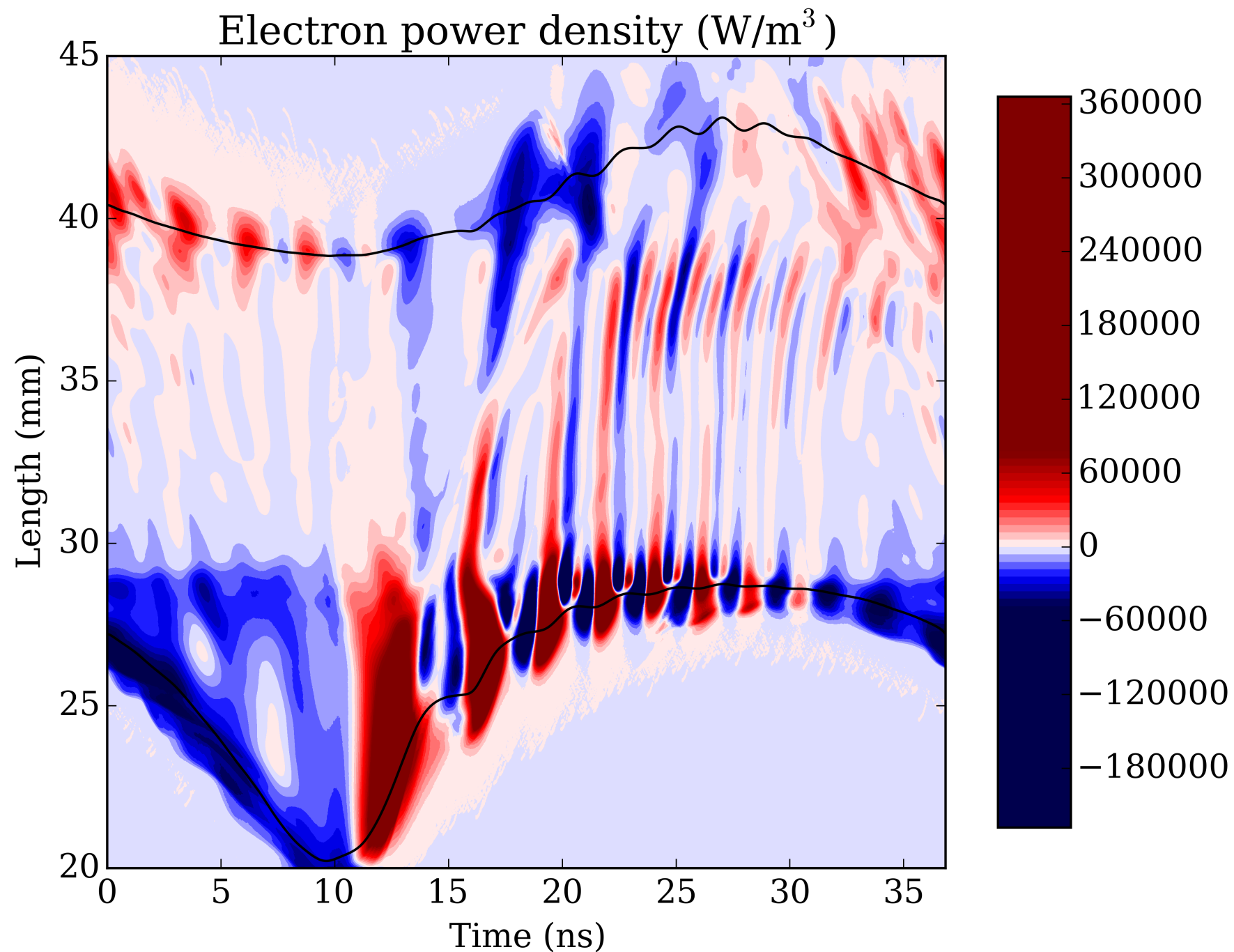


Summary: Nonlocal and nonlinear dynamics



- nonlocal transport of beam electrons plays an important role at low pressures
- process optimization by controlling the impingement phase of beam electrons
- penetration of beam electrons leads to local field reversals in front of the sheath
- backflow of bulk electrons generates harmonics in the RF-current
- generation of multiple beam electrons as well as the excitation of electrostatic waves represent the nonlinearity of CCRF discharges

Conclusion



- electron heating in CCRF discharges at low pressures is still not fully understood
- nonlinear and nonlocal dynamics contribute significantly to the electron dynamics