

The Difference Between Electron Heating and Power Absorption In Capacitively Coupled Plasmas

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Electron dynamics in low pressure capacitively coupled radio frequency discharges

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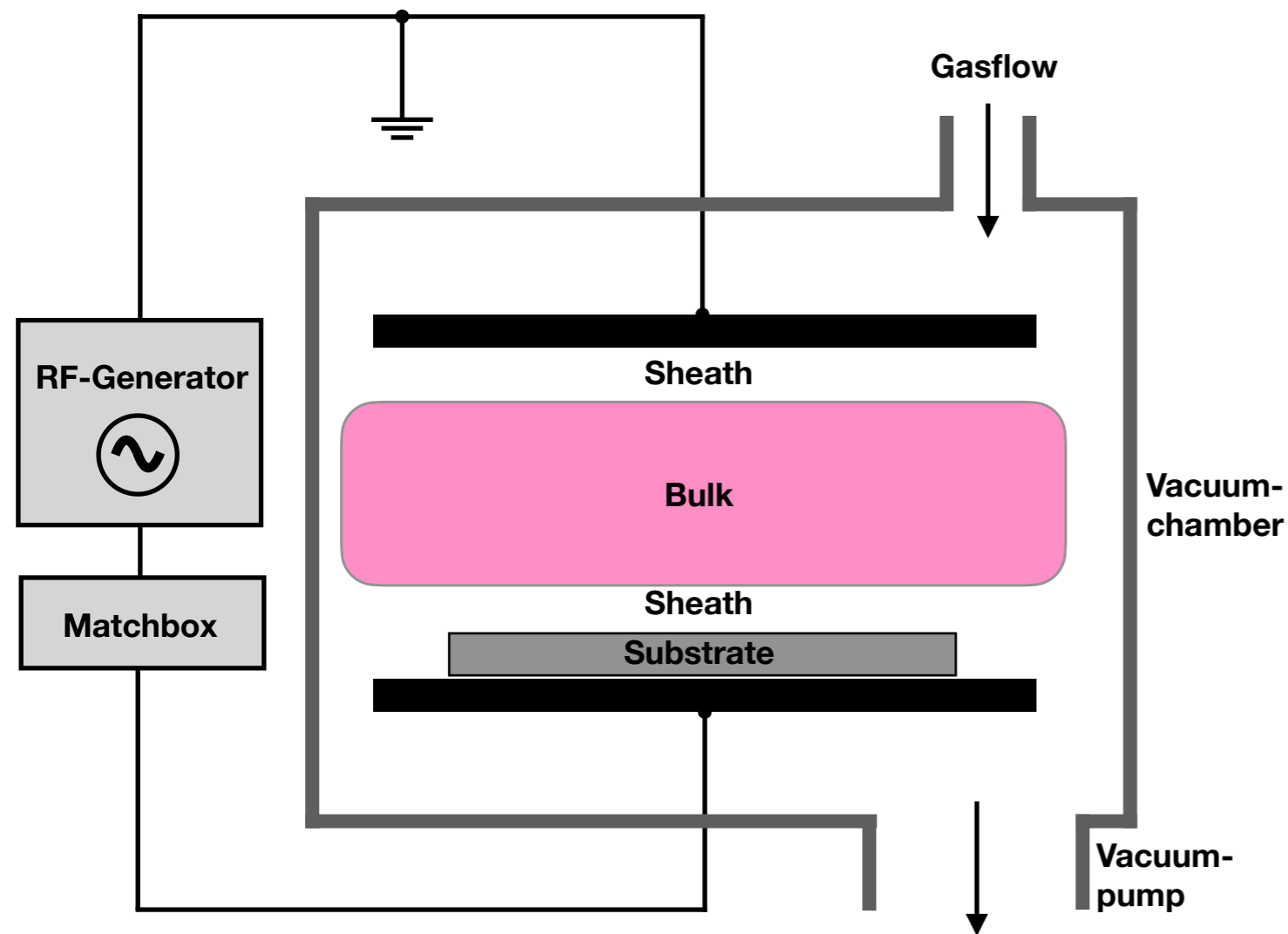


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This presentation is based on the tutorial „Electron dynamics in low pressure capacitively coupled radio frequency discharges“, which has already been published as a featured article in Journal of Applied Physics

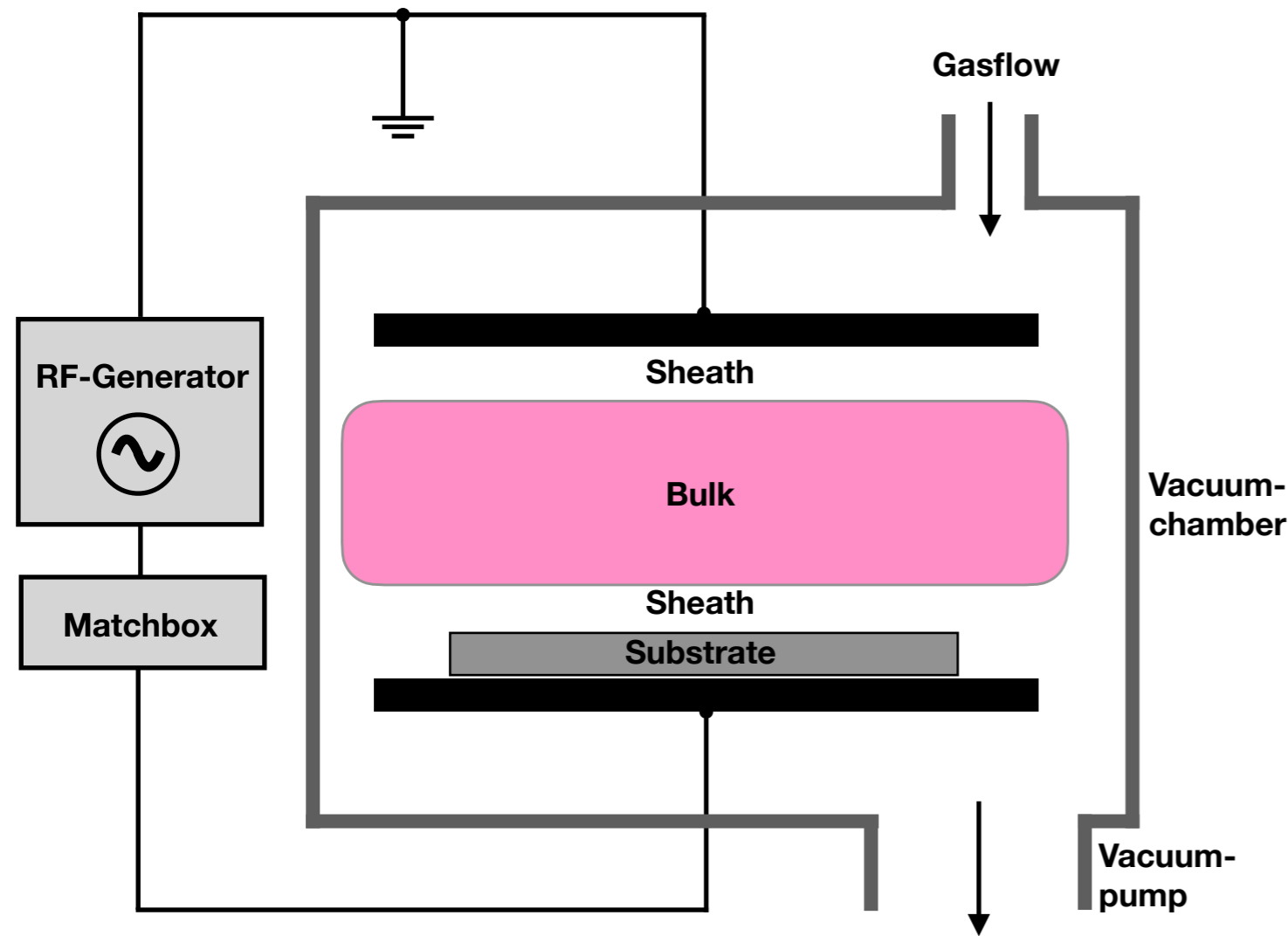
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Motivation: Electron Dynamics

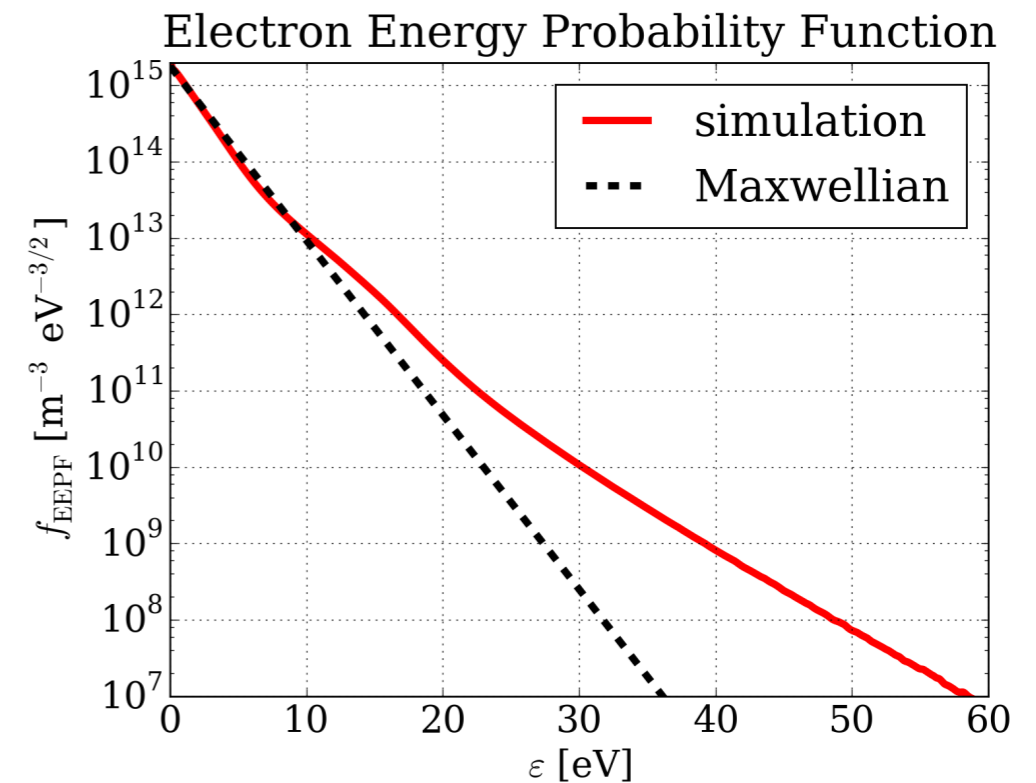


- control of the electrons in order to optimize the industrial relevant discharges
- however, electrons at low pressures (< 10 Pa) indicate a strong anisotropy

Motivation: Electron Dynamics



$p = 3 \text{ Pa (argon)}$ $V_0 = 500 \text{ V}$
 $f = 13.56 \text{ MHz}$ $L_{\text{gap}} = 50 \text{ mm}$

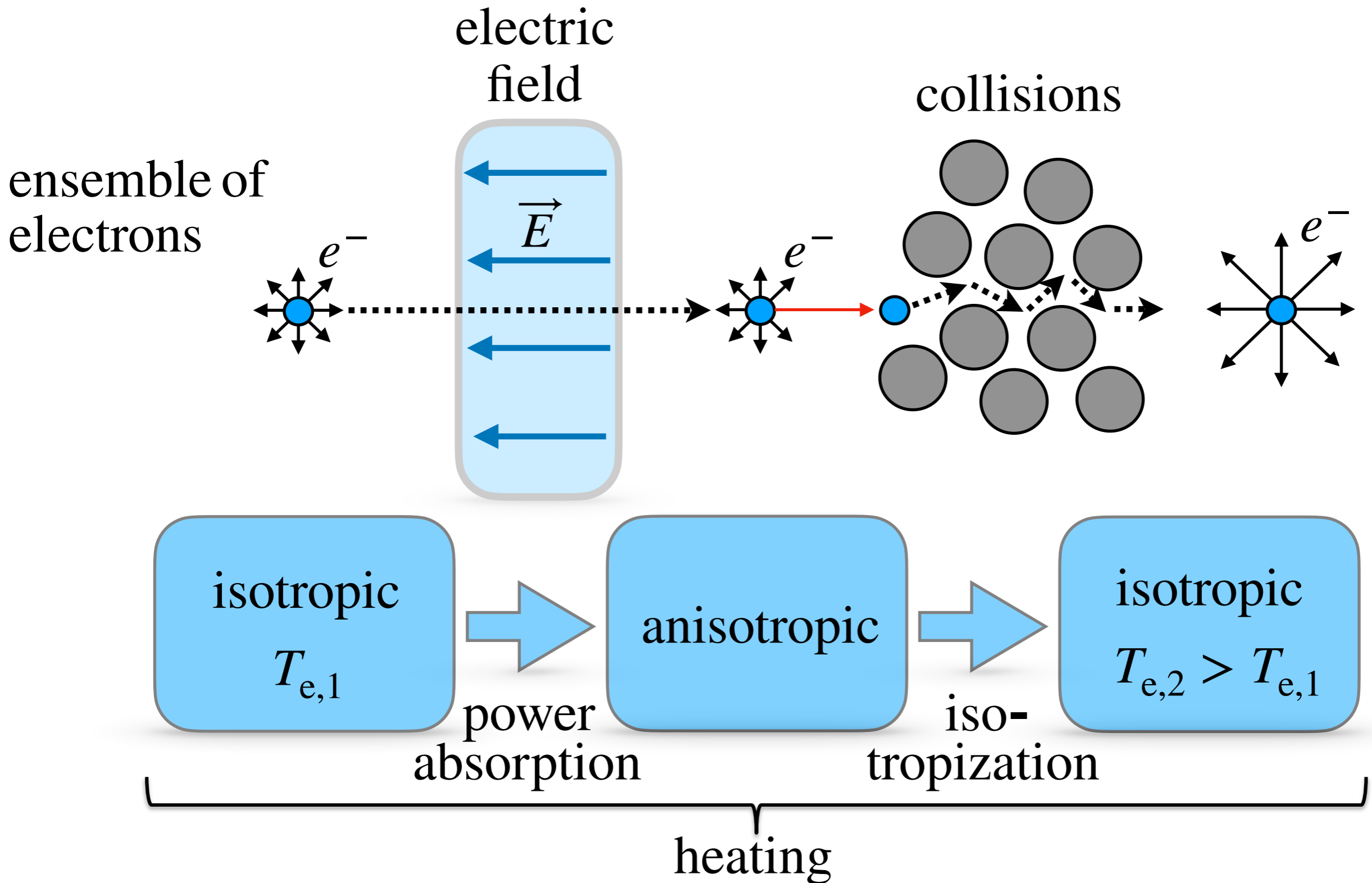


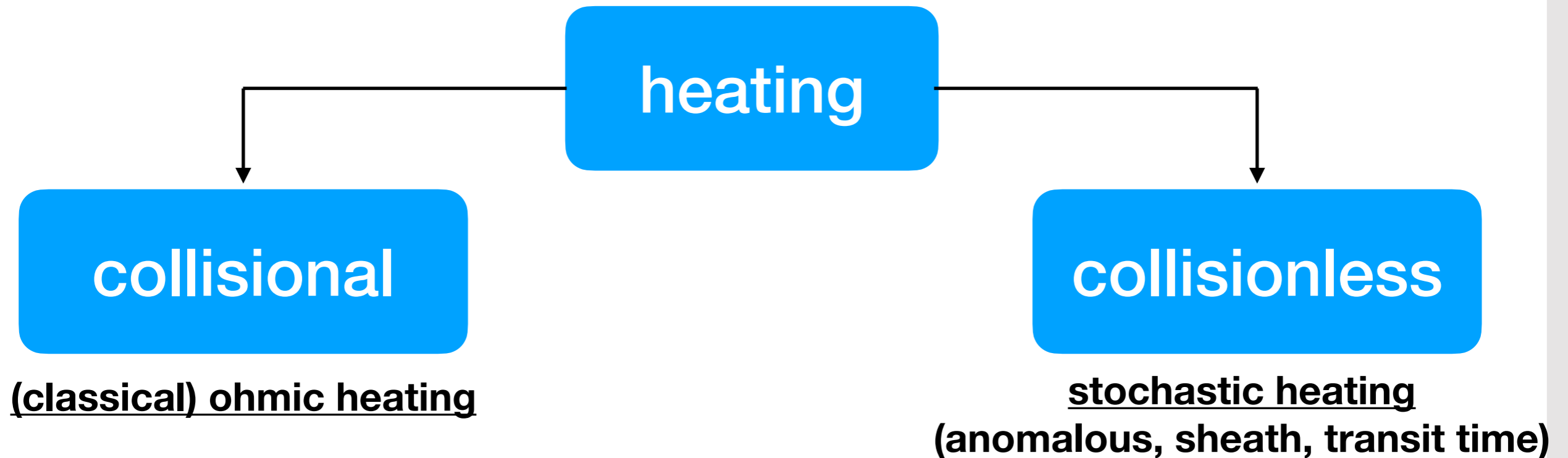
- control of the electrons in order to optimize the industrial relevant discharges
- however, electrons at low pressures ($< 10 \text{ Pa}$) indicate a strong anisotropy
- electron distribution function strongly differs from a Maxwellian distribution
- challenging to understand and control the electron dynamics

Goal of this Work

- 1. How do the electrons gain and lose their energy in an electric field?
Traditionally, how does the electron heating really work?**
- 2. How to deal with the thermodynamic concept of the electron temperature in such a very nonlocal and anisotropic regime?**

What is actually Electron Heating?





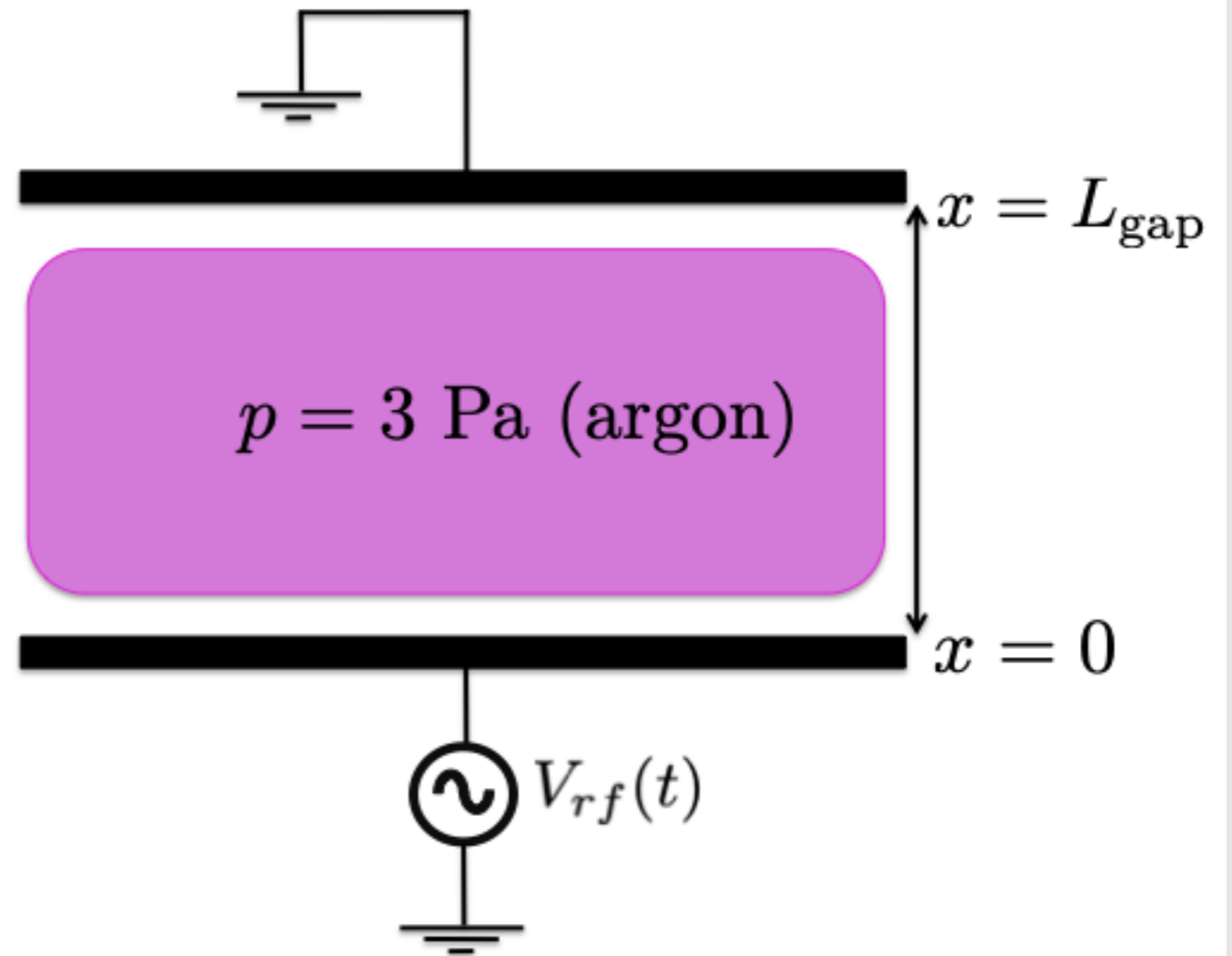
Further terminologies:

- nonlinear electron resonance heating
- pressure heating
- ambipolar heating
- bounce-resonance-heating
- secondary electron heating
- nonlinear wave-particle heating

**Finally, most of the terms describe the same mechanism:
The particle interaction with a time-varying electric field!
However, no coherent terminologies!**

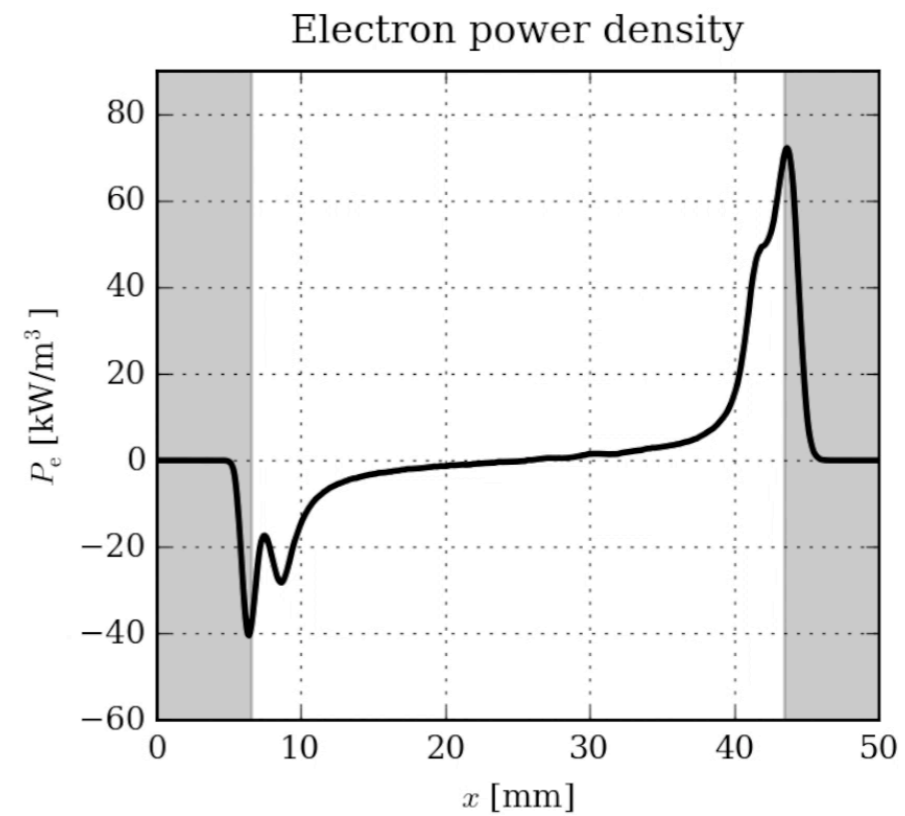
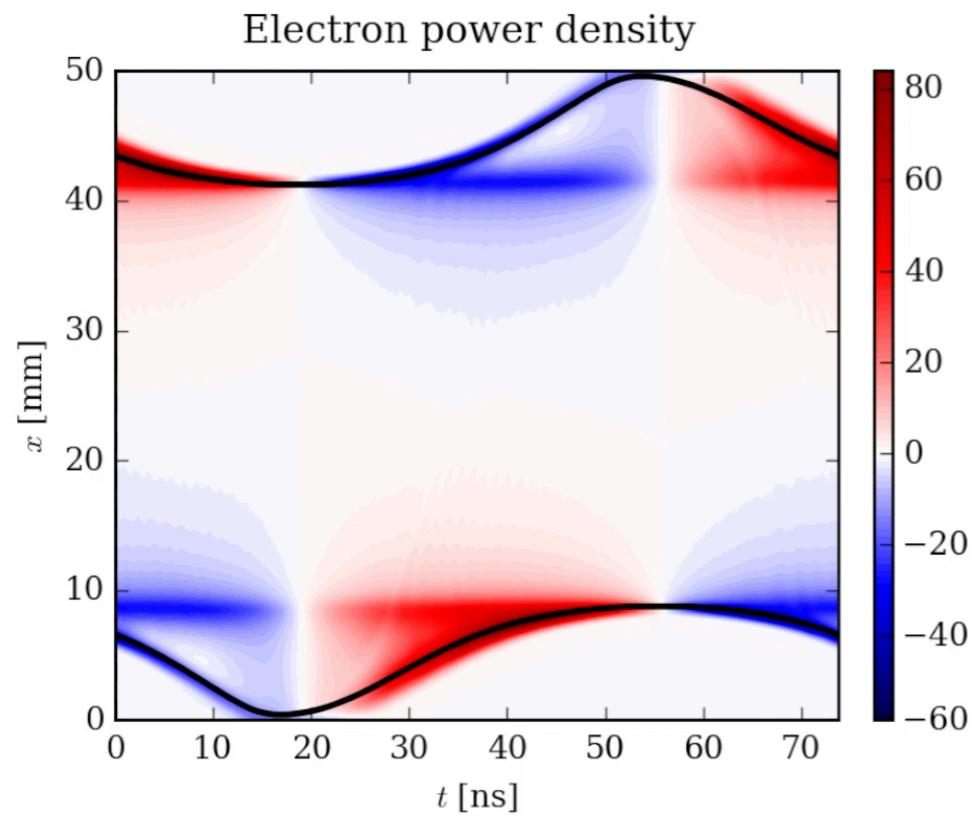
Simulation Setup

- 1d3v PIC/MCC simulation
- planar, parallel and infinite electrodes
- axial symmetric, translational invariant in y and z
- only parallel and perpendicular directions
- argon gas pressure: 3 Pa
- gap size: 50 mm
- driving frequency: 13.56 MHz
- voltage amplitude: 500 V
- no surface models



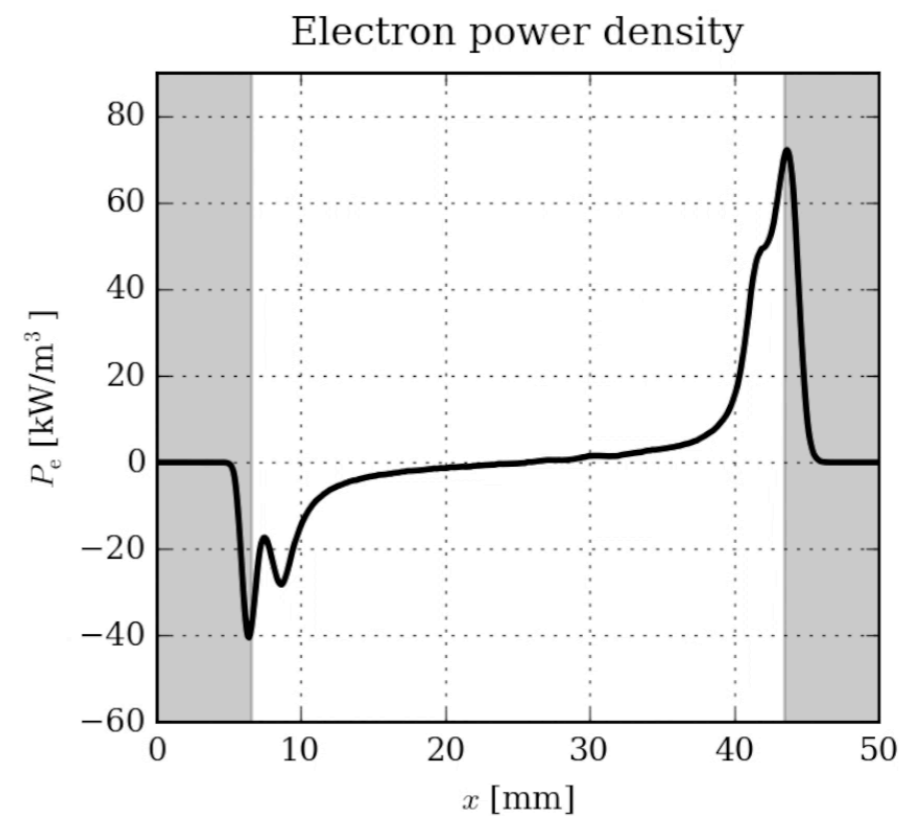
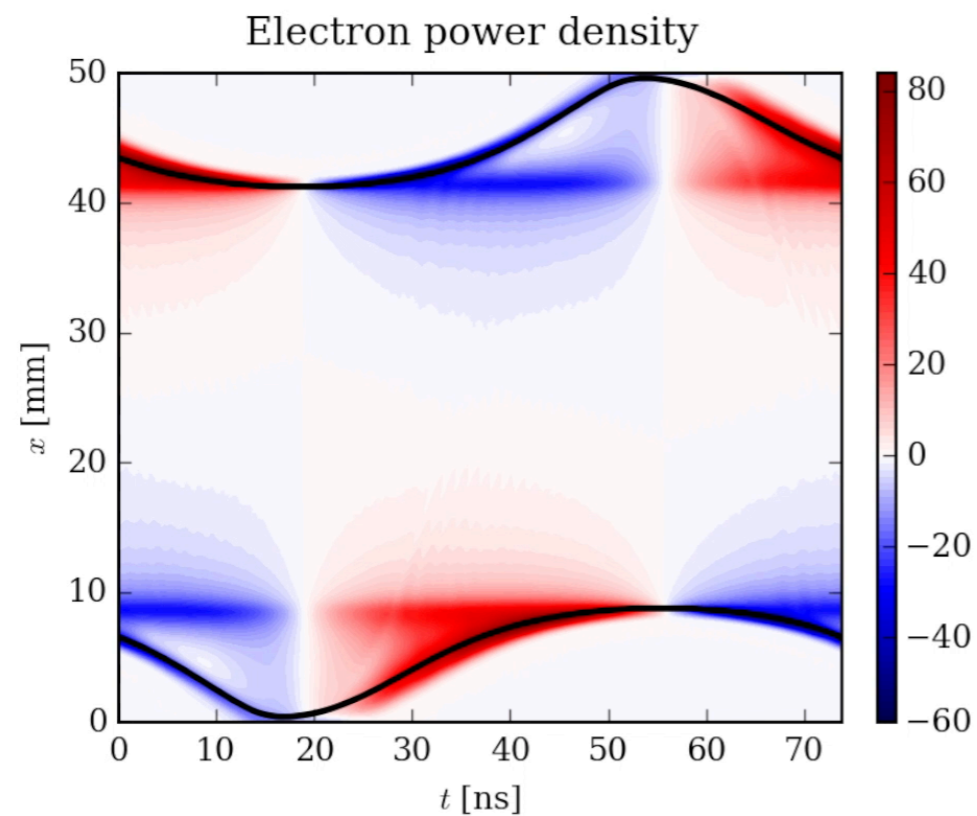
Electron Power Absorption

- electron power density: $P_e = j_e \cdot E$



Electron Power Absorption

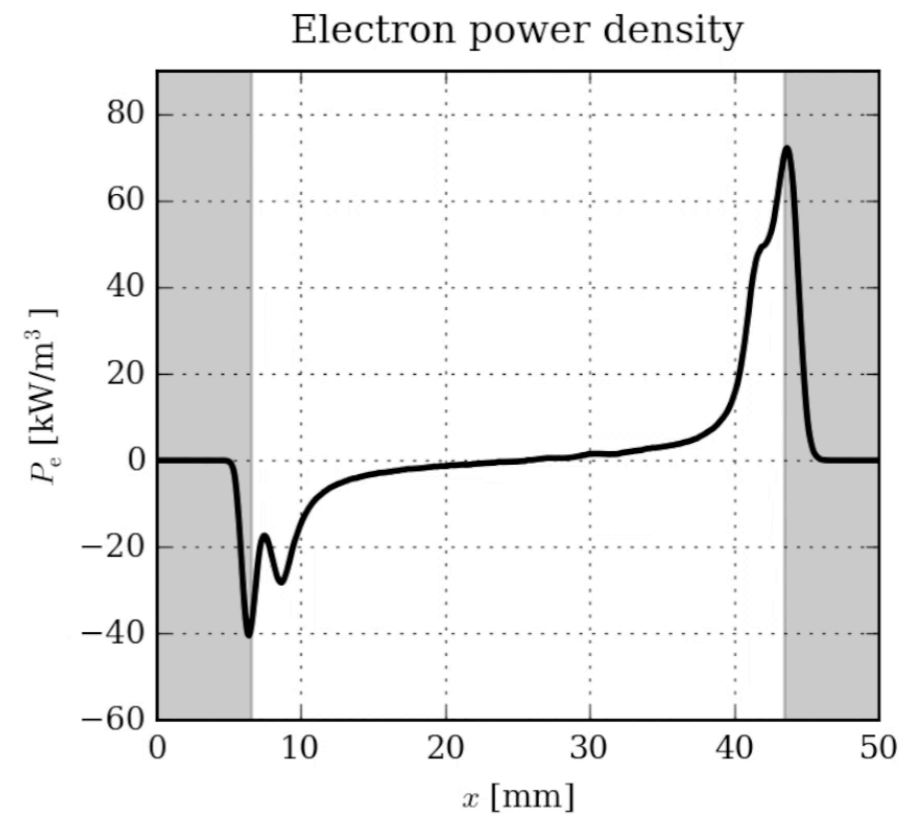
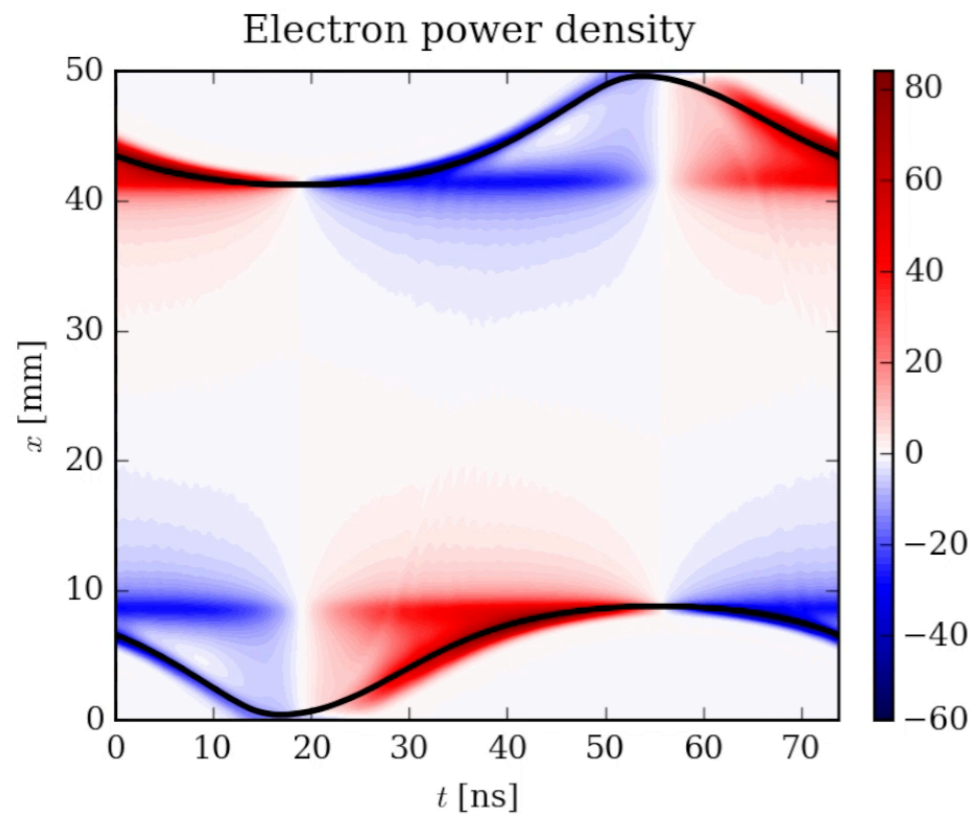
- electron power density: $P_e = j_e \cdot E$
- dominant power absorption near sheath edge (black solid line)
- dominant power absorption in the ambipolar region in front of the sheath edge
- how to study the electron power absorption mechanism in detail?



Electron Power Absorption

momentum balance
in x-direction (parallel):

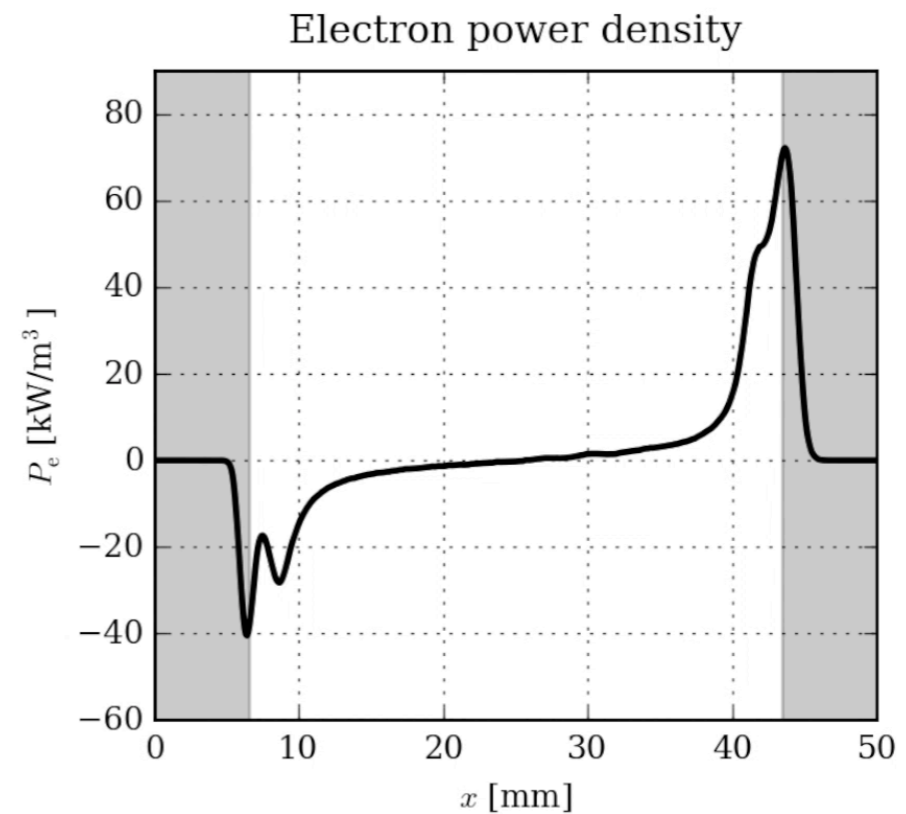
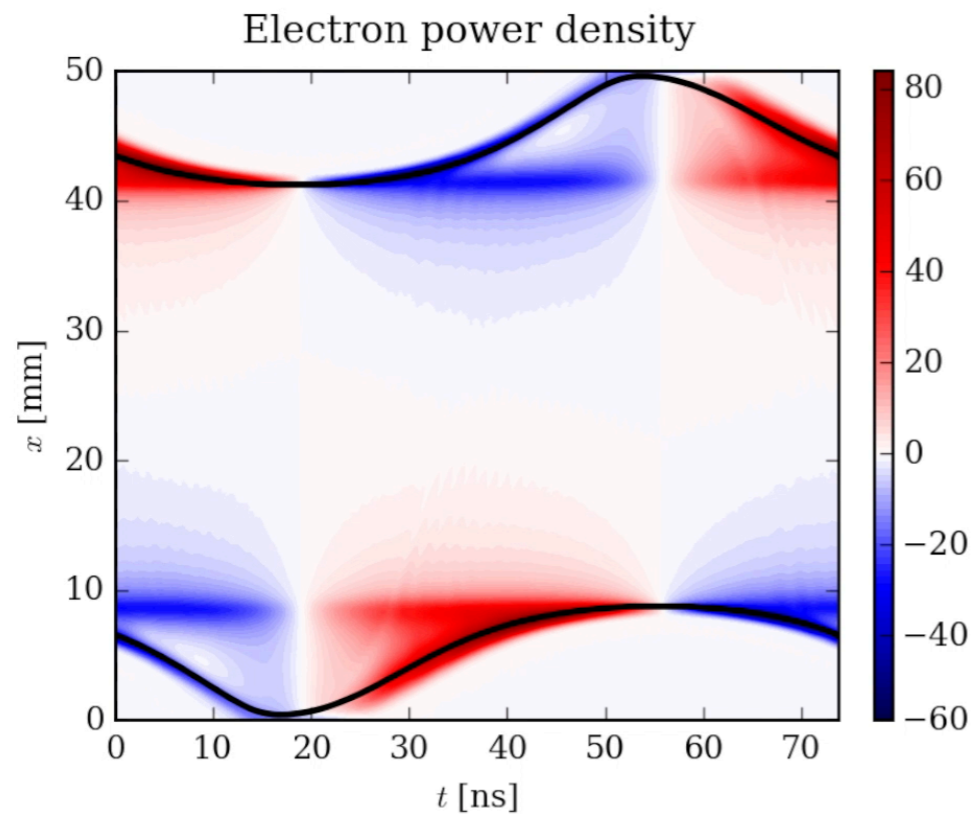
$$m_e \frac{\partial(n_e u_{\parallel})}{\partial t} + m_e \frac{\partial(n_e u_{\parallel}^2)}{\partial x} + \frac{\partial p_{\parallel}}{\partial x} = -en_e E_{\parallel} - \Pi_{c\parallel}$$



Electron Power Absorption

solving for the electric field:

$$E_{\parallel} = \underbrace{-\frac{m_e}{n_e} \left(\frac{\partial(u_{\parallel} n_e)}{\partial t} + \frac{\partial(n_e u_{\parallel}^2)}{\partial x} \right)}_{E_{\text{in}}} - \underbrace{\frac{1}{en_e} \frac{\partial p_{\parallel}}{\partial x}}_{E_{\text{pr}}} - \underbrace{\frac{1}{n_e e} \Pi_{c\parallel}}_{E_{\text{Ohm}}}$$

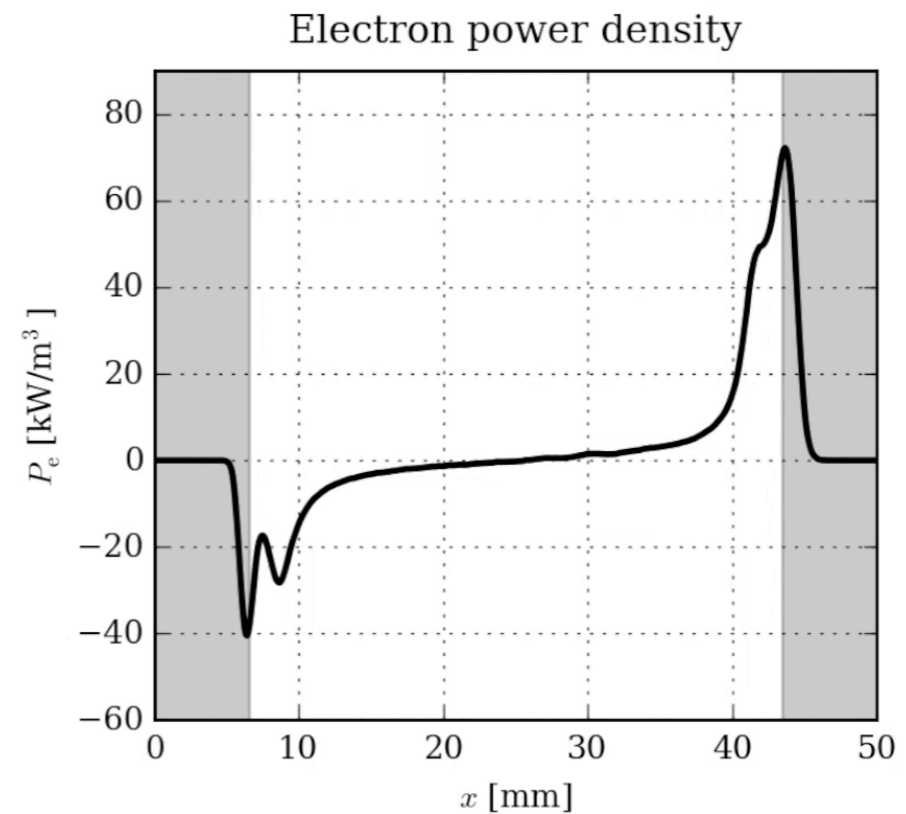
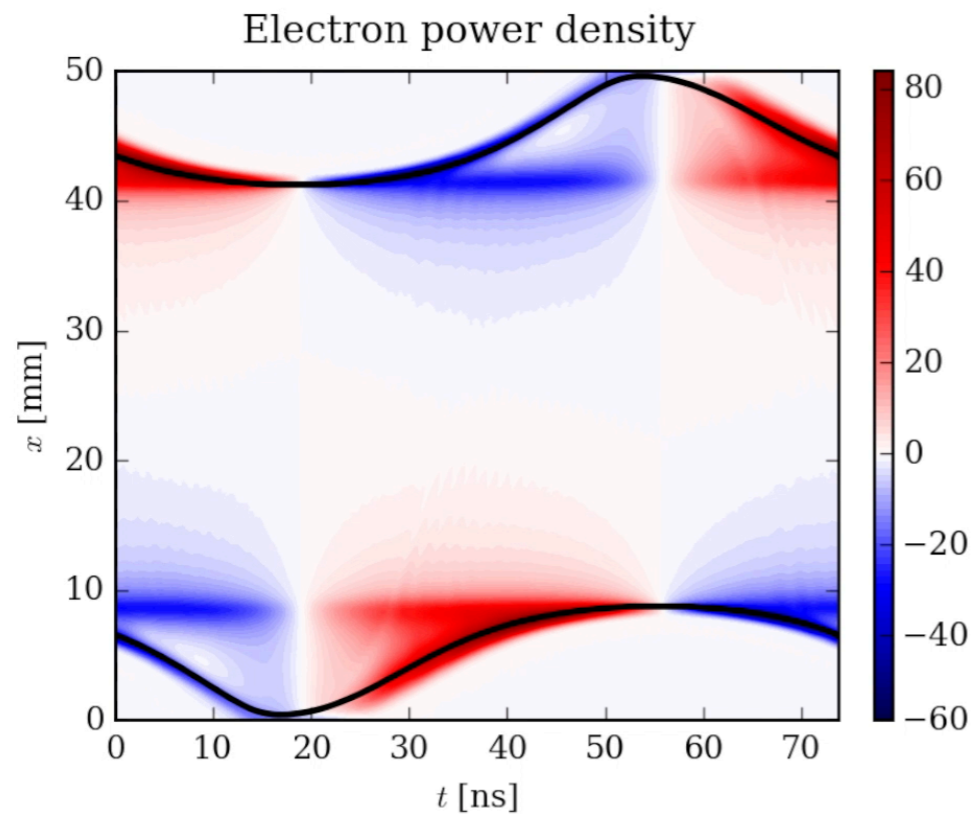


Electron Power Absorption

multiply by the
electron current:

$$\underbrace{j_e E}_{P_e} = \underbrace{j_e E_{in}}_{P_{in}} + \underbrace{j_e E_{pr}}_{P_{pr}} + \underbrace{j_e E_{Ohm}}_{P_{Ohm}}$$

$$\underbrace{P_{in} + P_{pr}}_{P_{collisionless}} \quad \underbrace{P_{Ohm}}_{P_{collisional}}$$

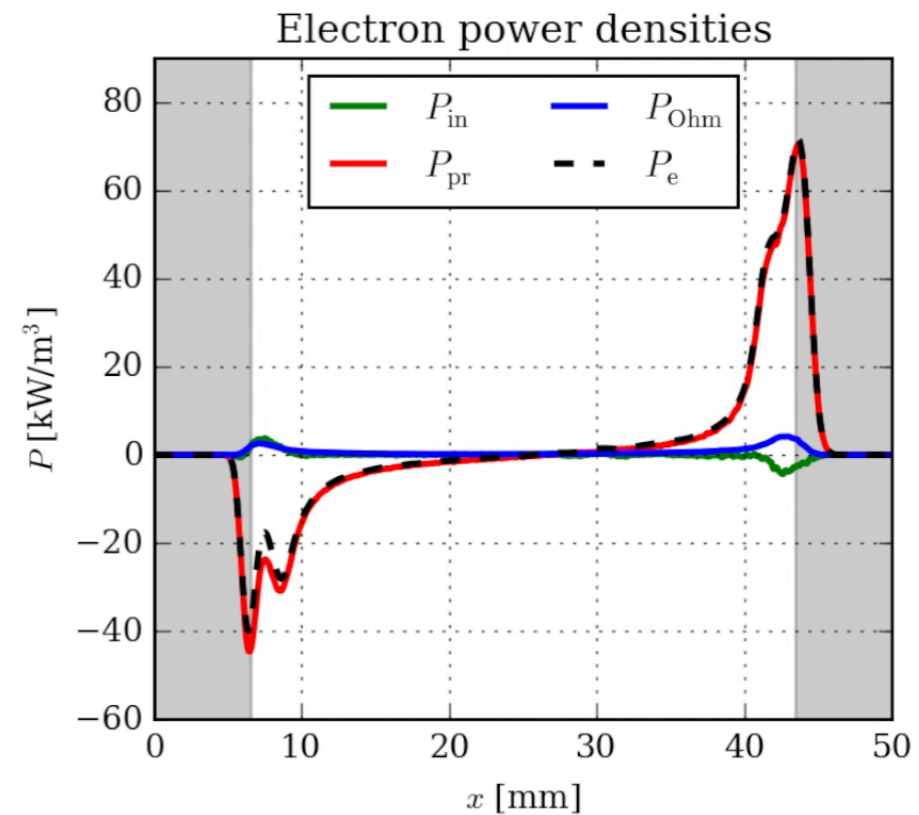
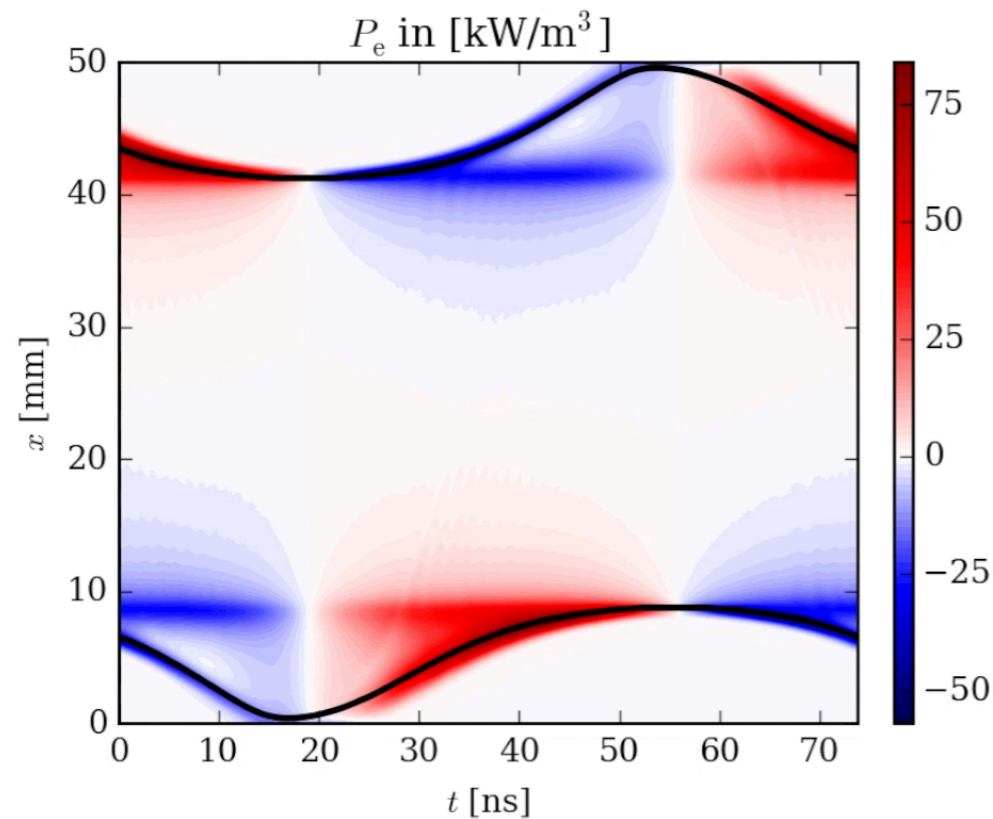


Boltzmann Term Analysis

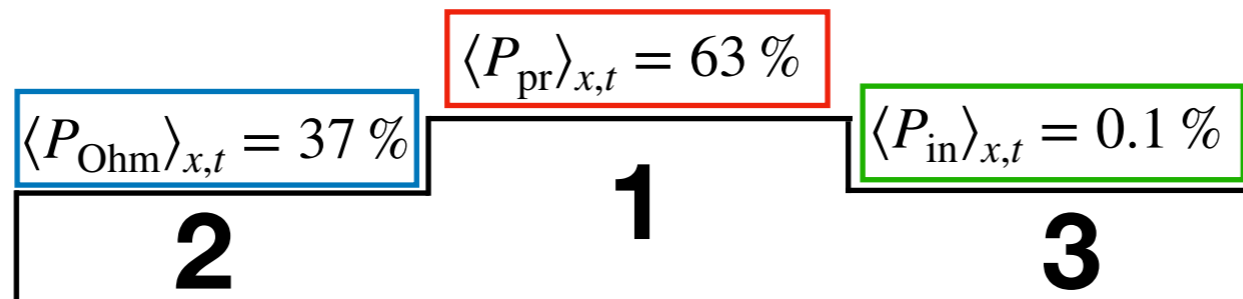
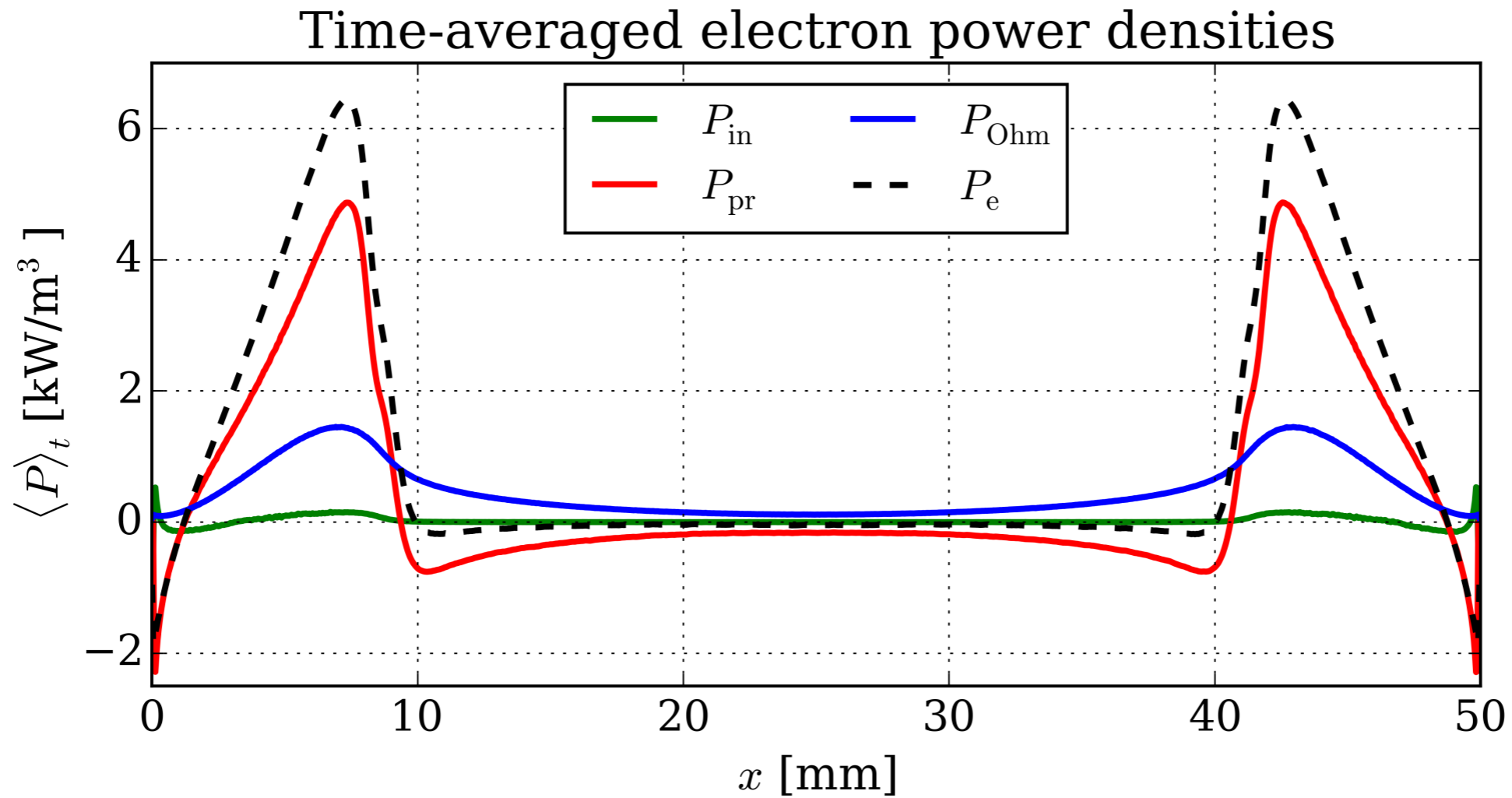
multiply by the
electron current:

$$\underbrace{j_e E}_{P_e} = \underbrace{j_e E_{in}}_{P_{in}} + \underbrace{j_e E_{pr}}_{P_{pr}} + \underbrace{j_e E_{Ohm}}_{P_{Ohm}}$$

$\underbrace{P_{in} + P_{pr}}_{P_{collisionless}} + \underbrace{P_{Ohm}}_{P_{collisional}}$

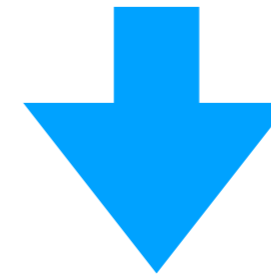
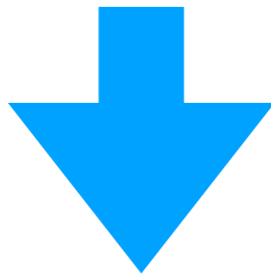


Boltzmann Term Analysis



Mechanism of Electron Power Absorption

more information about the electron power absorption by Máté Vass



Session PW2: Capacitively Coupled Plasmas II

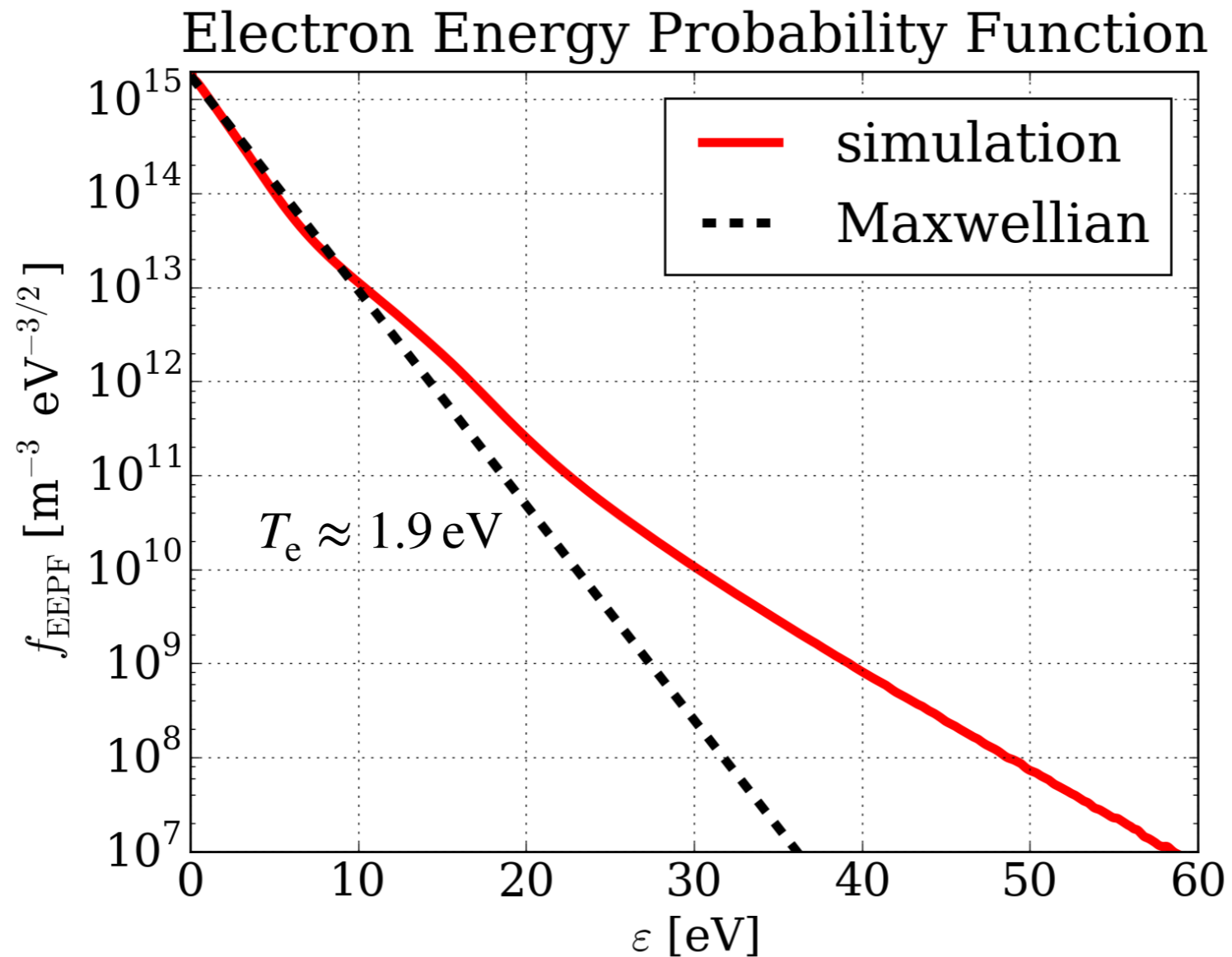
1:00 PM–2:30 PM, Wednesday, October 7, 2020

Chair: Xiaopu Li, Applied Materials

Abstract: PW2.00001 : Electron power absorption in capacitive RF plasmas based on a moment analysis of the Boltzmann equation

1:00 PM–1:30 PM **Live**

Electron Temperature



- electron temperature (thermodynamic relation): $T_e = \frac{2}{3} \langle \epsilon \rangle \approx 1.9 \text{ eV}$
- good approximation to represent the low energetic electrons (99% population)
- strong anisotropy for the high energetic electrons (1% population)
- provide a kinetic concept of the temperature to discuss the anisotropy

Electron Temperature

**momentum balance
in x-direction (parallel):**

$$m_e \frac{\partial(n_e u_{\parallel})}{\partial t} + m_e \frac{\partial(n_e u_{\parallel}^2)}{\partial x} + \frac{\partial(n_e T_{\parallel})}{\partial x} = -en_e E_{\parallel} - \Pi_{c\parallel}$$

$$\bar{\bar{p}} = \begin{pmatrix} p_{xx} & p_{xy} & p_{xz} \\ p_{yx} & p_{yy} & p_{yz} \\ p_{zx} & p_{zy} & p_{zz} \end{pmatrix} \Rightarrow \begin{pmatrix} p_{xx} & 0 & 0 \\ 0 & p_{yy} & 0 \\ 0 & 0 & p_{zz} \end{pmatrix} \Rightarrow \begin{pmatrix} p_{\parallel} & 0 & 0 \\ 0 & p_{\perp} & 0 \\ 0 & 0 & p_{\perp} \end{pmatrix}$$

Electron Temperature

**momentum balance
in x-direction (parallel):**

$$m_e \frac{\partial(n_e u_{\parallel})}{\partial t} + m_e \frac{\partial(n_e u_{\parallel}^2)}{\partial x} + \frac{\partial(n_e T_{\parallel})}{\partial x} = -en_e E_{\parallel} - \Pi_{c\parallel}$$

**momentum balance in
perpendicular-direction:**

~~$$m_e \frac{\partial(n_e u_{\perp})}{\partial t} + m_e \frac{\partial(n_e u_{\perp}^2)}{\partial x} + \frac{\partial(n_e T_{\perp})}{\partial x} = -en_e E_{\perp} - \Pi_{c\perp}$$~~

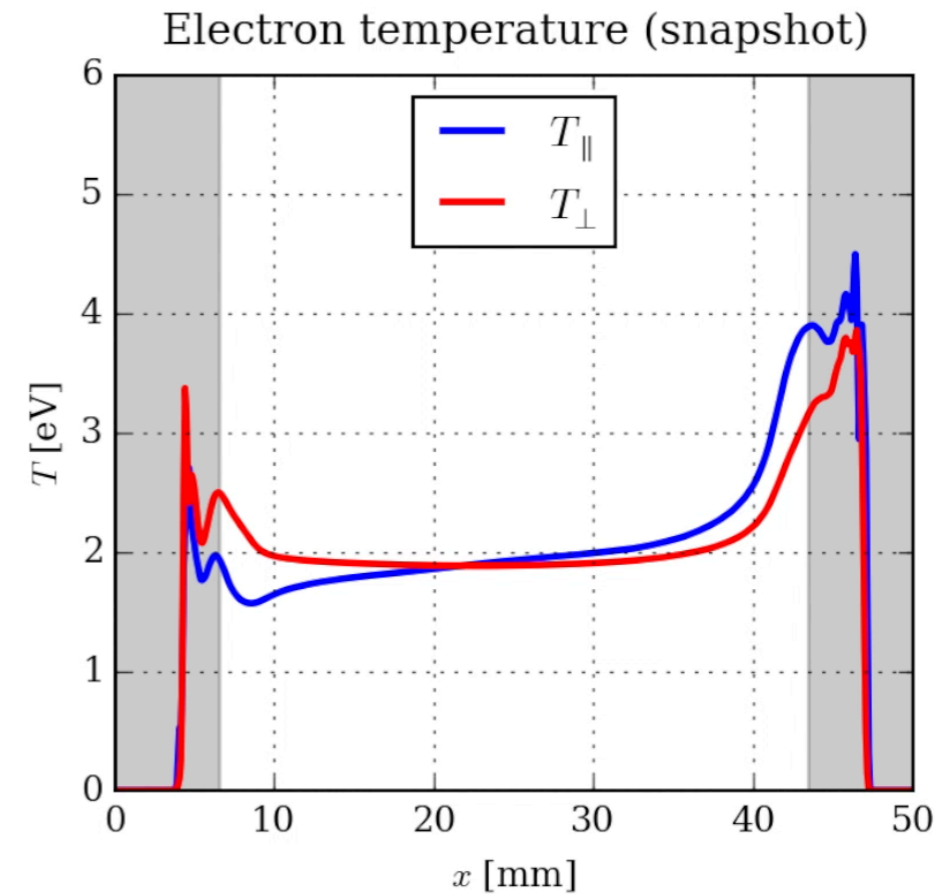
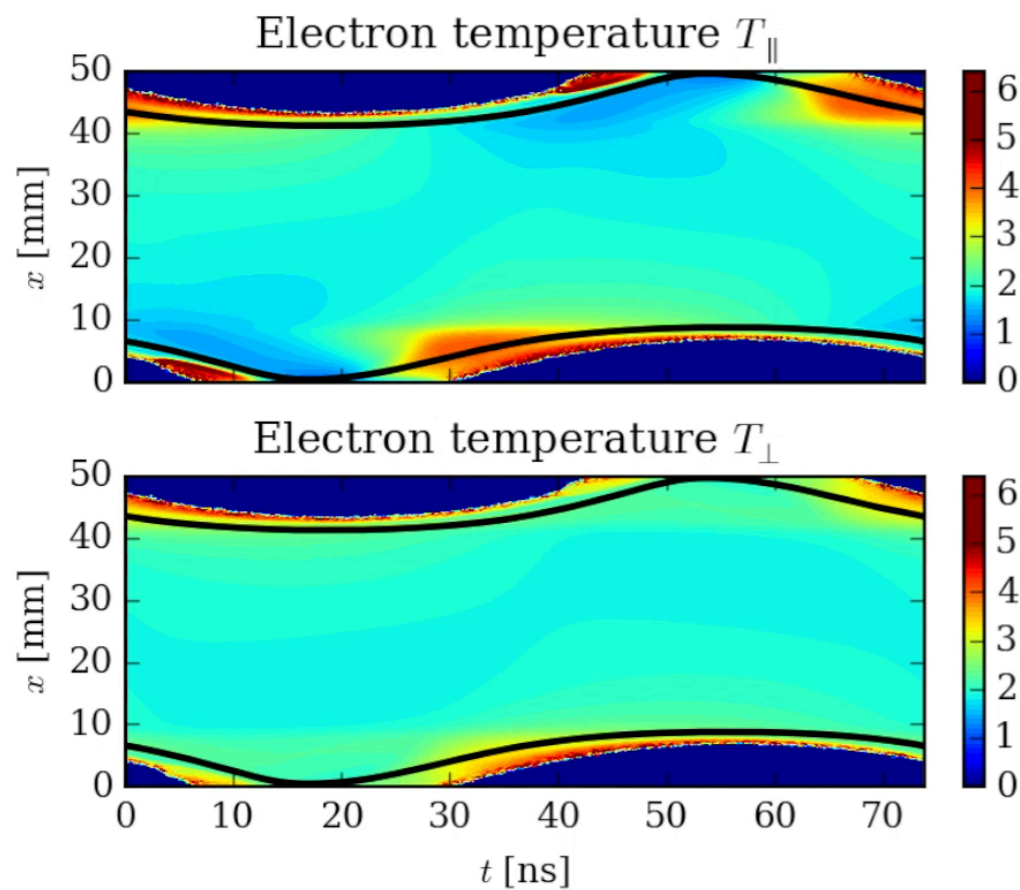


the parallel and perpendicular temperature communicate via collisions

Electron Temperature

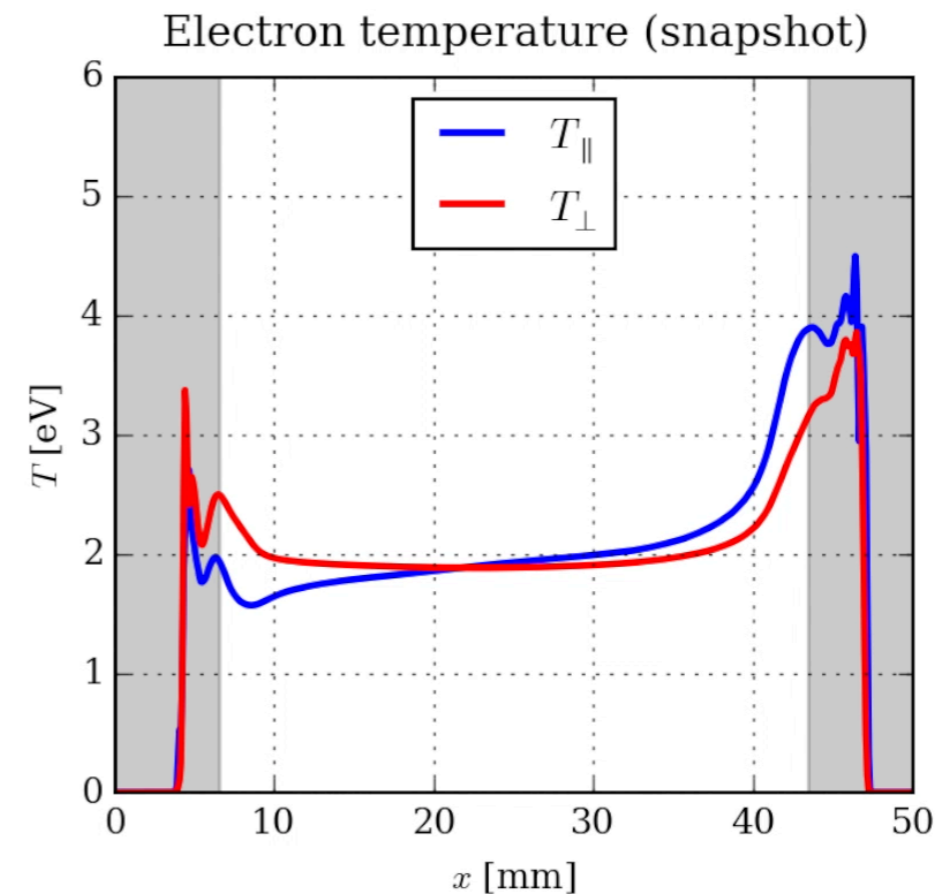
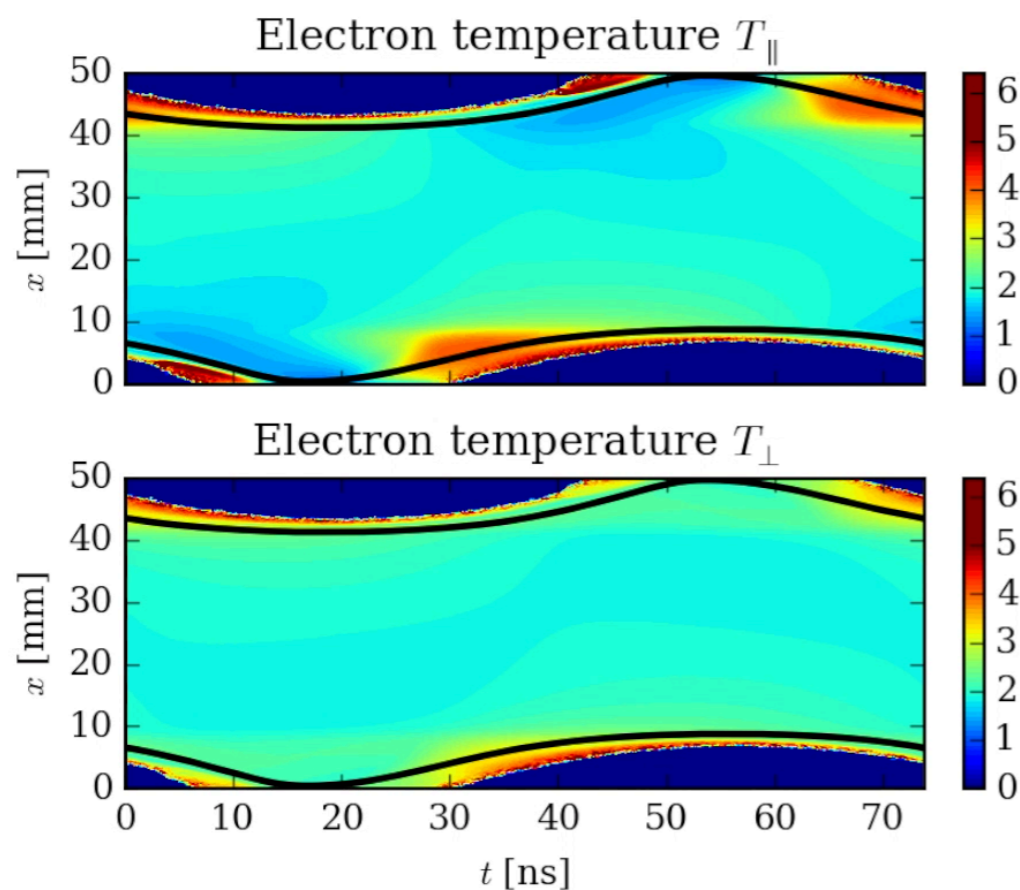
$$T_{\parallel} = \frac{p_{\parallel}}{n_e} = m_e \left(\langle v_{\parallel}^2(x, t) \rangle - u_{\parallel}^2(x, t) \right)$$

$$T_{\perp} = \frac{p_{\perp}}{n_e} = m_e \left(\langle v_{\perp}^2(x, t) \rangle - u_{\perp}^2(x, t) \right)$$



Electron Temperature

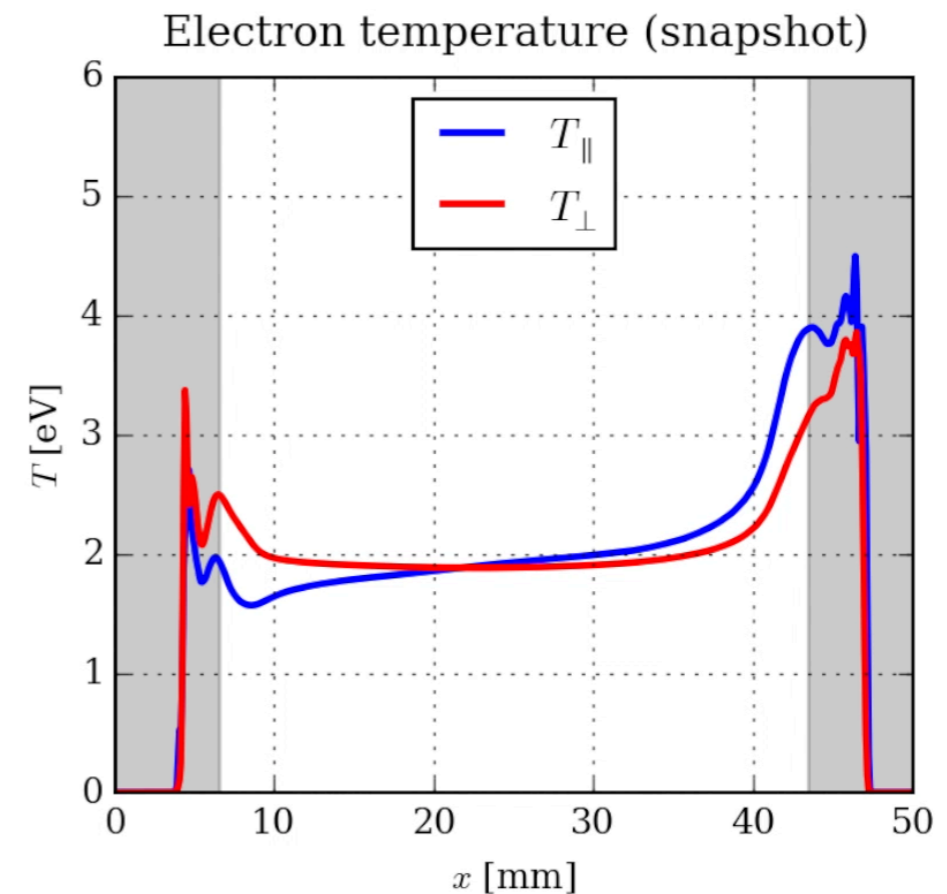
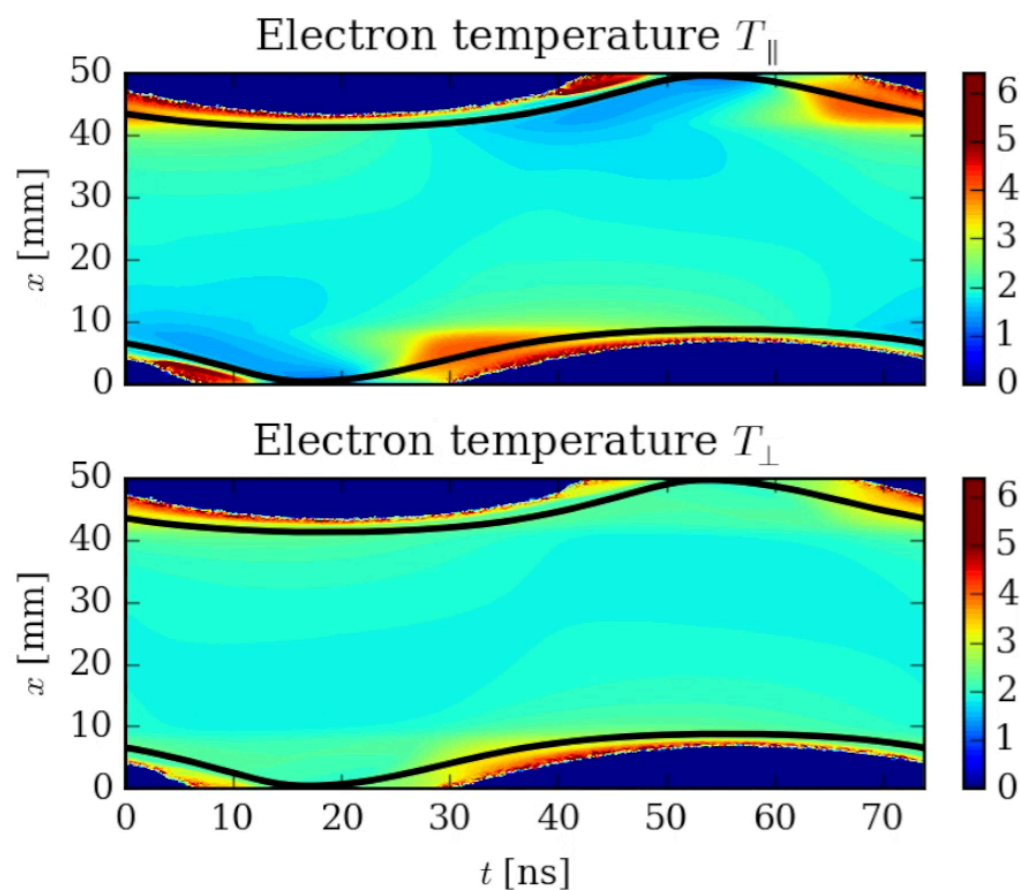
- almost isotropic in the center of the discharge ($T = 1.9$ eV)
- parallel electron temperature increases during sheath expansion
- perpendicular electron temperature temporally lags behind
- it needs a certain time to redistribute the energy due to collisions



Electron Temperature

- this concept of the electron temperature clearly shows the degree of anisotropy
- both temperatures act like an energy reservoir and contribute to the energy density

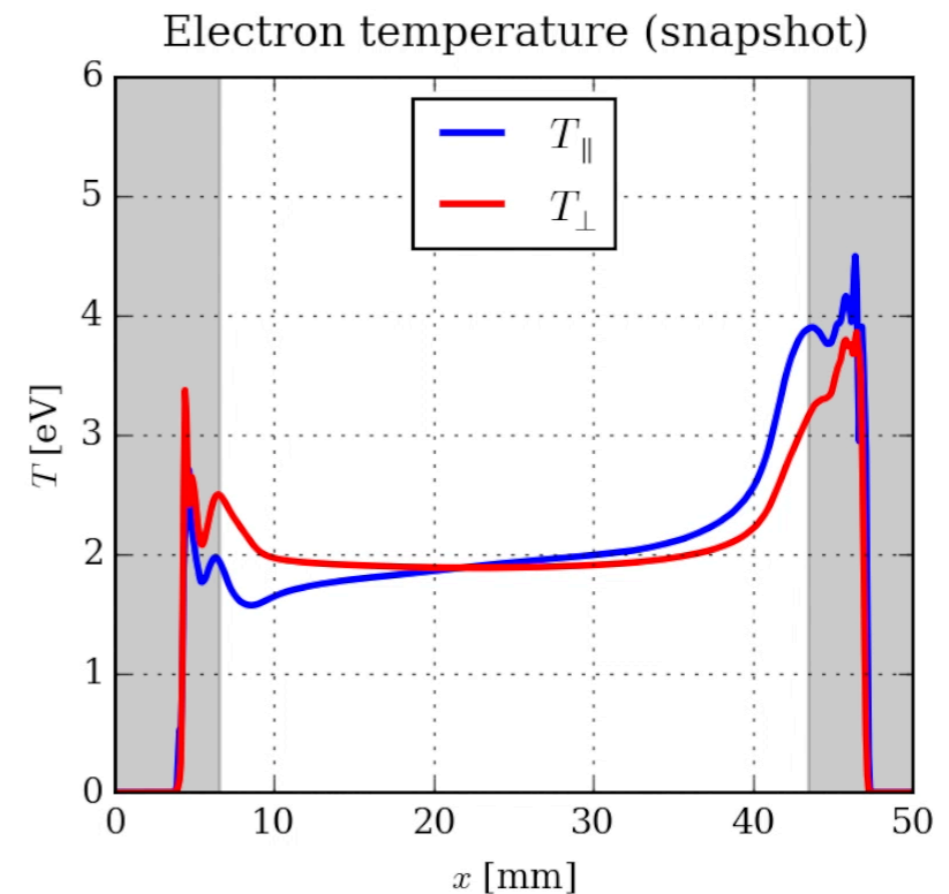
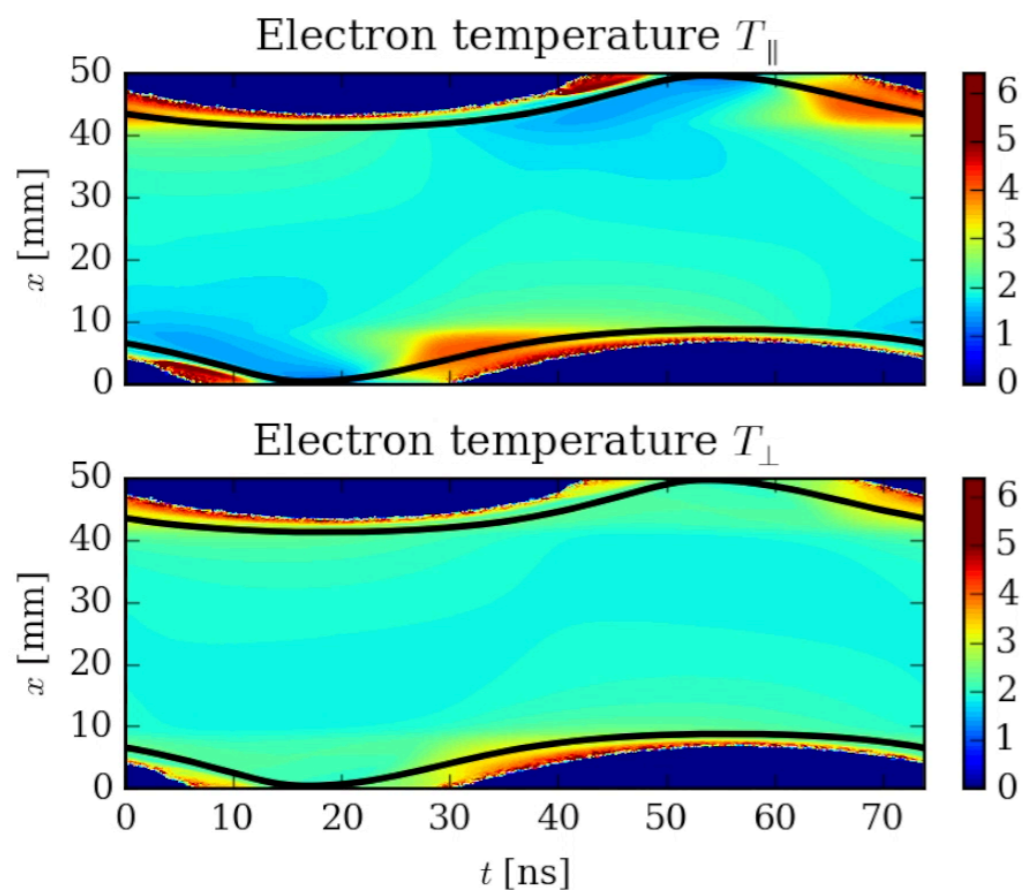
$$w = \frac{1}{2}n_e(m_e u_{\parallel}^2 + T_{\parallel} + 2T_{\perp})$$



Electron Temperature

- this concept of the electron temperature clearly shows the degree of anisotropy
- both temperatures act like an energy reservoir and contribute to the energy density

energy conservation:
$$\frac{\partial}{\partial t} w + \nabla \cdot \vec{Q} = P_{\text{tot}} - \varepsilon_c$$



Conclusion

- CCRF discharges at low pressures ($p < 10$ Pa), work in a very nonlocal regime
- the Boltzmann term analysis shows a coherent terminology of how to study the electron power gain and loss mechanism
- mostly the pressure heating term dominates at low pressures
- the concept of the kinetic electron temperature (parallel and perpendicular) indicates that electron power absorption and electron heating are physically two different mechanisms
- the difference of both temperatures demonstrates the degree of anisotropy

Poster Session Today:

**Abstract: LT2.00023 : Basic research of electron dynamics
in low pressure capacitively coupled plasmas**